CASE STUDY FROM BIO-MATHEMATICAL MODELLING

A PHYSIOLOGICAL FLOW PROBLEM

(After Lin and Segel, Chapter 8)

How is water and salt expelled from the body, e.g. from the kidneys?

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J. Diamond (1967): Salt is expelled from the body in a *non-direct way* by means of so-called *secondary channels*, which are consistently found in fluid secreting tissue.

• At the inner end of the secondary channels *chemical pumps* enter salt into the channel leading to a local high concentration of salt, and a salt concentration gradient towards the opening of the channels

- Water enters the channel by osmosis through the walls
- Salt is moving in the channel by diffusion and convection
- At the outer end of the channel, the salt concentration is C_0 (body average)

What determines the amount of salt coming out of the secondary channel under stationary (time independent) conditions?

QUALITATIVE MODEL



PHYSICAL MECHANISMS

OSMOSIS:



Salt ions are too large to pass through the membrane!

If the ion concentrations on each side of the membrane are C_1 and C_2 , then the amount of water passing through the membrane per area and time units is

$$J = P(C_2 - C_1)$$

The constant *P* is called the *permeability*.

$$J = P(C_2 - C_1)$$

Units:

$$[J] = \frac{Volume}{Area \times Time} = \frac{m^3}{m^2 s} = \frac{m}{s} \quad \text{(Same unit as velocity!)}$$
$$[C] = \frac{\# \text{ions}}{\text{Volume}} = \frac{osmol}{m^3}, \text{ (osmol = Avogadro's number)}$$
$$[P] = \frac{m/s}{osmol/m^3} = \frac{m^4}{osmol \cdot s}$$

(Note since the concentration of ions is about twice the concentration of salt, we may just as well think of *C* as the salt concentration)

DIFFUSION

Motion of salt in an otherwise stationary solution is due to concentration differences:

$$F = -D\frac{\partial C}{\partial x}$$

F is called the *flux* of salt (in the *x*-direction).

 $\partial C / \partial x$ is the *concentration gradient*.

D is called the *diffusion coefficient*.

Flux (in general a vector quantity!) is *amount passing through an imaginary surface in the fluid per time and area unit.* Units:

$$[F] = \frac{amount}{area \times time} = \frac{\text{osmol}}{\text{m}^2 \text{s}}$$
$$[C] = \frac{amount}{volume} = \frac{\text{osmol}}{\text{m}^3}$$
$$[D] = \frac{\text{m}^2}{\text{s}}$$

CHEMICAL PUMPS

Enter salt into the solution from the tissue by means of certain chemical mechanisms using energy (details not known!)

$$[N_0] = \frac{Amount}{Unit \ area \ wall \times time} = \frac{\text{osmol}}{\text{m}^2 \text{s}}$$

THE NEXT STEP IS DEFINING THE GEOMETRY:

- The channels are *long and narrow*. Thus, we consider a 1d model.
- The inner end of the channel is closed.



Water entering the channel through osmosis (per area unit):

 $J = P[C^*(x^*) - C_0]$

Transport of salt in the channel:

A. Diffusion:
$$F_D = -DdC^*/dx^*$$

B. Convection: Passive transport due to the motion of the fluid



$$F^* = F_D + F_c = v * C * -D\frac{dC *}{dx *}$$

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Chemical pumps (we do not include the end): $N_0(\delta c)$

The problem presented to us is to determine the *emergent osmolarity* defined by $F^*(L) = v^*(L)Os^*$:

$$Os^* = \frac{F^*(L)}{v^*(L)} = \frac{v^*(L)C^*(L) - D\frac{dC^*}{dx^*}(L)}{v^*(L)}$$

Emergent osmolarity = the *apparent* concentration needed in order to have the same amount of output with the given fluid velocity and *no contribution from diffusion* (fictitious channel with constant concentration)

THE EQUATIONS: CONSERVATION OF WATER AND SALT!

WATER (Since the salt concentration is small, the density of water is assumed to be constant)



$$Q_{out} - Q_{in} = Q_{os}$$

Av*(x*+\Delta x*) - Av*(x*) = P(C*(x*+\delta x*) - C_0)

$$(\Delta x^* \to 0) \Rightarrow \frac{dv^*}{dx^*} = \frac{Pc}{A} (C^*(x^*) - C_0)$$

SALT

We use a conservation argument similar to for water:

$$Q_{out}^{salt} = AF^*(x^* + \Delta x^*)$$
$$Q_{in}^{salt} = AF^*(x^*)$$
$$Q_{ch.p.}^{salt} = N^*(x^*) \cdot (c\Delta x^*)$$

By letting Δx^* tend to 0:

$$A\frac{dF^*}{dx^*} = \begin{cases} N_0c, & x^* \le \delta\\ 0, & x^* > \delta \end{cases}$$

The previous simple equation (for the salt) may be solved immediately: $x^* \le \delta$:

$$A\frac{dF^*}{dx^*} = N_0 c \implies F^* = \frac{N_0 c}{A} x^* \quad (\text{Since } F^*(0) = 0)$$

 $x^* > \delta$:

$$A\frac{dF^*}{dx^*} = 0 \implies F^* = \frac{N_0c}{A}\delta$$
 (Since F^* is continuous)

Together, with the expression for the salt flux, this gives the following differential equation for C^* :

$$C^* v^* - D\frac{dC^*}{dx^*} = \begin{cases} \frac{cN_0}{A} x^*, & 0 \le x^* \le \delta \\ \frac{cN_0}{A} \delta, & \delta \le x^* \le L \end{cases}$$

BOUNDARY CONDITIONS AND MATCHING CONDITIONS AT $x^*=\delta$:

Closed end channel:

$$v^*(0) = 0,$$

 $F^*(0) = 0.$

In fact,

$$\frac{dC^*}{dx^*}(0) = 0 \quad \left(\text{Since } F^*(0) = C^*(0)v^*(0) - D\frac{dC^*}{dx^*}(0) \right)$$

At the right end we assume that $C^*(L) = C_0$.

Finally, *matching* conditions at $x^* = \delta$:

$$F^*(\delta +) = F^*(\delta -)$$

$$v^*(\delta +) = v^*(\delta -)$$

$$C^*(\delta +) = C^*(\delta -)$$

$$\frac{dC^*}{dx^*}(\delta +) = \frac{dC^*}{dx^*}(\delta -) \left(\text{Since } F^* = C^*v^* - D\frac{dC^*}{dx^*}\right)$$

(Not all conditions are needed in the final formulation)

FINAL FORMULATION

Differential equations (non-linear and coupled!):

$$\frac{dv^*}{dx^*} = \frac{Pc}{A} \left(C^* (x^*) - C_0 \right), \ 0 \le x^* \le \delta$$
$$C^* v^* - D \frac{dC^*}{dx^*} = \begin{cases} \frac{cN_0}{A} x^*, & 0 \le x^* \le \delta\\ \frac{cN_0}{A} \delta, & \delta \le x^* \le L \end{cases}$$

Boundary conditions:

$$v^{*}(0) = 0, \quad C^{*}(L) = C_{0}$$

 v^{*}, C^{*} continuous for $x^{*} = \delta$

Determine:

$$Os^* = \frac{F^*(L)}{v^*(L)} = \frac{cN_0\delta}{Av^*(L)}$$

SCALING

| Parameter | Unit | Min. value | Typical value | Max. value |
|-----------|-------------------------|--------------------|--------------------|--------------------|
| r | cm | 10-6 | 5×10 ⁻⁶ | 10-4 |
| L | cm | 4×10 ⁻⁴ | 10 ⁻² | 2×10 ⁻² |
| δ | cm | 4×10 ⁻⁵ | 10-3 | 2×10 ⁻³ |
| D | cm ² /s | 10-6 | 10-5 | 5×10 ⁻⁵ |
| N_0 | mOsm/cm ² s | 10 ⁻¹⁰ | 10-7 | 10-5 |
| Р | cm ⁴ /s mOsm | 10-6 | 2×10 ⁻⁵ | 2×10 ⁻⁴ |
| C_0 | mOsm/cm ³ | - | 3×10 ⁻¹ | |

(Table from Lin & Segel, p. 264)

Length scale: δ Concentration scale: C_0 Velocity scale: $U = \frac{cN_0\delta}{C_0A}$ Derivation of the velocity scale:

$$\frac{cN_0\delta}{\underline{A}}A = C_0AU \implies U = \frac{cN_0\delta}{AC_0}$$

Salt out per time unit

Scaled variables:

$$x^* = \delta x$$

$$C^* = C_0 C$$

$$v^* = Uv = \frac{cN_0\delta}{AC_0}v$$

THE SCALED EQUATIONS

$$\varepsilon \frac{dv}{dx} = C - 1, \ 0 \le x \le \lambda$$

$$Cv - \eta \frac{dC}{dx} = \begin{cases} x & 0 \le x \le 1\\ 1 & 1 \le x \le \lambda \end{cases}$$

$$v(0) = 0, \ C(\lambda) = 1$$

$$v, C, dC / dx \text{ continuous at } x = 1.$$

$$\varepsilon = \frac{N_0}{PC_0^2}, \ \eta = \frac{AC_0D}{N_0\delta^2c}, \ \lambda = \frac{L}{\delta}$$
Determine $Os = \frac{1}{v(\lambda)}$

Impossible to solve analytically!

Typical values of the non-dimensional parameters:

| Parameter | Minimum | Typical | Maximum |
|-----------|--------------------|--------------------|------------------|
| ε | 10 ⁻⁵ | 2.10 ⁻² | 10 ² |
| η | 4.10 ⁻³ | 75 | 10 ¹⁰ |
| λ | 10(?) | 10 | 500 |

SOME NUMERICAL SOLUTIONS

- obtained by means of a simple shooting method for $\lambda = 10, \quad \kappa = 1, \quad \varepsilon = .5$ (The parameter κ will be defined below) 1.04 1.02 С 1 0.98 2 0 4 6 8 10 X 0.4 0.3 V 0.2 0.1 0 2 4 6 8 10 0 Χ (Note: C(0) = 1.0375)

MATLAB CODE:

```
x0 =[0, 1.0355];
xspan = [0,10];
[t,x]=ode23('kidney',xspan,x0);
subplot(2,1,1); plot(t,x(:,2));
subplot(2,1,2); plot(t,x(:,1));
```

```
%L&S Eqn. 16a,b, Section 8.3: v=x(1), C = x(2):
function xdot = kidney(t,x)
lambda = 10.;
kappa = 1;
eps = .5;
f = (eps*kappa^2/lambda^2);
xdot(1)=( x(2)-1 )/eps;
xdot(2)=f*x(1)*x(2)-f*min(t,1);
```

$$\lambda = 10, \quad \kappa = 1, \quad \varepsilon = 0.1:$$

C(0) = 1.0375 (*Misses the target* C(1) = 1!):



C(0)=1.0173 (*Hits the target* C(1)=1!):



A FIRST-TRY PERTURBATION EXPANSION $\varepsilon \frac{dv}{dx} = C - 1, \ 0 \le x \le \lambda$ $Cv - \eta \frac{dC}{dx} = \begin{cases} x & 0 \le x \le 1\\ 1 & 1 \le x \le \lambda \end{cases}$ $C = C_0 + \varepsilon C_1 + \varepsilon^2 C_2 + \dots$ $v = v_0 + \varepsilon v_1 + \varepsilon^2 v_2 + \dots$ To order ε^0 : $C_0 - 1 = 0 \Longrightarrow C_0 = 1$ $1v_0 + \eta \frac{dC_0}{dr} = \begin{cases} x \\ 1 \Rightarrow v_0 = \begin{cases} x \\ 1 \end{cases},$ To order ε^1 :

$$C_1 = \frac{dv_0}{dx} = \begin{cases} 1\\ 0 \Rightarrow C_1 \text{ is discontinuous!} \end{cases}$$

This is against our (physical) requirement that C should be continuous.

The problem is that η tends to be large when ε is small!

We introduce the modified parameter κ and eliminate η :

$$\eta = \frac{\left(\lambda^2 / \kappa^2\right)}{\varepsilon}$$

(The somewhat strange form is basically for making the expressions simpler!)

MODIFIED PERTURBATION EXPANSION $C - 1 = \varepsilon \frac{dv}{dx}, \quad 0 \le x \le \lambda$ $\varepsilon \kappa^{2} C v - \lambda^{2} \frac{dC}{dx} = \varepsilon \kappa^{2} \begin{cases} x, & 0 \le x \le 1\\ 1, & 1 \le x \le \lambda \end{cases}$ $\varepsilon <<1, \quad \kappa, \lambda \sim O(1)$

$$C = C_0 + \varepsilon C_1 + \varepsilon^2 C_2 + \dots$$
$$v = v_0 + \varepsilon v_1 + \varepsilon^2 v_2 + \dots$$

To order ε^0 :

$$\begin{array}{c} C_0 - 1 = 0 \\ \lambda^2 C_0 = 0 \end{array} \Longrightarrow C_0 = 1,$$

To order ε^1 :

$$\begin{bmatrix} C_1 = v_0' \\ \kappa^2 v_0 C_0 - \lambda^2 C_1' = \kappa^2 \begin{pmatrix} x \\ 1 \end{pmatrix} \end{bmatrix} \Rightarrow \kappa^2 v_0 - \lambda^2 v_0'' = \kappa^2 \begin{pmatrix} x \\ 1 \end{pmatrix}$$

Solution for
$$0 \le x \le 1$$
:
 $v_0 = x - K_1 \sinh\left(\frac{\kappa}{\lambda}x\right)$
 $C_1 = v_0' = 1 - K_1 \frac{\kappa}{\lambda} \cosh\left(\frac{\kappa}{\lambda}x\right)$

Solution for $1 \le x \le \lambda$:

$$v_0 = 1 - K_2 \cosh\left(\frac{\kappa}{\lambda}x + \Psi\right)$$
$$C_1 = v_0' = -K_2 \frac{\kappa}{\lambda} \sinh\left(\frac{\kappa}{\lambda}x + \Psi\right)$$

(since $C_1(\lambda) = 0$, $\Psi = -\kappa$)

The constants K_1 and K_2 are determined from the continuity requirements at x = 1:

$$K_{1} = \frac{\lambda}{\kappa} \frac{\cosh(\kappa / \lambda - \kappa)}{\cosh(\kappa)}$$
$$K_{2} = \frac{\lambda}{\kappa} \frac{\sinh(\kappa / \lambda)}{\cosh(\kappa)}$$

The dimensionless osmolarity at $x = \lambda$:

$$Os = \frac{1}{v(\lambda)} \approx \frac{1}{v_0(\lambda)} = \frac{1}{1 - K_2}$$

Assume $\kappa/\lambda < 1$:

$$Os \approx \frac{1}{1 - K_2} \approx \frac{1}{1 - \frac{1}{\cosh(\kappa)}} = \frac{\cosh(\kappa)}{\cosh(\kappa) - 1}$$

When κ is small compared to 1, then

$$Os \approx \frac{\cosh(\kappa)}{\cosh(\kappa) - 1} \approx \frac{1 + \kappa^2 / 2 + \dots}{\kappa^2 / 2 + \dots} \approx \frac{2}{\kappa^2}$$



ALL DIMENSIONLESS PARAMETERS MEAN SOMETHING IMPORTANT AND USEFUL!

$$\frac{\kappa^2}{2} = \frac{cPC_0L^2}{AD2} \frac{\overline{C} - C_0}{\overline{C} - C_0} = \frac{\left(cLP(\overline{C} - C_0)\right)C_0\frac{1}{A}}{D\left(\overline{C} - C_0\right)/(L/2)}$$

Now,

$$(cLP(\overline{C}-C_0))C_0\frac{1}{A}\approx F_{conv.}$$

water in by osmosis

$$F_{diff} \approx -D \frac{dC^*}{dx^*} \Big|_{ave} \approx D \frac{\overline{C} - C_0}{L/2}$$

That is,

$$\frac{\kappa^2}{2} \approx \frac{F_{conv.}}{F_{diff.}}$$

Thus,

$$Os = \frac{Os^*}{C_0} = \frac{F^*(L)}{v^*(L)C_0} = \frac{F_{conv.} + F_{diff}}{F_{conv.}} = 1 + \frac{F_{diff}}{F_{conv.}} \approx 1 + \frac{2}{\kappa^2}$$

THE STEPS OF THE ANALYSIS:

- 1. A qualitative idea of the model
- 2. Identify physical mechanisms
- 3. Sort out the geometry
- 4. The modelling: Based on conservation principles!
- 5. Scaling
- 6. Perturbation analysis (not working!)
- 7. Modified scaling
- 8. Modified perturbation analysis (working!)
- 9. Numerical experiments
- 10. Final analysis
- 11. Conclusions