



1 Dimensional matrix  $A$ :

	$d$	$\tau$	$\mu$	$u$	$r$
m	1	-1	-1	1	1
s	0	-2	-1	-1	0
kg	0	1	1	0	0

The  $d, \mu, u$  columns are linearly independent, so  $d, \mu, u$  are dimensional independent and core and possible core variables. By Buckingham's Pi-theorem there are

$$5 - \text{rank } A = 5 - 3 = 2$$

dimensional independent combinations, e.g.

$$\pi_1 = \frac{\tau}{d^a \mu^b u^c} = \dots = \frac{\tau d}{\mu u},$$
$$\pi_2 = \frac{r}{d}.$$

By Buckingham's Pi-theorem again, any dimensionally consistent relation

$$\tau = \Phi(d, \mu, u, r)$$

is equivalent to a relation

$$\Psi(\pi_1, \pi_2) = 0,$$

and solving for  $\pi_1$  we find that  $\pi_1 = C(\pi_2)$ , or  $\tau = C\left(\frac{d}{r}\right)\mu$ .

2 We introduce the scales  $c^* = Cc$ ,  $x^* = Xx$ ,  $t^* = Tt$ . The natural scales are  $C = M$  and  $X = L$ . The scaled equation is

$$\frac{M}{T}c_t = \frac{DM}{L^2}c_{xx} + \frac{UM}{L}c_x + rMc.$$

By assumption  $c_t, c_{xx}, c_x, c, x, t \sim 1$ , hence after dividing by  $M$ ,

$$\frac{1}{T} \sim \frac{D}{L^2} + \frac{U}{L} + r = \frac{1}{L^2}(D + uL + rL^2).$$

Case (i) :  $D \gg UL + rL^2$

The time scale is then given by  $\frac{1}{T} \sim \frac{D}{L^2}$ , and the scaled equation is

$$c_t = C_{xx} + \varepsilon_1 c_x + \varepsilon_2 c,$$

where  $\varepsilon_1 = \frac{UT}{L} = \frac{UL}{D} \ll 1$  and  $\varepsilon_2 = rT = \frac{rL^2}{D} \ll 1$ .

Case (ii) :  $rL^2 \gg D + UL$

The time scale is given by  $\frac{1}{T} \sim r$ , and the scaled equation is

$$c_t = \varepsilon_1 c_{xx} + \varepsilon_2 c_x + c,$$

where  $\varepsilon_1 = \frac{Dr}{L^2} \ll 1$  and  $\varepsilon_2 = \frac{Ur}{L} \ll 1$ .

3 (i) could model two competing populations of size  $n_1$  and  $n_2$ . When  $n_2 = 0$ , the  $n_1$ -equation is the logistic equation with bounded population, while if  $n_1 = 0$ , the  $n_2$ -equation gives exponential growth.

(ii) could model two populations that are mutually beneficial, e.g. as in symbiosis.

We find equilibrium points given by

$$\mathbf{F} = \frac{d}{dt} \begin{pmatrix} n_1 \\ n_2 \end{pmatrix} = \begin{pmatrix} n_1(1 - n_1 - an_2) \\ cn_2(1 - n_1) \end{pmatrix} = 0,$$

that is,  $(n_1, n_2) \in \{(0, 0), (1, 0)\}$ . We examine the stability of the equilibrium points. To that end consider the Jacobi matrix of  $\mathbf{F}$ ,

$$D\mathbf{F}(n_1, n_2) = \begin{bmatrix} 1 - 2n_1 - an_2 & -an_1 \\ -cn_2 & c(1 - n_1) \end{bmatrix}.$$

By inserting the first equilibrium point we get

$$D\mathbf{F}(0, 0) = \begin{bmatrix} 1 & 0 \\ 0 & c \end{bmatrix},$$

with at least one positive eigenvalue ( $\lambda = 1$ ). This implies that  $(0, 0)$  is unstable. The second equilibrium point gives

$$D\mathbf{F}(1, 0) = \begin{bmatrix} -1 & -a \\ 0 & 0 \end{bmatrix},$$

where  $\max\{\text{Re}(\lambda)\} = 0$  and no conclusion.

4 This is a singular perturbation problem. The only boundary point is  $t = 0$ , and the boundary layer is there. We solve the problem as follows.

- (i) Find the outer solution.
- (ii) Find the boundary layer thickness and the rescaled equations.
- (iii) Find the inner solution.
- (iv) Matching in the intermediate region.
- (v) The uniform approximation.

Implementing the above strategy.

- (i) The outer solution solves the equations with  $\varepsilon = 0$ ,

$$\begin{aligned} \dot{x}_O &= -x_O + (x_O + 1)(y_O - 1), \\ 0 &= x_O - (x_O + 1)y_O. \end{aligned}$$

The solution is

$$\begin{aligned}x_O(t) &= Ce^{-t} - (1 - e^{-t}), \\y_O(t) &= \frac{x_O(t)}{x_O(t) + 1},\end{aligned}$$

where we have used the technique of integrating factor to determine  $x_O(t)$ .

- (ii) We rescale the equation  $t = \delta\tau$ ,  $X(\tau) = x(t)$  and  $Y(\tau) = y(t)$ . Inserting into the equations,

$$\begin{aligned}\frac{1}{\delta} \frac{d}{d\tau} X &= -X + (X + 1)(Y - 1), \\ \frac{\varepsilon}{\delta} \frac{d}{d\tau} Y &= X - (X + 1)Y.\end{aligned}$$

If we balance terms in the first equation we get  $\delta = 1$  and recover the outer solution. If we, on the other hand, balance terms in the second equation we get  $\delta = \varepsilon$ . Hence  $\delta = \varepsilon$  is the other time scale.

- (iii) The rescaled equations are

$$\begin{aligned}\dot{X} &= \varepsilon(-X + (X + 1)(Y - 1)), & X(0) &= 1 \\ \dot{Y} &= X - (X + 1)Y, & Y(0) &= 0.\end{aligned}$$

We get the inner solution by letting  $\varepsilon = 0$ , hence

$$\begin{aligned}X_I(t) &= 1, \\ Y_I(t) &= \frac{1}{2}(1 - e^{-2t}).\end{aligned}$$

- (iv) The matching condition is

$$\lim_{t \rightarrow 0^+} (x_O(t), y_O(t)) = \lim_{\tau \rightarrow \infty} (X_I(\tau), Y_I(\tau)).$$

We calculate the limits using the solutions and get that  $C = 1$ .

- (v) The uniform approximation is calculated by summing the inner and outer solution and subtracting the intermediate value. Thus

$$\begin{aligned}x_u(t) &= 2e^{-t} - 1, \\ y_u(t) &= 1 - \frac{1}{2}e^t - \frac{1}{2}e^{-2\frac{t}{\varepsilon}}.\end{aligned}$$

- 5** The kinematic speed is

$$c(\rho) = j'(\rho) = 1 - 2\rho,$$

and the characteristics are given by

$$\begin{aligned}\dot{x} &= c(z), & x(t_0) &= x_0, \\ \dot{z} &= 0, & z(t_0) &= \rho(x_0, t_0).\end{aligned}$$

The solution of this system is

$$\begin{aligned}z(t) &= \rho(x_0, t_0), \\ x(t) &= x_0 + (t - t_0)c(\rho(x_0, t_0)).\end{aligned}$$

In  $x \geq 0$ , inflow boundary corresponds to  $c \geq 0$ . To find a boundary condition at  $x = 0$ , we convert to Dirichlet condition

$$\frac{3}{16} = j(\rho) = \rho(1 - \rho),$$

which gives  $\rho = \frac{1}{4}$  or  $\rho = \frac{3}{4}$ . Since  $c(\frac{1}{4}) > 0$  and  $c(\frac{3}{4}) < 0$ , and we can only impose boundary conditions at inflow, we take

$$\rho = \frac{1}{4} \text{ at } x = 0, t > 0.$$

The characteristics are

$$x(t) = \begin{cases} x_0 + tc\left(\frac{1}{8}\right), & t_0 = 0, x_0 \geq 0, t > 0, \\ (t - t_0)c\left(\frac{1}{4}\right), & x_0 = 0, t > t_0 > 0, \end{cases}$$

$$z(t) = \begin{cases} \frac{1}{8}, & t_0 = 0, x_0 \geq 0, t > 0, \\ \frac{1}{4}, & x_0 = 0, t > t_0 > 0. \end{cases}$$

In the dead sector there is a rarefaction wave solution.

**6** a) Conservation of mass in  $R$

$$(1) \quad \frac{d}{dt} \int_R \rho \, dV = \int_{\partial R} \rho \mathbf{v} \cdot \mathbf{n} \, d\sigma.$$

Observe that

$$\int_R \rho \, dV = \rho \int_0^1 dy \int_{\bar{z}}^{\bar{z} + \Delta z} \int_0^{d(z, \bar{t})} dx dz,$$

and since  $d$  is smooth,

$$\frac{d}{dt} \int_R \rho \, dV = \rho \int_{\bar{z}}^{\bar{z} + \Delta z} d_t(z, \bar{t}) \, dz.$$

As  $\mathbf{v} \cdot \mathbf{n} = -u \mathbf{e}_z \cdot \mathbf{n}$ ,

$$\begin{aligned} \int_{\partial R} \rho \mathbf{v} \cdot \mathbf{n} \, d\sigma &= \int_{\partial R \cap \{z = \bar{z}\}} \rho(-u)(-1) \, d\sigma \\ &\quad + \int_{\partial R \cap \{z = \bar{z} + \Delta z\}} \rho(-u)(-1) \, d\sigma \\ &= \rho((ud)(\bar{z}) - (ud)(\bar{z} + \Delta z)). \end{aligned}$$

If we divide (1) by  $\rho \Delta z$ , we find

$$\frac{1}{\Delta z} \rho \int_{\bar{z}}^{\bar{z} + \Delta z} d_t \, dz = \frac{\rho((ud)(\bar{z}) - (ud)(\bar{z} + \Delta z))}{\Delta z}.$$

Let  $\Delta z \rightarrow 0$ , use definition of derivative and the fundamental theorem of calculus to see that

$$\frac{\partial}{\partial t} d = \frac{\partial}{\partial z} (ud).$$

- b) Fix  $\bar{t} > 0$  and  $\bar{z} \in [0, 1)$ , and let  $\Delta z > 0$  be such that  $\bar{z} + \Delta z < 1$ . If  $d(\bar{z}, \bar{t}) = 0$  then

$$\int_R -\rho g \mathbf{e}_z dV = - \int_{R \cap \{x=0\}} C \frac{\mu u}{d} \mathbf{e}_z d\sigma.$$

This implies that

$$\rho g d(\bar{z}, \bar{t}) = C \mu \left( \frac{u}{d} \right) (\bar{z}, \bar{t}),$$

and hence

$$u = \frac{\rho g}{C \mu} d^2 = K d^2.$$

By a),  $d$  then satisfies

$$d_t - (K d^3)_z = 0.$$

Method of characteristics ( $z(t) = d(x(t), t)$ )

$$\begin{aligned} \dot{x} &= -2K z^2 = c(z), & x(0) &= x_0, \\ \dot{z} &= 0, & z(0) &= d(x_0, 0). \end{aligned}$$

The characteristics are then

$$\begin{aligned} x &= x_0 + t c(d(x_0, 0)), \\ z &= d(x_0, 0). \end{aligned}$$

Since  $z_1 < z_2$  implies that  $d(z_1, 0) < d(z_2, 0)$ , which again implies that  $c(d(z_1, 0)) > c(d(z_2, 0))$ , the characteristics will develop shock. The shock speed via the Rankine-Hugoniot condition is

$$\dot{s} = -2K \frac{\rho^2(z_1, 0) - \rho^2(z_2, 0)}{\rho(z_1, 0) - \rho(z_2, 0)} < 0$$

so the shock moves downwards.