

On the stability of the Magnus method

Jitse Niesen

Ustaoset, 15 February 2005

Introduction

Anto, Hans et al. wrote in *Acta Numerica* 2000:

[M]uch needs to be done in the realm of stability investigations [. . .] Interesting results abound which cannot yet be fitted into a general theory.

Introduction

Anto, Hans et al. wrote in *Acta Numerica* 2000:

[M]uch needs to be done in the realm of stability investigations [...] Interesting results abound which cannot yet be fitted into a general theory.

Outline of the talk:

- ▶ An example of a stiff equation
- ▶ Classical stability for RK-methods
- ▶ The behaviour of the Magnus method

The equation

We consider the following linear ODE:

$$y' = \begin{bmatrix} & -\mu & 1 \\ \mu^2 & -\operatorname{sech}^2 t & \\ & & -\mu \end{bmatrix} y, \quad y(-10) = \begin{bmatrix} 1 \\ \mu \end{bmatrix},$$

where $\mu \gg 0$.

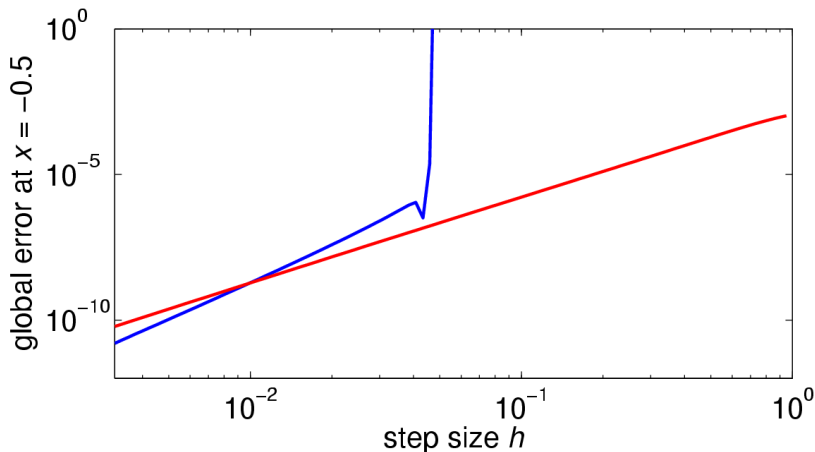
The eigenvalues of the matrix on the right-hand side are:

$$\lambda_1 = -\mu + \sqrt{\mu^2 - \operatorname{sech}^2 x} \approx 0 \quad \text{with } v_1 = \begin{bmatrix} 1 \\ \sqrt{\quad} \end{bmatrix} \approx \begin{bmatrix} 1 \\ \mu \end{bmatrix},$$
$$\lambda_2 = -\mu - \sqrt{\mu^2 - \operatorname{sech}^2 x} \approx -2\mu \quad \text{with } v_2 = \begin{bmatrix} 1 \\ -\sqrt{\quad} \end{bmatrix} \approx \begin{bmatrix} 1 \\ -\mu \end{bmatrix}.$$

So there is one neutral and one very negative eigenvalue, and the solution will approximately follow v_1 .

Global error versus step size

We take $\mu = 30$ and solve the equation on $[-10, -0.5]$.



Blue line = RK4: four stages, fourth order, explicit.

Red line = Radau IIA: two stages, third order, implicit.

Outline of the talk

- ▶ An example of a stiff equation
- ▶ Classical stability for RK-methods
- ▶ The behaviour of the Magnus method

A-stability

We restrict ourselves to **one-step** methods $y_{n+1} = \psi_h(y_n)$.

A-stability theory is based on **Dahlquist's test equation** $y' = \lambda y$.

$$|y(t)| \leq |y(0)| \quad \text{for } t \geq 0 \quad \text{if } \lambda \in \mathbb{C}^- = \{z \in \mathbb{C} : \operatorname{Re} z \leq 0\}.$$

Runge–Kutta methods take the form $\phi_h(y) = R(h\lambda)y$.

The **stability domain** is $S = \{z \in \mathbb{C} : |R(z)| \leq 1\}$.

A method is **A-stable** if the stability domain contains \mathbb{C}^- .

For A-stable methods, we have: $|y_n| \leq |y_0|$ if $\lambda \in \mathbb{C}^-$.

Explicit Runge–Kutta methods are never A-stable.

All Gauss, Radau and Lobatto methods are A-stable.

B-stability

Suppose that the equation $y' = f(t, y)$ satisfies

$$\operatorname{Re}\langle f(t, y) - f(t, z), y - z \rangle \leq 0. \quad (*)$$

Then the equation is **contractive**: any two solutions y and z satisfy

$$|y(t) - z(t)| \leq |y(0) - z(0)| \quad \text{for } t \geq 0.$$

A method ψ_h is said to be **B-stable** if $(*)$ implies that

$$|\psi_h(y) - \psi_h(z)| \leq |y - z|.$$

All Gauss, Radau and Lobatto III C methods are B-stable, but the Lobatto III A and III B methods are not B-stable.

AN-stability

Recall condition (*): $\operatorname{Re}\langle f(t, y) - f(t, z), y - z \rangle \leq 0$.

The linear equation $y' = \lambda(t)y$ satisfies (*) if $\operatorname{Re} \lambda(t) \leq 0$.

We say the method is **AN-stable** if $|\psi_h(y) - \psi_h(z)| \leq |y - z|$ for these equations.

AN-stability

Recall condition (*): $\operatorname{Re}\langle f(t, y) - f(t, z), y - z \rangle \leq 0$.

The linear equation $y' = \lambda(t)y$ satisfies (*) if $\operatorname{Re} \lambda(t) \leq 0$.

We say the method is **AN-stable** if $|\psi_h(y) - \psi_h(z)| \leq |y - z|$ for these equations.

$$\begin{array}{ccc} \text{B-stability} & \implies & \text{AN-stability} & \implies & \text{A-stability.} \\ \left[\begin{array}{l} y' = f(t, y) \\ \text{condition (*)} \end{array} \right] & & \left[\begin{array}{l} y' = \lambda(t)y \\ \operatorname{Re} \lambda(t) \leq 0 \end{array} \right] & & \left[\begin{array}{l} y' = \lambda y \\ \operatorname{Re} \lambda \leq 0 \end{array} \right] \end{array}$$

For **non-confluent** RK-methods: AN-stability \iff A-stability.

B-convergence

Suppose that f satisfies a one-sided Lipschitz condition:

$$\operatorname{Re}\langle f(t, y) - f(t, z), y - z \rangle \leq \nu |y - z|^2.$$

Then any two solutions y and z satisfy

$$|y(t) - z(t)| \leq |y(0) - z(0)| e^{\nu t} \quad \text{for } t \geq 0.$$

The method is **B-convergent** of order r if

$$|y_n - y(t_n)| \leq h^r \gamma(t_n, \nu) \max_{j=1, \dots, \ell} \max_{t \in [0, t_n]} |y^{(j)}(t)| \quad \text{for } h\nu < \alpha.$$

The global error can be bounded in terms of ν and derivatives of the exact solution; the stiffness does not enter.

B-stability is a necessary condition for B-convergence.

Outline of the talk

- ▶ An example of a stiff equation
- ▶ Classical stability for RK-methods
- ▶ The behaviour of the Magnus method

Stability and the Magnus methods

Potential **problem**:

The Magnus series for $y' = A(t)y$ is only guaranteed to converge if $\int_t^{t+h} \|A(\tau)\| d\tau < \pi$; this might limit the step size h .

Results are obtained for the following equations:

- ▶ oscillatory equations (some eigenvalues on imaginary axis)
Magnus performs quite well (see work of **Per Christian, Simon et al.**, ...)
- ▶ Schrödinger equation (∞ eigenvalues near imaginary axis)
Hochbruck & Lubich consider $y' = -i(I - \Delta + V(t))y$.
They find that the above **problem** does not arise.
- ▶ parabolic PDEs (∞ eigenvalues near negative real axis)
... as explained by **Mechthild et al.**

Stability and the Magnus methods

$$\begin{array}{ccc} \text{B-stability} & \implies & \text{AN-stability} & \implies & \text{A-stability.} \\ \left[\begin{array}{l} y' = f(t, y) \\ \text{contractive} \end{array} \right] & & \left[\begin{array}{l} y' = \lambda(t) y \\ \text{Re } \lambda(t) \leq 0 \end{array} \right] & & \left[\begin{array}{l} y' = \lambda y \\ \text{Re } \lambda \leq 0 \end{array} \right] \end{array}$$

Magnus methods are **A-stable** and **AN-stable**, because they solve $y' = \lambda(t) y$ exactly.

A linear equation $y' = A(t) y$ is **contractive** if $\mu_2(A) \leq 0$, where

$$\mu_2(A) = \lim_{\varepsilon \downarrow 0} \frac{\|I + \varepsilon A\| - 1}{\varepsilon} \quad (\text{logarithmic norm})$$

Magnus $y_{k+1} = \exp(\Omega) y_k$ is **B-stable** if this implies $\mu_2(\Omega) \leq 0$. It does when $\Omega = \int A(t) dt$ (Magnus-2). Higher orders?

Back to our equation

Consider the fourth-order Magnus method applied to

$$y' = \begin{bmatrix} & -\mu & 1 \\ \mu^2 & -\operatorname{sech}^2 t & \\ & & -\mu \end{bmatrix} y, \quad y(-10) = \begin{bmatrix} 1 \\ \mu \end{bmatrix},$$

where $\mu \gg 0$.

The eigenvalues of the matrix on the right-hand side are:

$$\lambda_1 = -\mu + \sqrt{\mu^2 - \operatorname{sech}^2 x} \approx 0 \quad \text{with } v_1 = \begin{bmatrix} 1 \\ \sqrt{\quad} \end{bmatrix} \approx \begin{bmatrix} 1 \\ \mu \end{bmatrix},$$
$$\lambda_2 = -\mu - \sqrt{\mu^2 - \operatorname{sech}^2 x} \approx -2\mu \quad \text{with } v_2 = \begin{bmatrix} 1 \\ -\sqrt{\quad} \end{bmatrix} \approx \begin{bmatrix} 1 \\ -\mu \end{bmatrix}.$$

So there is one neutral and one very negative eigenvalue, and the solution will approximately follow v_1 .

The exact solution

Define new coordinates by $y = \begin{bmatrix} 1 & 1 \\ \mu & -\mu \end{bmatrix} \bar{y}$, then

$$\bar{y}' = \begin{bmatrix} -\frac{1}{2}\mu^{-1} \operatorname{sech}^2 t & -\frac{1}{2}\mu^{-1} \operatorname{sech}^2 t \\ \frac{1}{2}\mu^{-1} \operatorname{sech}^2 t & -2\mu + \frac{1}{2}\mu^{-1} \operatorname{sech}^2 t \end{bmatrix} \bar{y}, \quad \bar{y}(t_0) = \begin{bmatrix} 1 \\ 0 \end{bmatrix}.$$

Rationale: The matrix on the left-hand side is almost diagonal.

The exact solution is

$$\bar{y}(t) = \begin{bmatrix} 1 - \frac{1}{2}\mu^{-1} \int_{t_0}^t \operatorname{sech}^2 \tau \, d\tau + \mu^{-2} \int \dots \, d\tau \\ \frac{1}{4}\mu^{-2} \operatorname{sech}^2 t \end{bmatrix} + \mathcal{O}(\mu^{-3}).$$

To find this, substitute the ansatz

$$\bar{y}(t) = \bar{y}_0(t) + \mu^{-1} \bar{y}_1(t) + \mu^{-2} \bar{y}_2(t) + \dots$$

in the equation and equate powers of μ .

The local error

We use Magnus-4 with quadrature at the Gauss–Legendre points:

$$y_{k+1} = \exp(\Omega)y_k, \quad \text{with} \quad \Omega = \frac{1}{2}h(A_1 + A_2) - \frac{\sqrt{3}}{12}h^2[A_1, A_2],$$
$$A_{1,2} = A\left(t_k + \left(\frac{1}{2} \pm \frac{1}{6}\sqrt{3}\right)h\right).$$

The **local error** in the limit $h \rightarrow 0$, $\mu h \rightarrow \infty$ is

$$\bar{L}_k = \mu^{-1} \left[\begin{array}{c} h^5(\dots) + \mathcal{O}(h^7, \mu^{-1}) \\ \frac{1}{12}h^2 \operatorname{sech}' t_{k+1/2} + \mathcal{O}(h^4, \mu^{-1}) \end{array} \right].$$

So, we have **order reduction** in the stiff (second) component:

The local error is not h^5 but h^2 .

The above formula was found by diagonalizing Ω and with the help of some coffee and perseverance.

The global error

$$\bar{L}_k = \mu^{-1} \left[\begin{array}{l} h^5(\dots) + \mathcal{O}(h^7, \mu^{-1}) \\ \frac{1}{12} h^2 \operatorname{sech}' t_{k+1/2} + \mathcal{O}(h^4, \mu^{-1}) \end{array} \right].$$

The error in the stiff (second) component does not propagate, so the **global error** is

$$\bar{G}_k = \mu^{-1} \left[\begin{array}{l} h^4(\dots) + \mathcal{O}(h^6, \mu^{-1}) \\ \frac{1}{12} h^2 \operatorname{sech}' t_{k+1/2} + \mathcal{O}(h^4, \mu^{-1}) \end{array} \right].$$

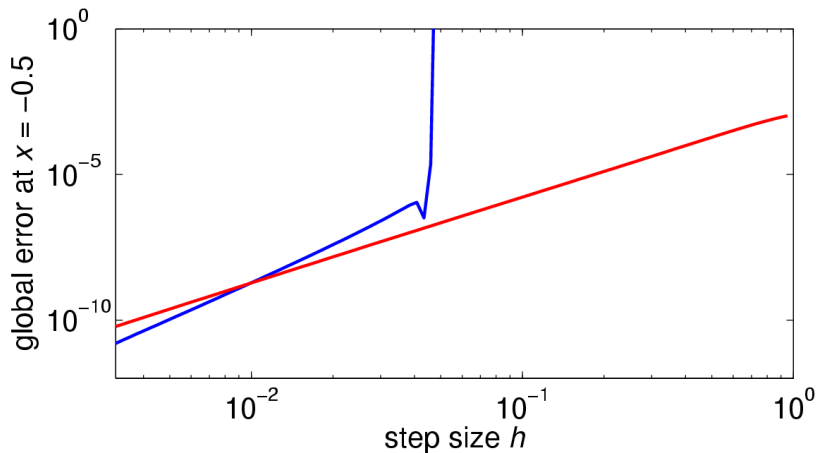
In the original coordinates, this becomes

$$G_k = \left[\begin{array}{l} \frac{1}{12} \mu^{-1} h^2 \operatorname{sech}' t_{k+1/2} + \mathcal{O}(h^4, \mu^{-2}) \\ \frac{1}{12} h^2 \operatorname{sech}' t_{k+1/2} + \mathcal{O}(h^4, \mu^{-1}) \end{array} \right].$$

Hence, the method is **second order** when $\mu \gg h$.

Global error versus step size

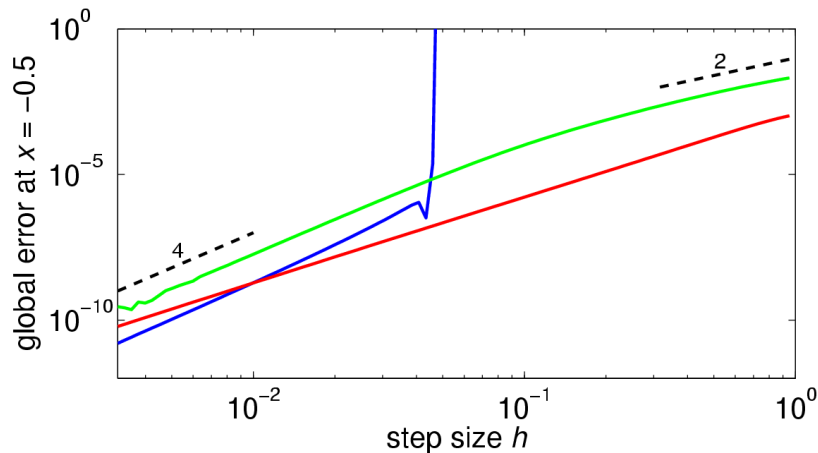
We take $\mu = 30$ and solve the equation on $[-10, -0.5]$.



Blue = RK4; red = 2-stage Radau IIA.

Global error versus step size

We take $\mu = 30$ and solve the equation on $[-10, -0.5]$.



Blue = RK4; red = 2-stage Radau IIA; green = Magnus-4.