

Exercise 6, solutions.

a) Multiply by a test function w , and use partial integration on the first term:

$$-\mu \int_0^L u_x w_x dx - \int_0^L u_x w + b \int_0^L u_x w dx = \int_0^L f w dx$$

For the integrals to exist, we require $u, w \in H^1(\Omega)$ and $f \in L^2(\Omega)$. Further, the second term is zero if $w(0) = w(L) = 0$, that is $w \in H_0^1(\Omega)$. So we get

Find $u \in H^1(\Omega)$, $u(0) = 0$, $u(L) = L$, s.t.

$$\mu \int_0^L u_x w_x dx + b \int_0^L u_x w dx = \int_0^L f w dx, \quad w \in H_0^1(\Omega).$$

To make the trial and test space the same, let

$$u = \overset{\circ}{u} + R_g, \quad \text{where } R_g(0) = 0, R_g(L) = 1$$

for some chosen lifting function R_g .
The weak formulation becomes

Find $\overset{\circ}{u} \in H_0^1(\Omega)$ s.t. $a(u, w) = F(w) \quad \forall w \in H_0^1(\Omega)$.

with

$$a(u, w) = \mu \int_0^L u_x w_x dx + b \int_0^L u_x w dx$$

$$F(w) = \int_0^L f w dx - \mu \int_0^L R_{g,x} w_x dx - b \int_0^L R_{g,x} w dx.$$

Notice the lack of symmetry in this case, $a(u, w) \neq a(w, u)$.

b) The Galerkin approximation is given by:

Find $\overset{\circ}{u}_h \in V_h$ s.t. $a(\overset{\circ}{u}_h, \varphi_i) = F(\varphi_i), \quad i = 1, \dots, N$.

Let $\overset{\circ}{u}_h = \sum_{j=1}^N \underline{u}_j \varphi_j$. Using the bi-linearity of a , this is

Find $\underline{u} \in \mathbb{R}^N$ s.t. $\sum_{j=1}^N a(\varphi_j, \varphi_i) \underline{u}_j = F(\varphi_i), \quad i = 1, \dots, N$

or simply

$$A_h \underline{u} = \underline{f},$$

Where

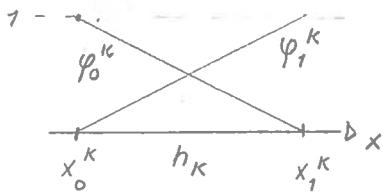
$$(A_h)_{ij} = \mu \int_0^L \frac{\partial \varphi_i}{\partial x} \frac{\partial \varphi_j}{\partial x} dx + b \int_0^L \frac{\partial \varphi_j}{\partial x} \cdot \varphi_i dx$$

$$f_i = \int_0^L f \cdot \varphi_i dx - \mu \int_0^L \frac{\partial R_g}{\partial x} \frac{\partial \varphi_i}{\partial x} dx - b \int_0^L \frac{\partial R_g}{\partial x} \varphi_i dx$$

And the solution is:

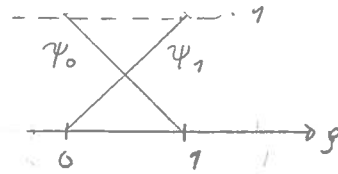
$$u_h = u_h^0 + R_g = \sum_{j=1}^N u_j \varphi_j + R_g$$

c) Element K



$$\varphi_\alpha^K = \psi_\alpha(\xi(x))$$

Reference element \hat{K}



$$\psi_0 = 1 - \xi$$

$$\psi_1 = \xi$$

Mapping $\hat{K} \Leftrightarrow K$: $x(\xi) = x_0^K + h_K \xi$

$$\xi(x) = \frac{x - x_0^K}{h_K}$$

So, element (α, β) of A^K , $\alpha, \beta = 0, 1$ is

$$\begin{aligned} A_{\alpha, \beta}^K &= \int_K a(\varphi_\alpha, \varphi_\beta) dx \\ &= \mu \int_0^L \frac{\partial \varphi_\alpha}{\partial x} \frac{\partial \varphi_\beta}{\partial x} dx + b \int_0^L \frac{\partial \varphi_\beta}{\partial x} \varphi_\alpha dx \\ &= \frac{\mu}{h_K} \int_0^1 \frac{\partial \psi_\alpha}{\partial \xi} \frac{\partial \psi_\beta}{\partial \xi} d\xi + b \int_0^1 \frac{\partial \psi_\beta}{\partial \xi} \psi_\alpha d\xi \end{aligned}$$

$$A_{00}^K = \frac{\mu}{h_K} \int_0^1 1 \cdot 1 d\xi + b \int_0^1 (-1) \cdot (1 - \xi) d\xi = \frac{\mu}{h_K} - \frac{b}{2}$$

$$A_{01}^K = \frac{\mu}{h_K} \int_0^1 1 \cdot (-1) d\xi + b \int_0^1 1 \cdot (1 - \xi) d\xi = -\frac{\mu}{h_K} + \frac{b}{2}$$

$$A_{10}^K = \frac{\mu}{h_K} \int_0^1 (-1) \cdot 1 d\xi + b \int_0^1 (-1) \cdot \xi d\xi = -\frac{\mu}{h_K} - \frac{b}{2}$$

$$A_{11}^K = \frac{\mu}{h_K} \int_0^1 (-1) \cdot (-1) d\xi + b \int_0^1 1 \cdot \xi d\xi = \frac{\mu}{h_K} + \frac{b}{2}$$

So, the element matrix becomes:

$$A^k = \frac{\mu}{h_k} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} + \frac{b}{2} \begin{bmatrix} -1 & 1 \\ -1 & 1 \end{bmatrix}$$

d) First, we need the mapping from the local to the global grid, which in this case simply is

$$i = \theta(\alpha, k) = k-1 + \alpha, \quad k=1, 2, \dots, M-1, \quad \alpha=0, 1$$

such that, for each $x_i = ih$, we have a basis function satisfying

$$\varphi_i(x_j) = \delta_{ij}, \quad \text{and } \varphi_i|_k = \varphi_{\theta(\alpha, k)}$$

The extended matrix \tilde{A}_h , which includes the boundary functions, becomes

$$\begin{aligned} \tilde{A}_h &= \frac{\mu}{h} \begin{bmatrix} 1 & -1 & & & & \\ -1 & (1+1) & -1 & & & \\ & -1 & (1+1) & -1 & & \\ & & \ddots & \ddots & \ddots & \\ & & & \ddots & (1+1) & -1 \\ & & & & -1 & 1 \end{bmatrix} \\ &+ \frac{b}{2} \begin{bmatrix} -1 & 1 & & & & \\ -1 & (1-1) & 1 & & & \\ & -1 & (1-1) & 1 & & \\ & & \ddots & \ddots & \ddots & \\ & & & \ddots & (1-1) & -1 \\ & & & & -1 & 1 \end{bmatrix} \\ &= \frac{\mu}{h} \begin{bmatrix} 1 & 1 & & & & 0 \\ 1 & 2 & & & & \\ & \ddots & \ddots & \ddots & & \\ 0 & & & 2 & 1 & \\ & & & 1 & 1 & \end{bmatrix} + \frac{b}{2} \begin{bmatrix} -1 & 1 & & & & 0 \\ -1 & 0 & & & & \\ & 0 & \ddots & \ddots & \ddots & \\ 0 & & & \ddots & 0 & 1 \\ & & & & -1 & 1 \end{bmatrix} \end{aligned}$$

In the derivation of the weak formulation, we could choose the lifting function $R_g \in H^1(\Omega)$, as long as it satisfies the boundary conditions. Now we choose

$$R_g = u(0)\varphi_0 + u(L)\varphi_{M-1}, \text{ such that}$$

$$\text{and } u_h = \sum_{i=0}^{M-1} u_i \varphi_i, \text{ with } u_0 = u(0) \text{ and } u_{M-1} = u(L),$$

The numerical solution satisfies

$$\tilde{A} \tilde{u} = \tilde{f}, \quad \tilde{u} = \begin{bmatrix} u_0 \\ u \\ u_{M-1} \end{bmatrix}$$

and $f_i = \int_0^L f \varphi_i dx, i = 0, \dots, M.$

Since $u_0 = 0$ and $u_M = 1$, we can remove the first and last row of the system, and move the first and last column to the rhs. Thus the global system is:

$$(\mu A_d + b A_t) \underline{u} = \underline{f} + \underline{b}$$

where

$$A_d = \text{tridiag} \{-1, 2, -1\} \cdot 1/h$$

$$A_t = \text{tridiag} \{-1, 0, 1\} \cdot 1/2$$

$$\underline{b} = (0, \dots, 0, \mu/h - b/2)^T$$

$$\underline{f} = (f_1, \dots, f_{M-1})^T.$$

e) Our scheme becomes, using $f = 0$, becomes

$$-\frac{\mu}{h} (-u_{i-1} + 2u_i - u_{i+1}) + \frac{b}{2} (-u_{i-1} + u_{i+1}) = 0$$

Try the solutions $u_i = r^i$, inserting this into the formula above gives

$$\left(-\frac{\mu}{h} + \frac{b}{2}\right) r^2 - \frac{2\mu}{h} r + \left(-\frac{\mu}{h} - \frac{b}{2}\right) = 0$$

Which has solutions

$$r_1 = 1, \quad r_2 = \frac{1 + \frac{bh}{2\mu}}{1 - \frac{bh}{2\mu}}$$

and the total solution is

$$u_i = C_1 r_1^i + C_2 r_2^i$$

the constants C_1 and C_2 are given from the boundary conditions. It is clear that oscillations will take place if $r_2 < 0$, that is, when

$$\left| \frac{bh}{2\mu} \right| > 1 \quad \text{or} \quad h > \frac{2\mu}{|b|}$$

Which in our case means

$$h > \frac{2 \cdot 0.04}{2} = 0.04, \quad \text{or} \quad M < \frac{L}{0.04} = 500.$$

f) Use artificial diffusion; solve the modified problem

$$-\mu(1 + \varphi_h) u_{xx} + b u_x = f$$

s.t.

$$\frac{|b|h}{2\mu(1 + \varphi_h)} < 1, \quad \text{f.ex. } \varphi_h = \frac{|b|h}{2\mu}$$

Which will lead to the standard upwind scheme.

b) You are not really expected to solve this unless you participated at the lecture 17.11. In this note, I will use an alternative approach.

Given

$$-\mu u_{xx} + b u_x = f$$

which can be rewritten as

$$-\mu \psi_x = f, \quad \text{with} \quad \psi = u_x - \beta u, \quad \beta = b/\mu.$$

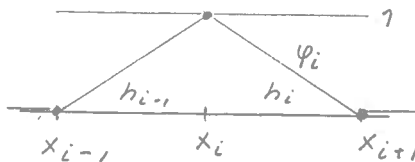
with some appropriate boundary conditions, (Dirichlet)
The weak formulation for this is

$$-\mu \int_{\Omega} \psi_x v \, dx = \int_{\Omega} \psi v_x \, dx = \int_{\Omega} f \cdot v \, dx, \quad \forall v \in H_0^1(\Omega)$$

and the Galerkin approximation becomes

$$(*) \quad \mu \int_{\Omega} \psi_h \frac{\partial \varphi_i}{\partial x} \, dx = \int_{\Omega} f \cdot \varphi_i \, dx$$

let $V_h = X_0^1(\Omega)$, or



Then, on each interval (x_i, x_{i+1}) , ψ_i is linear, and on each interval,

Next, approximate $\psi_h = \psi_{i+1/2}$ (const.) on (x_i, x_{i+1}) and solve

$$\mu \psi_x - \beta \psi = \psi_{i+1/2}, \quad \psi(x_i) = u_i, \quad \psi(x_{i+1}) = u_{i+1}$$

which gives

$$\psi_{i+1/2} = \beta \frac{e^{\beta h} u_i - u_{i+1}}{e^{\beta h} - 1}$$

Inserted into (*) this becomes

$$\frac{\mu \beta}{h} \cdot \frac{1}{e^{\beta h} - 1} (e^{\beta h} u_{i-1} - u_i - e^{\beta h} u_i + u_{i+1}) = \int_{x_{i-1}}^{x_{i+1}} f \cdot \varphi_i \, dx$$

for $i = 1, 2, \dots, M-1$.

Since u_0 and u_M are known, this is diagonal dominant system of equations, with a unique solution.

And it can be rewritten into the form

$$\frac{\mu}{h} (1 + \Phi_h) (-u_{i-1} + 2u_i - u_{i+1}) + \frac{1}{2} (-u_{i-1} + u_{i+1}) = \int f \cdot \varphi_i \, dx$$

with

$$\Phi_h = -1 + \frac{\beta h}{2} + \frac{\beta h}{e^{\beta h} - 1}.$$

It is pretty clear that this is the exact solution for $f=0$.

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i) If $h_k \neq h_{k+1}$, then a typical row is

$$\mu \left(\frac{1}{h_k} (-u_{k-1} + u_k) + \frac{1}{h_{k+1}} (u_k - u_{k+1}) \right) + \frac{b}{2} \left(\frac{1}{h_k} (-u_{k-1} + u_k) + \frac{1}{h_{k+1}} (-u_k + u_{k+1}) \right) = \int_{x_{k-1}}^{x_{k+1}} f \varphi_k dx$$