

Introduction: Slides (general information)

## 1 Dimensional analysis (TMA4195, supplementary notes, chap 1)

Two basic facts about physical (and geometric) relations:

- Come back to this
- \* All relations must be dimensionally correct.
  - \* All relations are valid independent of choice of (fundamental) units.

A physical quantity  $R$  has

- ∞ dimension
- ∞ unit
- ∞ numerical value

Dimension: Length (L), Time (T), Mass (M), Energy ( $ML^2T^{-2}$ ), ...

Unit: meter, feet, lightyear, months, Joule, Ampere, ...

$$R = 1m = 10^{-3} km = 39.37'' = 1.06 \cdot 10^{-15} \text{ lightyears}$$

In general:  $R = \underset{\substack{\uparrow \\ \text{numerical value}}}{V(R)} \underset{\substack{\uparrow \\ \text{unit}}}{[R]}$

### Fundamental units

This is a set of units  $F_1, \dots, F_m$  such that any other unit can be written as  $F_1^{a_1} F_2^{a_2} \dots F_m^{a_m}$  for rational numbers  $a_1, \dots, a_m$ , and such that  $F_1^{b_1} \dots F_m^{b_m} \neq 1$  for all  $b_1, \dots, b_m$ , except  $b_1 = b_2 = \dots = b_m = 0$ .

Example SI-system

Dimension	Dimension symbol	SI unit
Mass	M	kilo (kg)
Length	L	meter (m)
Time	T	second (s)
Electric current	I	Ampere (A)
Absolute temperature	Θ	Kelvin (K)
Amount	N	Mole (mol)
Luminous intensity	J	Candela (cd)

} Slide of wikipedia definition

Back to the two basic facts:

\* All relations must be dimensionally correct:

Forbidden:  $3m + 2kg = \dots$

$s = 4.9t^2$  (often seen in textbooks)

$\ln(s)$ , where  $[s] = m$  is not possible  $\infty m = f(x, t)$  mass cannot be expressed

↑ By a relation involving only time and length!

However,  $E = mc^2$  is fine since  $[E] = kgm^2s^{-2}$  in SI-units.

\* All relations are valid independent of choice of fundamental units:

This means that if  $f(R_1, \dots, R_n) = 0$  is dimensionally correct,

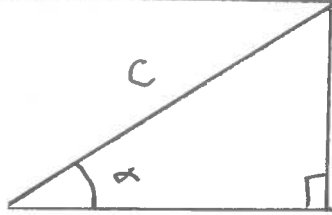
then  $f(v(R_1), v(R_2), \dots, v(R_n)) = 0$  is true

for any choice of fundamental units!

Thus, the function  $f$  must have some type of inherent symmetry! (3)

Dimensional analysis is a systematic procedure for exploiting this symmetry!

Example:



Right triangles

The area of a right triangle is uniquely defined by  $c$  and  $\alpha$ . That is, there is a relation

$$A = f(c, \alpha)$$

giving the area  $A$ . Now if we rescale the triangle with a factor  $a > 0$ , we know the area is rescaled with the factor  $a^2$ . Thus

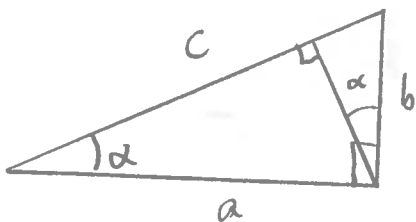
$$f(ac, \alpha) = a^2 f(c, \alpha) \quad (\text{angle } \alpha \text{ does not change})$$

$$\text{set } c=1: f(a, \alpha) = a^2 f(1, \alpha) = a^2 h(\alpha)$$

for some function  $h(\alpha)$ . Thus, just by considering how length and area scales we obtain

$$\underline{A = c^2 h(\alpha)}$$

Corollary: Pythagoras' theorem



From the figure we obtain by adding area

$$a^2 h(\alpha) + b^2 h(\alpha) = c^2 h(\alpha)$$

$$\Rightarrow \underline{a^2 + b^2 = c^2}$$

Let us derive  $A = c^2 h(\alpha)$  in a somewhat different manner. (4)

We know there is some relation  $g(A, c, \alpha) = 0$  (equivalent to  $A = f(c, \alpha)$ ).

Let us make a change of variables

Put  $\pi = \frac{A}{c^2}$  and keep variables  $c$  and  $\alpha$ : Then

$$g(A, c, \alpha) = 0 \Leftrightarrow g(\pi c^2, c, \alpha) = 0 \Leftrightarrow m(\pi, c, \alpha) = 0$$

Now comes the clue: The relation  $m(\pi, c, \alpha) = 0$  must be valid for any choice of length unit, i.e.

$$m(v(\pi), v(c), v(\alpha)) = 0$$

Now  $\pi$  and  $\alpha$  are dimensionless, they have no unit!

Thus  $\pi$  and  $\alpha$  have always the same value independent of length unit, while  $v(c)$  can be any (positive) number.

The function  $m$  is independent of  $c$ ! We obtain

$$m(\pi, \alpha) = 0 \Leftrightarrow \pi = h(\alpha) \Leftrightarrow \underline{A = c^2 h(\alpha)}$$

NB! The last way of deriving  $A = c^2 h(\alpha)$  is how we shall proceed in general in dimensional analysis.

Example The first atomic explosion

It might be plausible that the radius of the fireball of an atomic explosion depends roughly on

1. Time passed,  $t$
2. Total energy of explosion,  $E$
3. The density of the displaced medium (air),  $\rho$

Thus we assume a relation

$$f(r, t, E, \rho) = 0 \quad (*)$$

which must be valid for any choice of units for length, time, and mass. Again, let us try to make a new variable which is dimensionless. Using SI-units we have

$$[r] = m, [t] = s, [E] = \text{kg} \cdot \text{m}^2 \cdot \text{s}^{-2}, [\rho] = \text{kg} \cdot \text{m}^{-3}$$

We observe that  $[\frac{E}{\rho r^5}] = \text{m}^2 \cdot \text{s}^{-2} (\text{m}^{-3})^{-1} = \text{m}^5 \cdot \text{s}^{-2}$

Thus  $\pi = \frac{t^2 E}{\rho r^5}$  is dimensionless, and with this new variable instead of  $E$ , (\*) is equivalent to

$$h(\pi, r, t, \rho) = 0$$

valid for any choice of fundamental units. We realize that  $v(r)$ ,  $v(t)$ , and  $v(\rho)$  can be any numbers, thus

$h$  is independent of  $r, t$ , and  $\rho$  (!). We obtain

$$h(\pi) = 0 \Rightarrow \pi = c, \quad c \text{ constant and } [c] = 1.$$

$$\pi = c \Leftrightarrow \frac{t^2 E}{\rho r^5} = c \Leftrightarrow E = c \rho r^5 / t^2, \quad c \text{ const.}$$

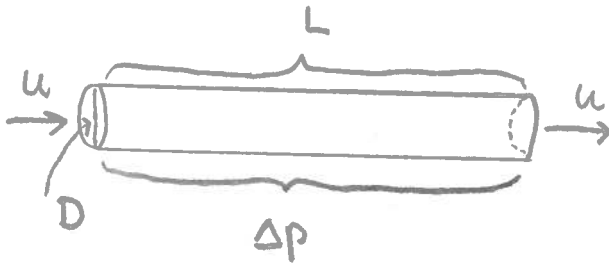
We have  $\rho \approx 1.2 \text{ kg/m}^3$ , and from public pictures we have

$r \approx 110 \text{ m}$  when  $t = 16 \cdot 10^{-3} \text{ s}$ . Thus

$$E \approx c \frac{1.2 \cdot (110)^5}{16^2 \cdot 10^{-6}} \text{ kg} \cdot \text{m}^2 \cdot \text{s}^{-2} = c \cdot 7.5 \cdot 10^{13} \text{ J}$$

Assuming  $c$  is of order 1, we have  $E \approx 10^{14} \text{ J}$

Example: steady state single-phase flow in a pipe



Length of pipe :  $L$ ,  $[L] = m$

Diameter of pipe :  $D$ ,  $[D] = m$

Flow velocity :  $u$ ,  $[u] = ms^{-1}$

Pressure drop :  $\Delta p$ ,  $[\Delta p] = kg\ m^{-1}\ s^{-2}$

Fluid viscosity :  $\mu$ ,  $[\mu] = kg\ m^{-1}\ s^{-1}$

Fluid density :  $\rho$ ,  $[\rho] = kg\ m^{-3}$

Pipe wall roughness :  $E$ ,  $[E] = m$  ( $E$  = "standard deviation of bumps")

We assume there is a relation

$$f(L, D, u, \Delta p, \mu, \rho, E) = 0 \quad (*)$$

1. We start by choosing core variables. The core variables must be dimensionally independent; that is their numerical value can be given any value by changing fundamental units. (more on this topic after the example). So let  $\rho, u, D$  be core variables (other choices also possible), and using the remaining 4 variables let us form 4 dimensionless new variables:

Let  $\pi_1 = \frac{L}{D}$ ,  $\pi_2 = \frac{E}{D}$  be the two first ones.

Now  $[\Delta p] = \text{kg m}^{-1} \text{s}^{-2} \Rightarrow \left[ \frac{\Delta p}{\rho} \right] = \text{m}^2 \text{s}^{-2}$ , so then

$$\underline{\pi_3 = \frac{\Delta p}{\rho u^2}} \text{ is dimensionless, Finally } [\mu] = \text{kg m}^{-1} \text{s}^{-1},$$

so  $\left[ \frac{\mu}{\rho} \right] = \text{m}^2 \text{s}^{-1}$ , thus  $\frac{\mu}{\rho D u}$  is dimensionless. We will

define  $\underline{\pi_4 = \frac{\rho D u}{\mu}}$  as the last dimensionless variable.

Thus (\*) is equivalent to a dimensionally valid relation

$$g(\pi_1, \pi_2, \pi_3, \pi_4, \rho, u, D) = 0,$$

and since  $v(\rho)$ ,  $v(u)$ ,  $v(D)$  can take any value, we must have

$$g(\pi_1, \pi_2, \pi_3, \pi_4) = 0, \text{ or}$$

$$\pi_3 = h(\pi_1, \pi_2, \pi_4).$$

Since we have steady state conditions, the pressure drop must be proportional to the pipe length. Thus

$$\pi_3 = \pi_1 k(\pi_2, \pi_4) \quad (**)$$

for some function  $k$ . Writing out (\*\*) we obtain

$$\Delta p = \frac{\rho u^2 L}{D} k\left(\frac{e}{D}, \frac{\rho u D}{\mu}\right) = \frac{\rho u^2 L}{D} k(e, Re)$$

where  $e$  is the relative roughness, and  $Re$  the famous Reynolds number

$f_F(Re, e) = \frac{1}{2} k(e, Re)$  is the friction factor

and can be determined from further analysis or measurements. That is, we need to cover a 2-dimensional space  $(Re, \epsilon)$  of experiments to determine the relation  $f(L, D, u, \rho, \mu, \rho, E) = 0$  (!)

Moral: Dimensional analysis saves a lot of work for laboratory experimentalists!

see page 13 in compendium for more on  $f$ .

Buckingham's pi-theorem

The last example prepares for the so-called Buckingham's pi-theorem in an excellent manner.

Assume we have a relation (dimensionally valid)

$$f(R_1, R_2, \dots, R_n) = 0 \quad (*)$$

where  $[R_j] = F_1^{a_{1j}} F_2^{a_{2j}} \dots F_m^{a_{mj}}$ ,  $F_1, \dots, F_m$  fundamental units. Then the dimension matrix is  $m \times n$ -matrix  $A$

	$R_1$	$R_2$	$\dots$	$R_n$
$F_1$	$a_{11}$	$a_{12}$	$\dots$	$a_{1n}$
$F_2$	$\vdots$	$\vdots$	$\dots$	$\vdots$
$F_m$	$a_{m1}$	$a_{m2}$	$\dots$	$a_{mn}$

$$A = \{a_{ij}\}_{i,j=1}^{m,n}$$

We can now use well-known ideas from linear algebra to explain dimensional analysis!



In particular, the column-space of  $A$ ,  $\text{col}(A)$  explains (9) "everything".  $\text{rank } A$  is the dimension of  $\text{col}(A)$ .

1.  $\text{rank}(A) =$  number of dimensionally independent variables.

2. Write  $A = [a_{i_1}, a_{i_2}, \dots, a_{i_n}]$ ,  $\text{rank}(A) = k$ , ( $a_i$ ; column vectors). Then  $R_{i_1}, R_{i_2}, \dots, R_{i_k}$  can be chosen as core variables if and only if  $a_{i_1}, \dots, a_{i_k}$  is a basis for  $\text{col}(A)$ .

3. Then any other column vector ( $n-k$  of them) is a linear combination of the basis vectors. This amounts to making  $n-k$  dimensionless variables  $\pi_1, \dots, \pi_{n-k}$  using the  $k$  core variables.

4. Thus we have

$$0 = f(R_1, \dots, R_n) \Leftrightarrow 0 = g(\pi_1, \dots, \pi_{n-k}, R_{i_1}, \dots, R_{i_n}) = 0$$

$\Leftrightarrow g(\pi_1, \dots, \pi_{n-k}) = 0$ , since  $v(R_{i_j})$  can be any value.

Example :

$F_1$	$R_1$	$R_2$	$R_3$	$R_4$	$= A$
$F_2$	1	2	1	1	
$F_3$	-1	0	0	1	
	0	0	1	0	

$\text{rank}(A) = 3$ , and  $a_3$  must be included in a basis (10)

for  $\text{col}(A)$ . The possible bases for  $\text{col}(A)$  are  $\{a_1, a_2, a_3\}$ ,  $\{a_1, a_3, a_4\}$ , and  $\{a_2, a_3, a_4\}$ .

Thus possible core variables are  $\{R_1, R_2, R_3\}$ ,  $\{R_1, R_3, R_4\}$ , and  $\{R_2, R_3, R_4\}$ .

But, there is a catch here!

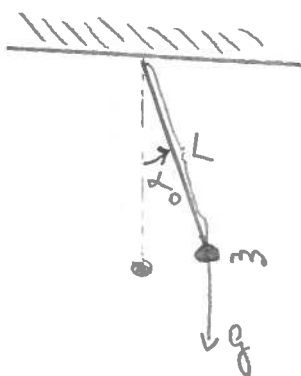
Here we can form  $4 - 3 = 1$  dimensionless variable  $\pi$ , and since  $R_3$  is the only variable containing  $F_3$  in its unit,  $\pi$  cannot be written as an expression containing  $R_3$ .

Thus  $f(R_1, \dots, R_4) = 0 \Leftrightarrow g(\pi) = 0$

means that  $R_3$  cannot be included in

$f(R_1, R_2, R_3, R_4) = 0$  in the first place (i.e. it is impossible to make a dimensionally valid expression  $f(R_1, R_2, R_3, R_4) = 0$  where  $R_3$  is actually included).

### Example Pendulum



We want to find an expression for the frequency  $\omega$  of a swinging pendulum, where we disregard air-resistance.

It is reasonable to assume

$$f(\omega, L, m, \alpha_0, g) = 0 \quad (*)$$

where  $L$  is length of pendulum,  $m$  the mass,  $\alpha_0$  initial angle to equilibrium position, and  $g$  the acceleration of gravity. We have dimension matrix: (11)

$$A = \begin{array}{c} \\ \\ R_g \end{array} \begin{array}{c} m \\ s \\ R_g \end{array} \begin{array}{c} \omega \\ L \\ m \\ \alpha_0 \\ g \end{array} \begin{array}{c} 0 \\ 1 \\ 0 \\ 0 \\ 1 \end{array} \begin{array}{c} -1 \\ 0 \\ 0 \\ 0 \\ -2 \end{array} \begin{array}{c} 0 \\ 0 \\ 1 \\ 0 \\ 0 \end{array}$$

We observe  $\text{rank}(A) = 3$ , however, mass  $m$  is the only variable with  $\text{kg}$  in its unit. Thus, mass cannot be included (\*) if (\*) is dimensionally correct!

Thus, dimension matrix is

$$A = \begin{array}{c} \\ \\ s \end{array} \begin{array}{c} m \\ s \end{array} \begin{array}{c} \omega \\ L \\ \alpha_0 \\ g \end{array} \begin{array}{c} 0 \\ 1 \\ 0 \\ 1 \end{array} \begin{array}{c} -1 \\ 0 \\ 0 \\ -2 \end{array} \quad , \text{rank}(A) = 2.$$

Thus, choose  $L$  and  $g$  as core variables.  $\alpha_0$  is already dimensionless, and we see  $\pi = \frac{\omega}{\sqrt{g/L}}$  is dimensionless

$$\text{Thus } f(\omega, L, \alpha_0, g) = 0 \Leftrightarrow g(\pi, \alpha_0, L, g) = 0$$

$$\Leftrightarrow g(\pi, \alpha_0) = 0 \Leftrightarrow \pi = h(\alpha_0) \Leftrightarrow \underline{\omega = \sqrt{g/L} h(\alpha_0)}$$

Thus, fixing a planet (i.e.  $g$ ) and a length  $L$

we can measure  $\pi = \frac{\omega}{\sqrt{g/L}} = h(\alpha_0)$  for varying values of  $\alpha_0$  to estimate  $h(\alpha_0)$  (again saving work in the lab!)

## Equivalent physics

(12)

$$\text{If } f(R_1, \dots, R_n) = 0 \Leftrightarrow g(\pi_1, \dots, \pi_{n-k}) = 0.$$

Then if two solutions to  $f(R_1, \dots, R_n) = 0$   $\{R_1, \dots, R_n\}$  and  $\{\hat{R}_1, \dots, \hat{R}_n\}$  give rise to the same values for  $\pi_1, \dots, \pi_{n-k}$ , then the two cases give "equivalent physics". This means that if you want to measure, say  $R_1$ , in terms of  $R_2, \dots, R_n$ , you can instead consider a "model setup" with values  $\hat{R}_2, \dots, \hat{R}_n$ , and measure  $\hat{R}_1$ .

For instance, in the exercises you are asked to demonstrate, using dimensional analysis, that ship design is difficult!

The values of the dimensionless numbers  $\pi_1, \dots, \pi_{n-k}$  in general define the physical regime for the problem at hand. E.g. for flow in pipe, the Reynolds number  $Re$  is strongly linked to how turbulent the flow is. Like  $Re < 2000$  means the flow is laminar (no turbulence). For much larger values of  $Re$  the flow is turbulent.