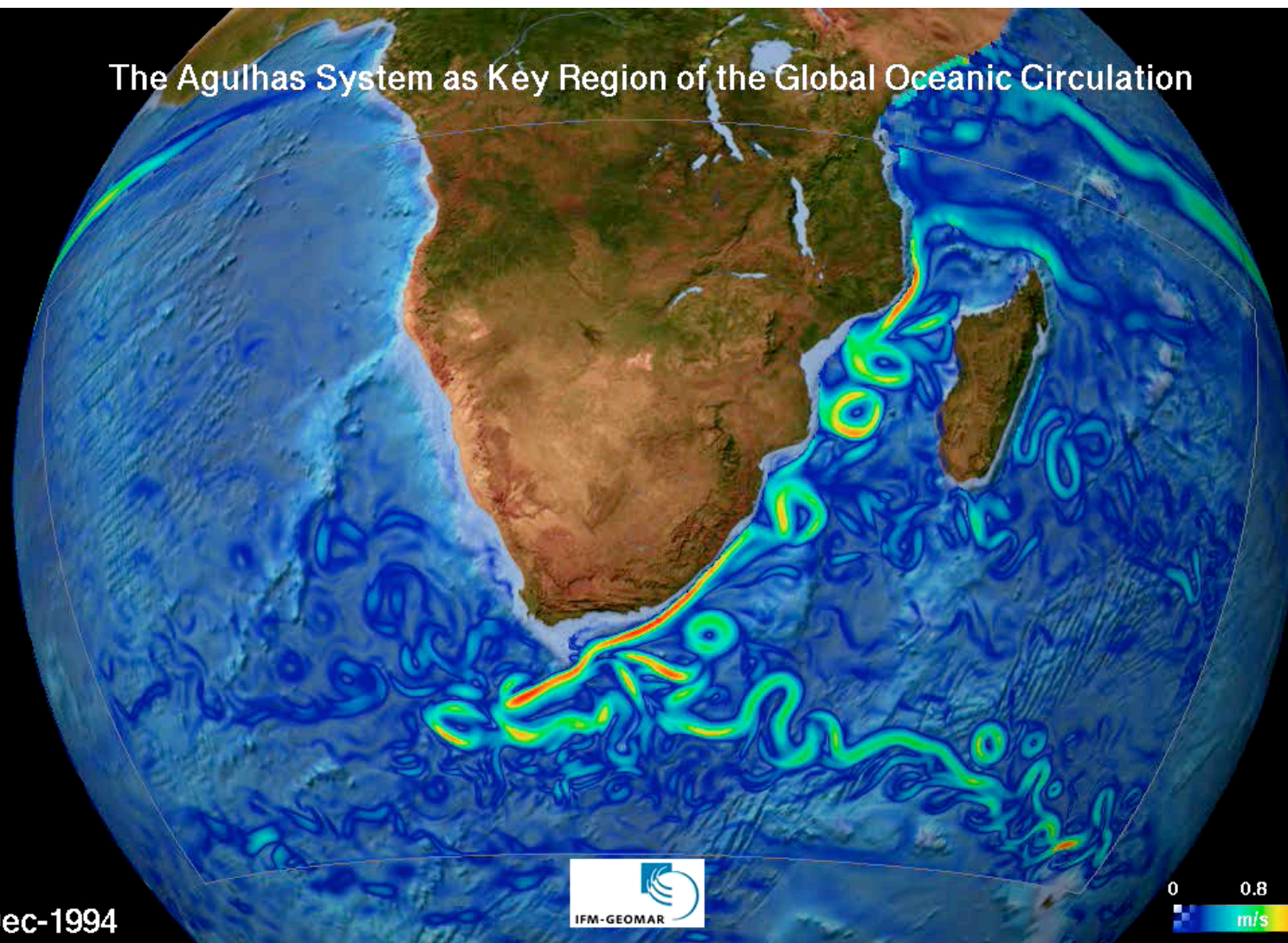


Particle Filters for high-dimensional problems

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University of Reading

Selbu 2012

The Agulhas System as Key Region of the Global Oceanic Circulation



31-Dec-1994



Near-Surface Speeds in a High-Resolution Model, Nested in a Global, Coarse-Resolution Ocean Model Biastoch and Böning, Ocean Modelling Group

How is DA used today in geosciences?

Present-day data-assimilation systems are based on **linearizations** and **state covariances** are essential.

4DVar, Representer method (PSAS):

- Gaussian pdf's for the state, solves only for **posterior mode**, needs error covariance of initial state (B matrix)

(Ensemble) Kalman filter:

- assumes Gaussian pdf's for the state, approximates posterior **mean and covariance, doesn't minimize anything in nonlinear systems**, needs inflation and localisation

Combinations of these: hybrid methods

Nonlinear filtering: Particle filter

$$p(x|y) = \frac{p(y|x)p(x)}{\int p(y|x)p(x) dx}$$

Use ensemble
↓

$$p(x) = \sum_{i=1}^N \frac{1}{N} \delta(x - x_i)$$

$$p(x|y) = \sum_{i=1}^N w_i \delta(x - x_i)$$

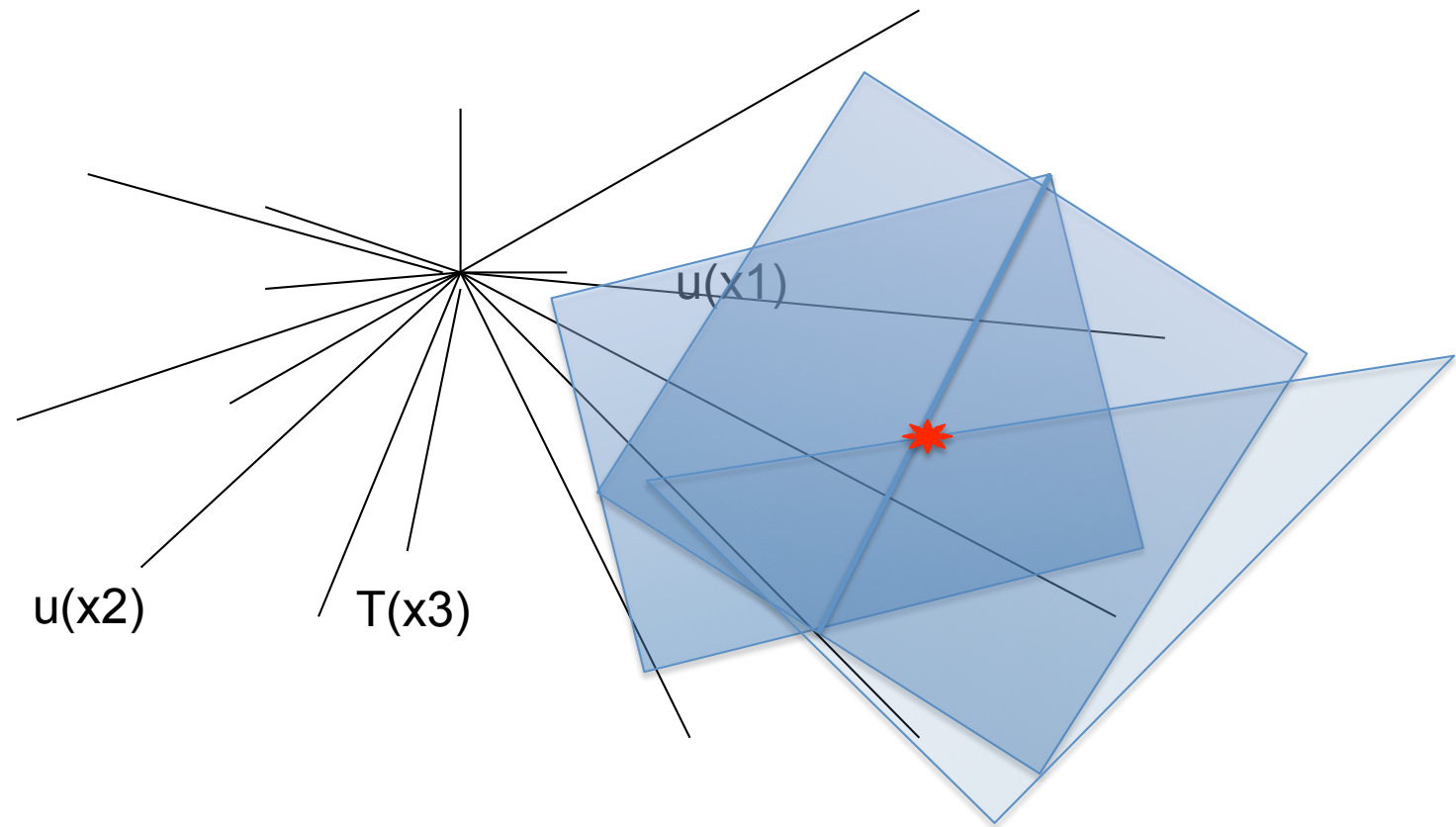
with

$$w_i = \frac{p(y|x_i)}{\sum_j p(y|x_j)}$$

the **weights**.

Why are particle filters degenerate I

Probability space in large-dimensional systems is 'empty': **the curse of dimensionality**



Why are Particle Filters degenerate II

- The volume of a hypersphere of radius r in an M dimensional space is

$$V \propto \frac{r^M}{\Gamma(M/2 - 1)}$$

- Taking for the radius $r \approx 3\sigma_y$ we find, using Stirling:

$$V \propto \left[\frac{9\sigma_y}{M/2} \right]^{M/2}$$

- So very small indeed.

Why are Particle Filters degenerate III

For the optimal proposal density we find, for Gaussian process model and Gaussian observation errors:

$$\begin{aligned} w_i &\propto p(y^n | x_i^{n-1}) \\ &\propto \exp \left[-\frac{1}{2} (y^n - H f(x_i^{n-1})) (H Q H^T + R)^{-1} \right. \\ &\quad \left. \times (y^n - H f(x_i^{n-1})) \right]. \end{aligned}$$

Ignoring covariances we find:

$$\text{var}[-\log(w_i)] \propto \frac{M}{2} \left(\frac{V_x}{V_\beta + V_y} \right)^2 \left(1 + 2 \left(\frac{V_y + V_\beta}{V_x} \right) \right)$$

Why are Particle Filters degenerate?

- ‘Number of particles needed grows exponentially with dimension of the state vector (Bickel et al, 2007).’
- A slightly different view: degeneracy due to number of independent observations.
- This is related to the extremely narrow likelihood, a tiny move of a particle gives a completely different weight.

The statistics

- The Stochastic PDE: $x^n = f(x^{n-1}) + \beta^{n-1}$
- Observations: y^n
- Relation between the two: $y^n = H(x^n) + \epsilon^n$

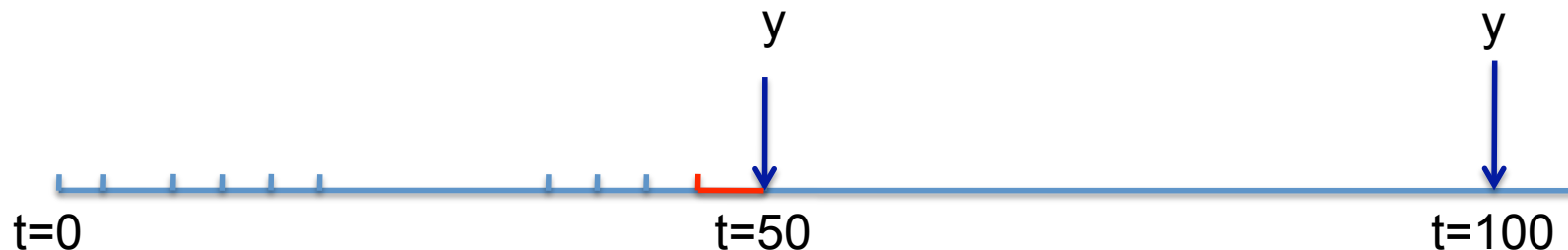
Assume: $\beta \sim N(0, Q)$

$\epsilon \sim N(0, R)$

H is linear

The Equivalent-Weights Particle Filter

- Use simple proposal at each time step, e.g. relaxation to observations.
- Use different proposal **at final time step** to ensure that weights are very similar.



Transition density for x^n

- Stochastic model

$$x^n = f(x^{n-1}) + \beta^{n-1}$$

- With $\beta^{n-1} \sim N(0, Q)$

- Hence transition density

$$p(x^n | x^{n-1}) = N\left(f(x^{n-1}), Q\right)$$

Bayes Theorem and the proposal density

Bayes Theorem can be written as:

$$\begin{aligned} p(x^n | y^n) &= \frac{p(y^n | x^n) p(x^n)}{p(y)} \\ &= \frac{p(y^n | x^n)}{p(y)} \int p(x^n | x^{n-1}) p(x^{n-1}) dx^{n-1} \end{aligned}$$

Multiply and divide this expression by a **proposal transition density** q :

$$p(x^n | y^n) = \frac{p(y^n | x^n)}{p(y)} \int \frac{p(x^n | x^{n-1})}{q(x^n | x^{n-1}, y^n)} q(x^n | x^{n-1}, y^n) p(x^{n-1}) dx^{n-1}$$

Proposal transition density

For each particle at time $n-1$ draw a sample from the proposal transition density q , to find:

$$p(x^n | y^n) = \frac{1}{N} \sum_{i=1}^N \frac{p(y^n | x_i^n)}{p(y)} \frac{p(x_i^n | x_i^{n-1})}{q(x_i^n | x_i^{n-1}, y^n)} \delta(x^n - x_i^n)$$

Which can be rewritten as:

$$p(x^n | y^n) = \sum_{i=1}^N w_i \delta(x^n - x_i^n)$$

with weights

$$w_i = \frac{p(y^n | x_i^n)}{p(y)} \frac{p(x_i^n | x_i^{n-1})}{q(x_i^n | x_i^{n-1}, y^n)}$$

Likelihood weight

Proposal weight

Proposal density between observations

We can explore the fact that the model needs several $O(100)$ time steps between observations, e.g. by using a relaxation term in the proposal:

$$q(x^n | x_i^{n-1}, y^m) = N \left(f(x_i^{n-1}) + S \left(y^m - H(x_i^{n-1}) \right), Q \right)$$

Corresponding to an evolution equation for each particle

$$x_i^n = f(x_i^{n-1}) + \hat{\beta}_i^n + S \left(y^n - H(x_i^{n-1}) \right)$$

Proposal density between observations

- One could also use the ‘optimal proposal density’ between observations.
- This can be implemented as a minimization method for each particle, and is also known as the Implicit Particle Filter.
- This is related to a method called 4DVar in meteorology and oceanography, which explores only the mode of the joint-in-time pdf.

Proposal density at observation time: the essence of the Equivalent-Weights Particle Filter

The proposal density depends on the maximum weight a particle can achieve using a **deterministic** time step. It is defined as:

$$q(x^n | x_i^{n-1}, y^n) = \begin{cases} q_1(x^n | x_i^{n-1}, y^n) & \text{if } w_i^{max} > w^{target} \\ q_2(x^n | x_i^{n-1}, y^n) & \text{if } w_i^{max} < w^{target} \end{cases}$$

The target weight is set by the user, as e.g. the weight that 80% of the particles can achieve.

The maximum weights

1. We know:

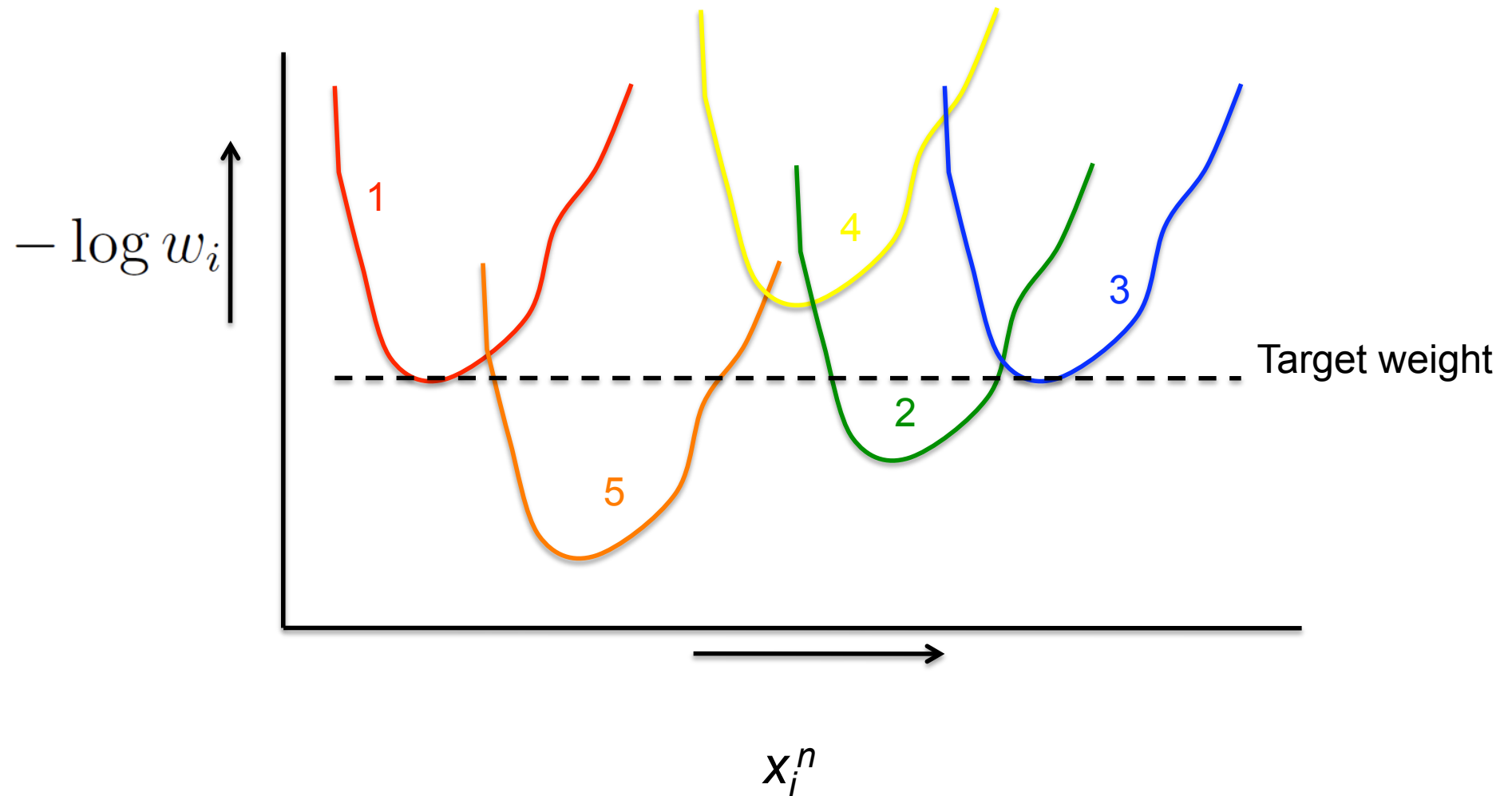
$$w_i = \frac{p(y^n | x_i^n)}{p(y^n)} \frac{p(x_i^n | x_i^{n-1})}{q(x_i^n | x_i^{n-1}, y^n)}$$

2. Write down expression for each weight ignoring proposal:

$$w_i \propto w_i^{rest} \exp \left[-\frac{1}{2} \left(x_i^n - f(x_i^{n-1}) \right)^T Q^{-1} \left(x_i^n - f(x_i^{n-1}) \right) - \frac{1}{2} \left(y^n - H(x_i^n) \right)^T R^{-1} \left(y^n - H(x_i^n) \right) \right]$$

3. When H is linear this is a quadratic function in x_i^n for each particle. Otherwise linearize.

The target weight



The Equivalent-Weights Particle Filter

The proposal density is chosen as:

$$q(x^n | x_i^{n-1}, y^n) = \begin{cases} q_1(x^n | x_i^{n-1}, y^n) & \text{if } w_i^{max} > w^{target} \\ q_2(x^n | x_i^{n-1}, y^n) & \text{if } w_i^{max} < w^{target} \end{cases}$$

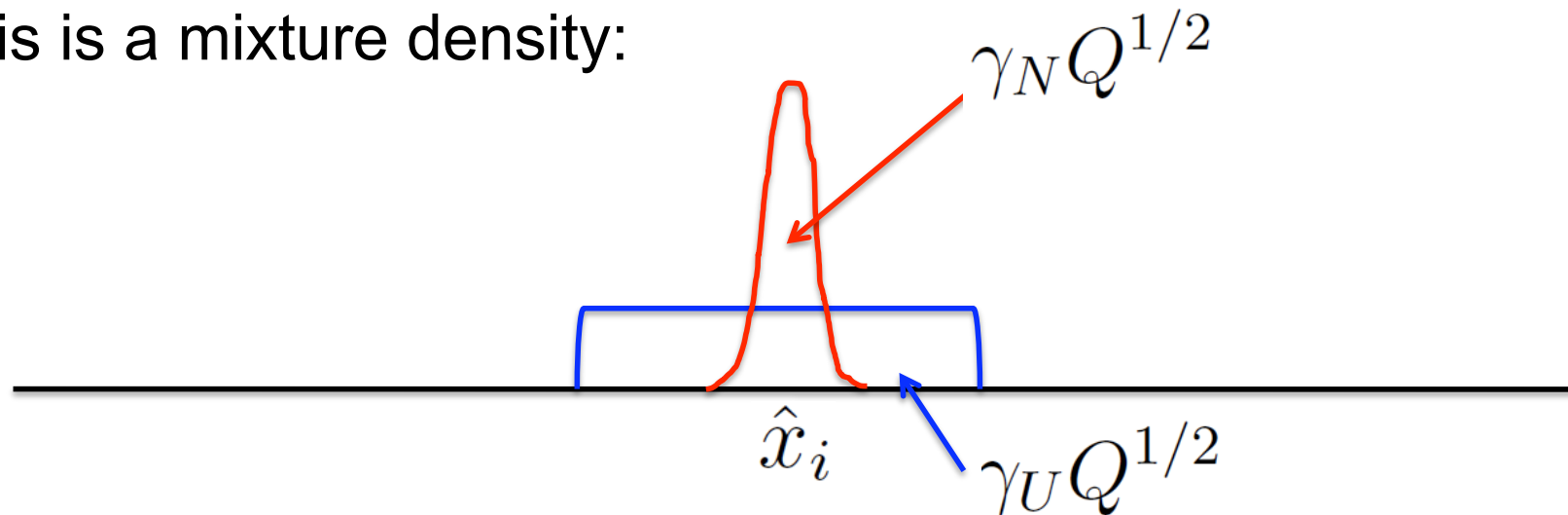
The target weight is set by the user, as e.g. the weight that 80% of the particles can achieve.

The two proposal densities

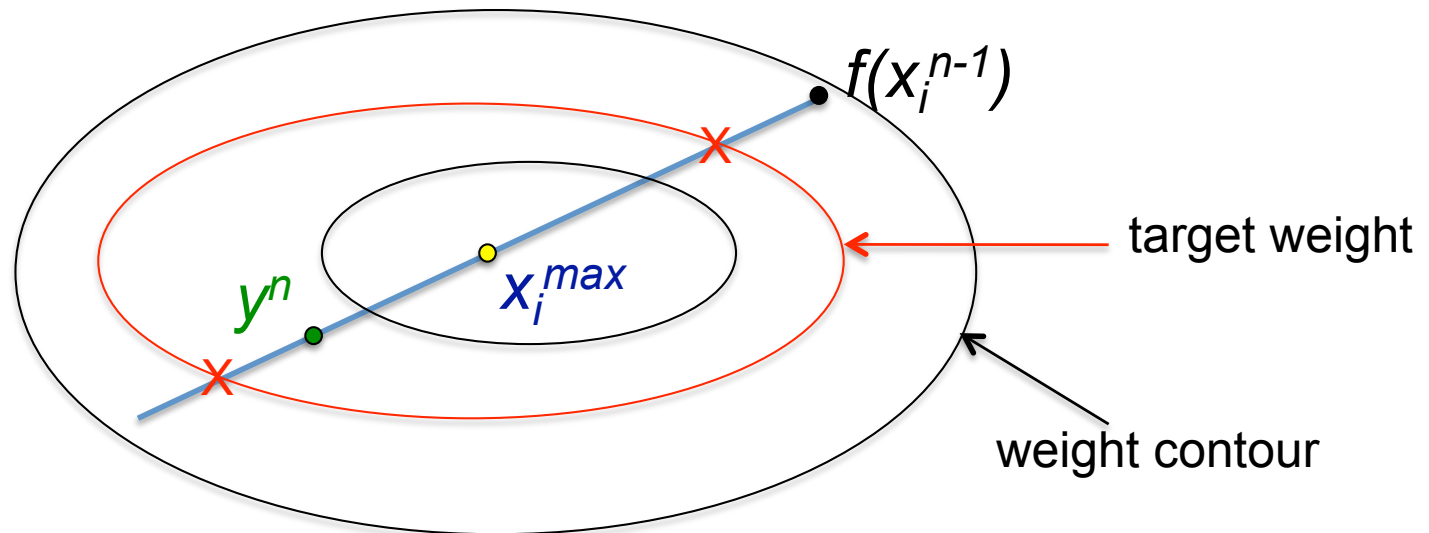
For particles that can reach the target weight we use:

$$q_1(x^n | x_i^{n-1}, y^m) = (1 - \epsilon)U(\hat{x}_i - \gamma_U Q^{1/2} \mathbf{1}, \hat{x}_i + \gamma_U Q^{1/2} \mathbf{1}) + \epsilon N(\hat{x}_i, \gamma_N^2 Q)$$

This is a mixture density:



The deterministic move



Determine α at crossing of line with target weight contour in:

$$\hat{x}_i = f(x_i^{n-1}) + \alpha_i K (y^n - H(f(x_i^{n-1})))$$

with

$$K = QH^T (HQH^T + R)^{-1}$$

The stochastic part of the proposal

A draw from the uniform density gives:

$$w_i = \frac{|Q|^{1/2} (2\gamma_U)^k}{1 - \epsilon} w_i^{rest} p(x_i^n | x_i^{n-1}) p(y^n | x_i^n)$$

A draw from the Gaussian density gives:

$$w_i = \frac{w_i^{rest} p(x_i^n | x_i^{n-1}) p(y^n | x_i^n)}{\frac{\epsilon}{(2\pi)^{k/2} |\gamma_N^2 Q|^{1/2}} \exp(-\frac{1}{2} \gamma_U d\beta_i^n (\gamma_U^2 Q)^{-1} \gamma_U d\beta_i^n)}$$

The ratio between the two is (ignoring the exp part):

$$\frac{(2\pi)^{k/2} |\gamma_N^2 Q|^{1/2}}{\epsilon} \frac{(1 - \epsilon)}{|Q|^{1/2} (2\gamma_U)^k}$$

which can be made equal to one when:

$$\gamma_N = \frac{2^{k/2} \epsilon}{\pi^{k/2} (1 - \epsilon)} \gamma_U^k$$

Equivalent-Weights Particle Filter

- Use relaxation-term proposal up to last time step
- Calculate w_i^{max} and target weight (e.g. 80%)
- Calculate deterministic moves for high-weight particles:

$$\hat{x}_i = f(x_i^{n-1}) + \alpha_i K \left(y^n - H(f(x_i^{n-1})) \right)$$

- Determine stochastic move

$$p(\hat{\beta}_i^{n-1}) \propto (1 - a)U[-b, b] + aN(0, \hat{Q})$$

- Calculate new weights and resample 'lost' particles

How essential are Gaussian assumptions?

- Allows for analytical expressions.
- But no real need.
- w_i^{max} calculations do not have to be very accurate.
- Same for w^{target} .
- Deterministic move has to be **very accurate**, good iterative schemes should be used.

Application: the barotropic vorticity equation

- Stochastic barotropic vorticity equation:

$$\frac{\partial q}{\partial t} + u \cdot \nabla q = F$$

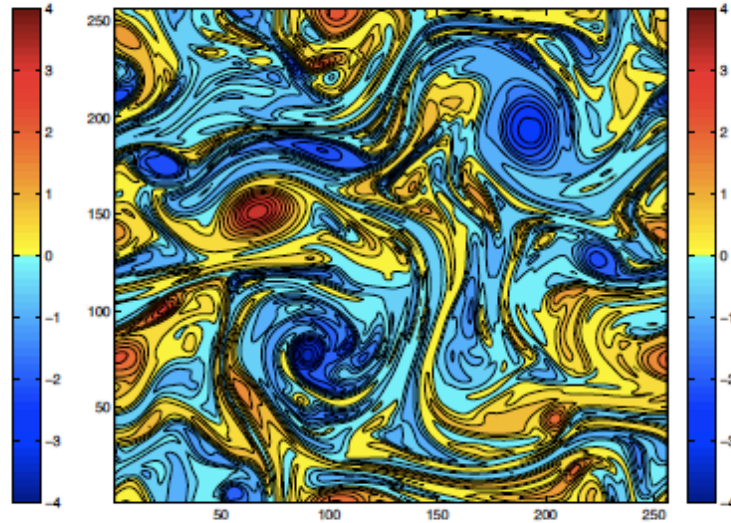
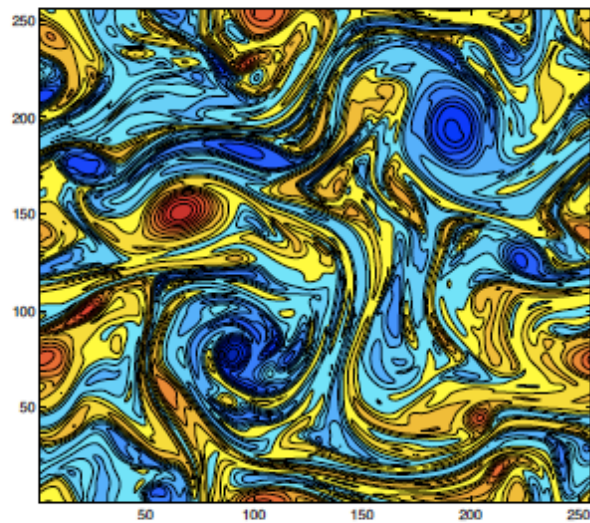
- 256 by 256 grid - 65,536 variables
- Double periodic boundary conditions
- Semi-Lagrangian time stepping scheme
- Twin experiments
- Observations every 50 time steps – decorrelation time of 42
- 32 particles

Fully observed system

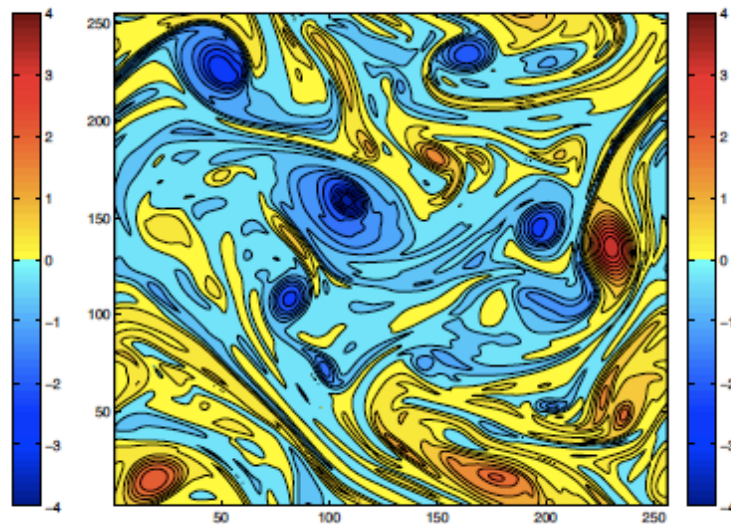
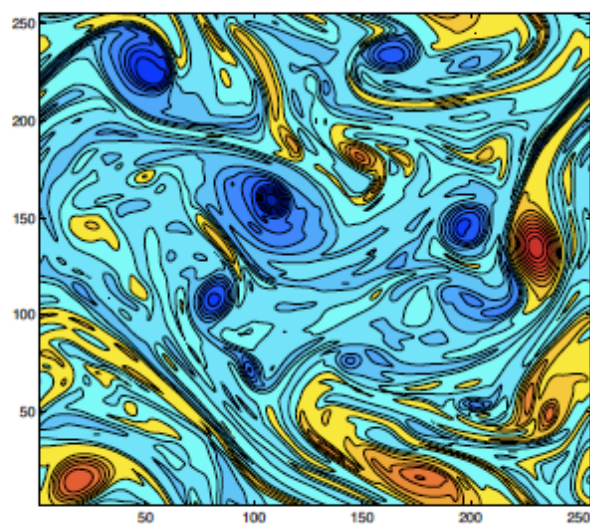
True model state

Mean of particles

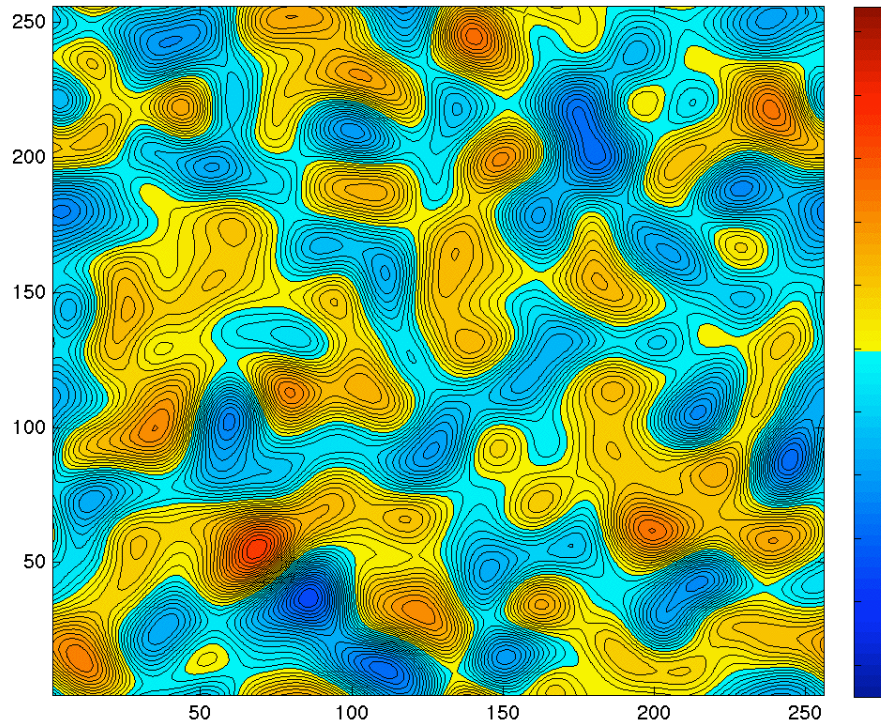
Time step 600



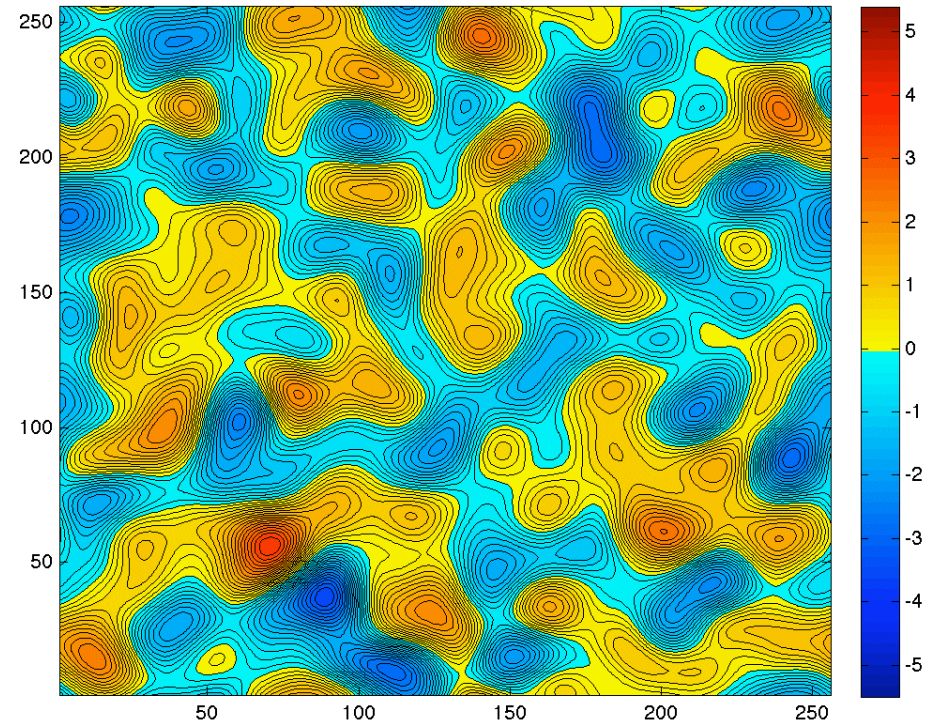
Time step 1150



$\frac{1}{4}$ Observations over half of state

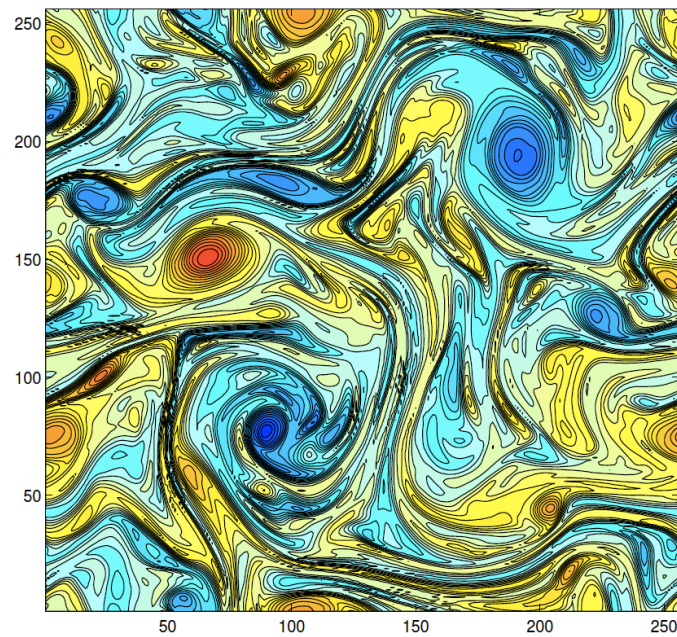


Truth

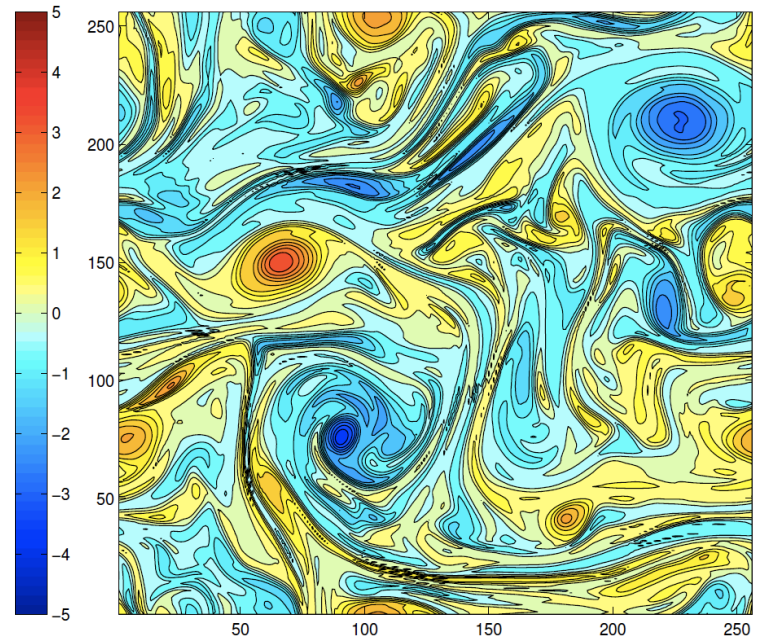


Mean of particle filter ensemble

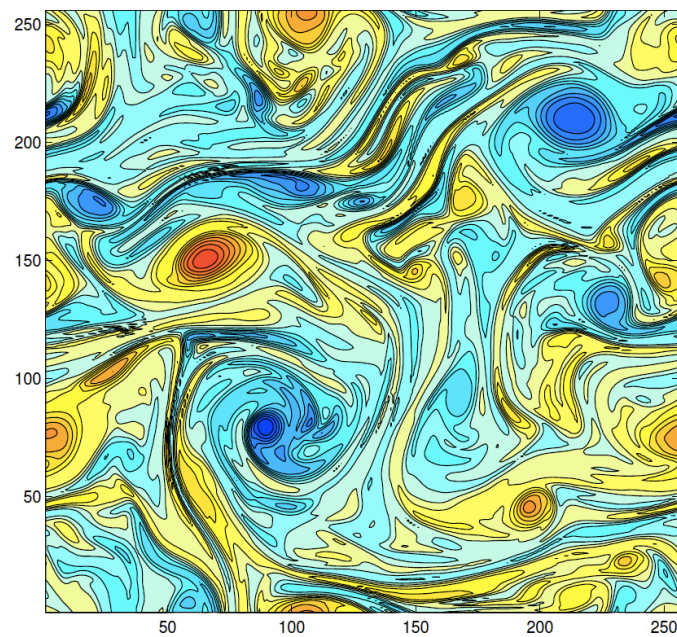
Individual particles are not smooth.



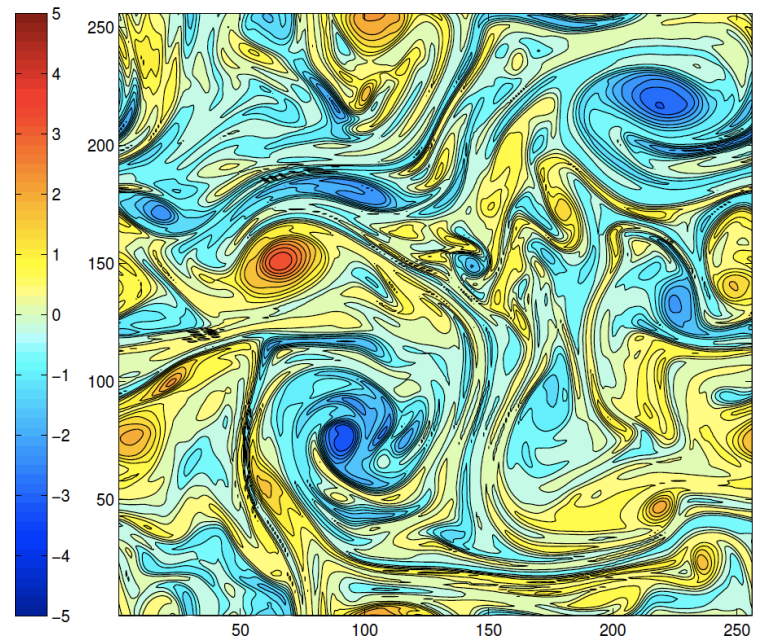
(a) truth



(b) Particle 2

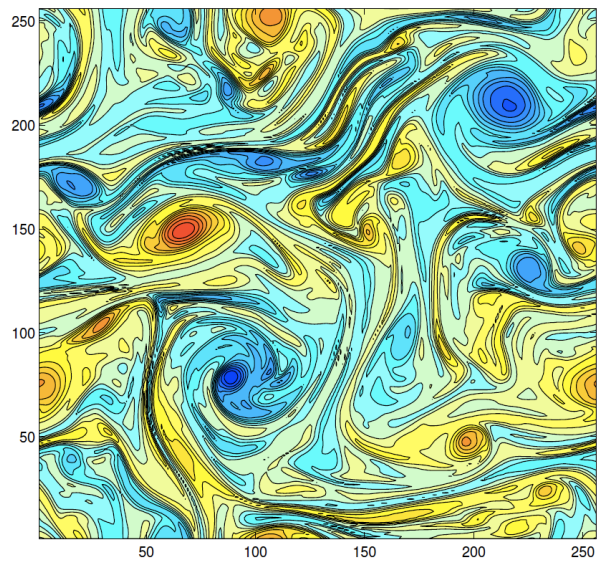


(c) Particle 23

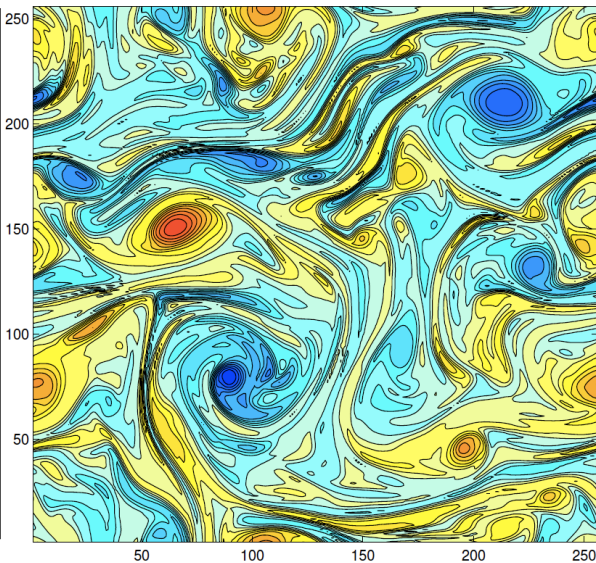


(d) Particle 28

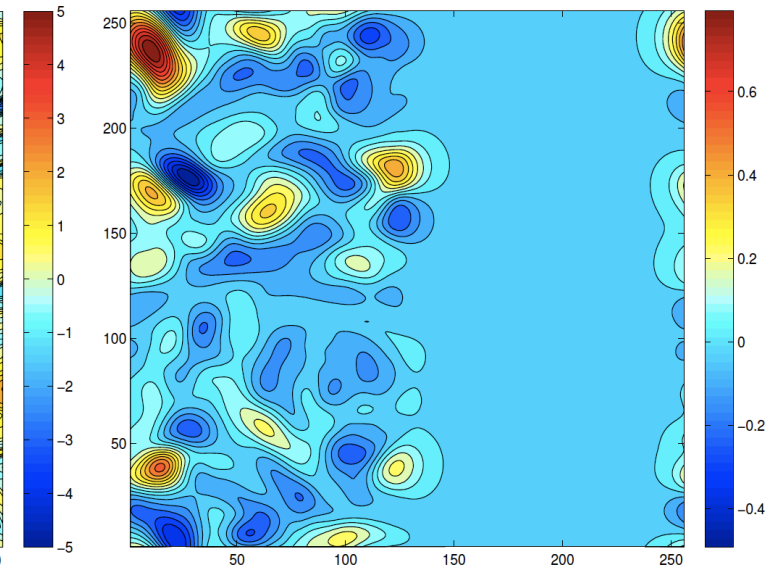
The update of the unobserved part



Particle 23 before update

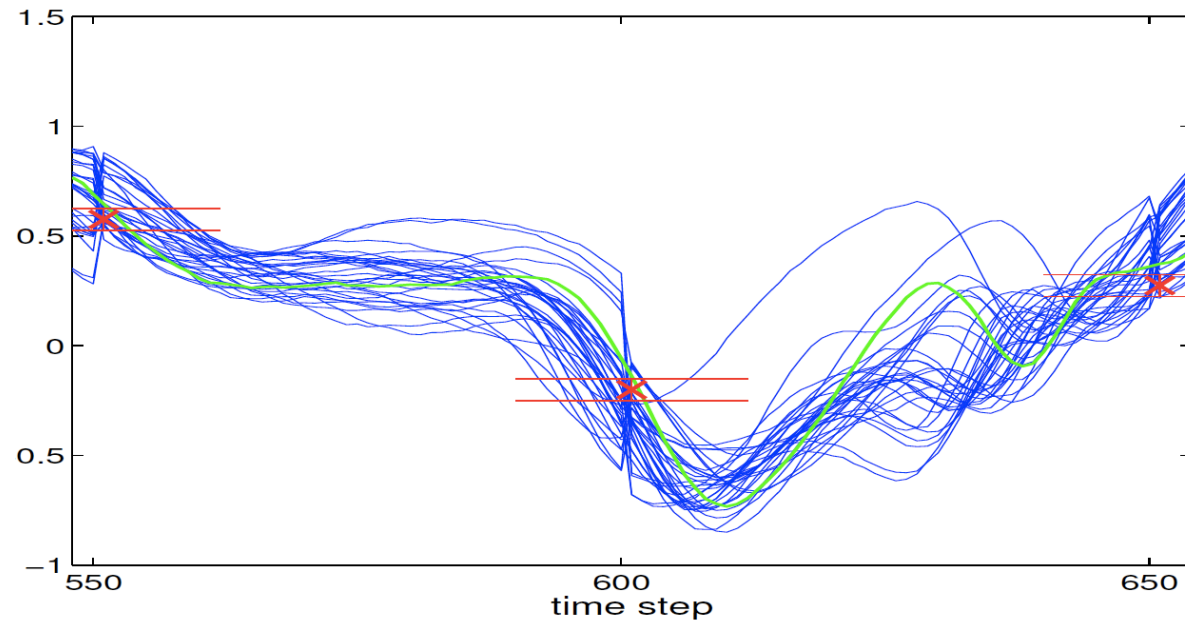


Particle 23 after update



Difference

Time evolution at a specific grid point

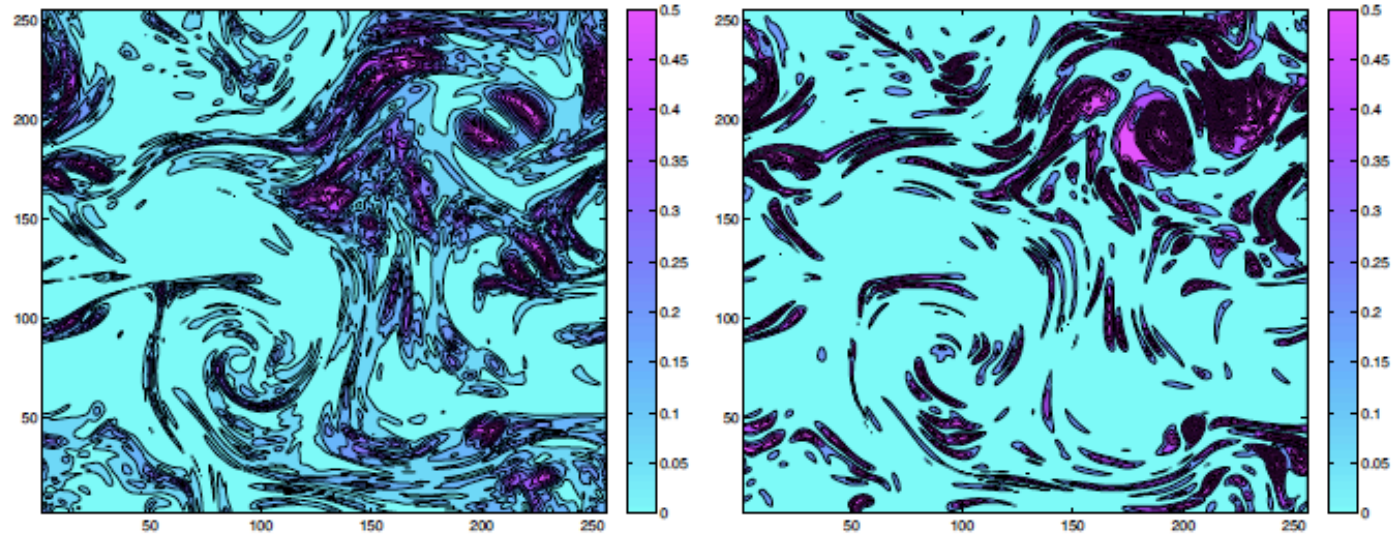


$\frac{1}{4}$ observations over half of state

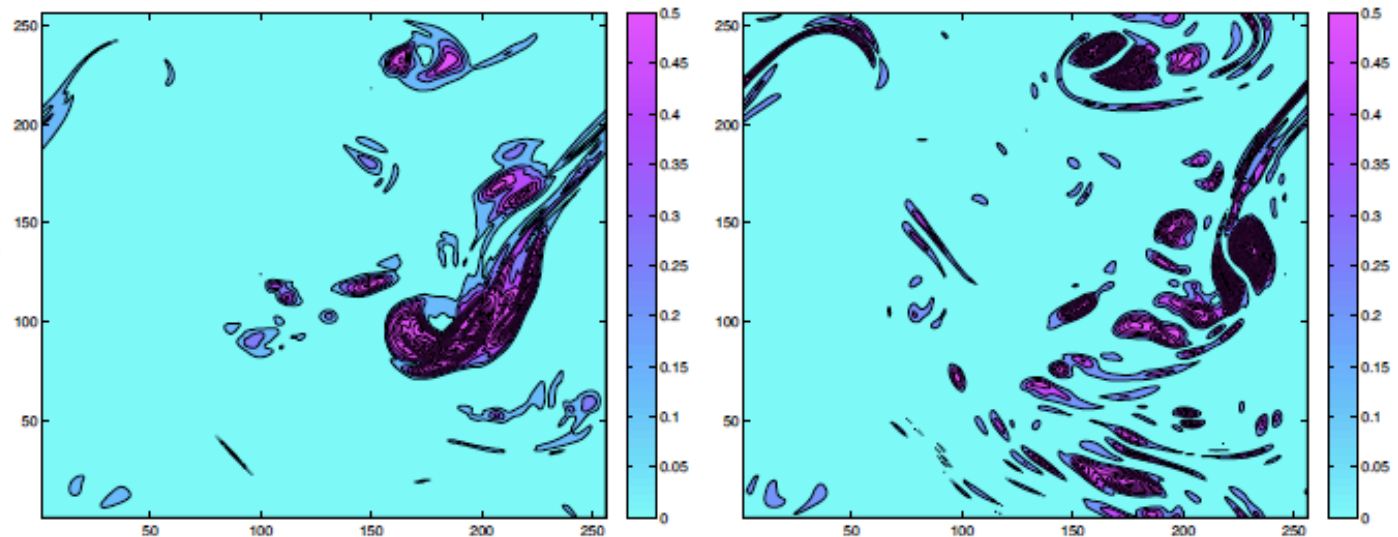
Variance

$(\text{Mean} - \text{Truth})^2$

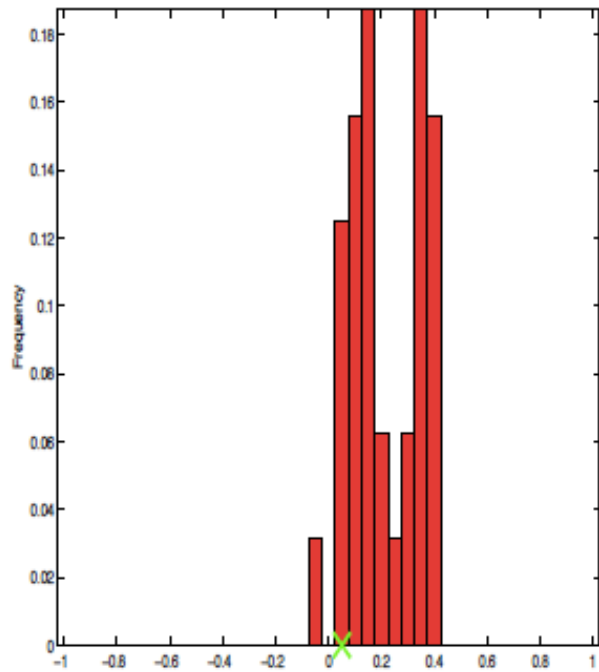
Time step 600



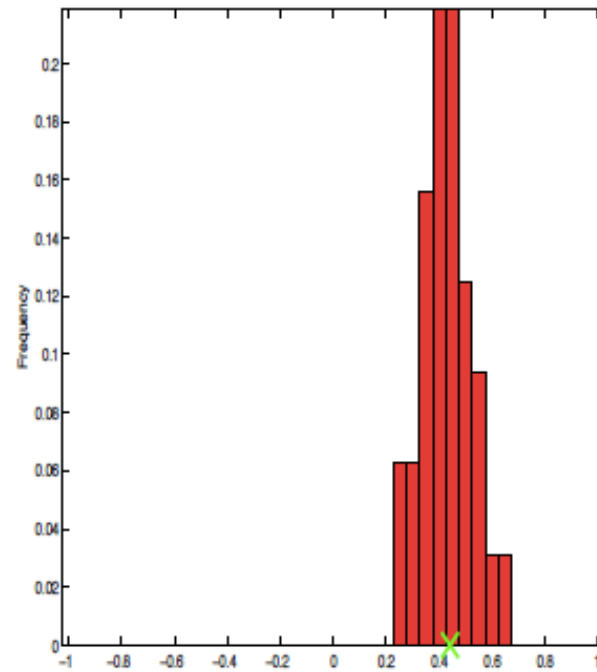
Time step 1150



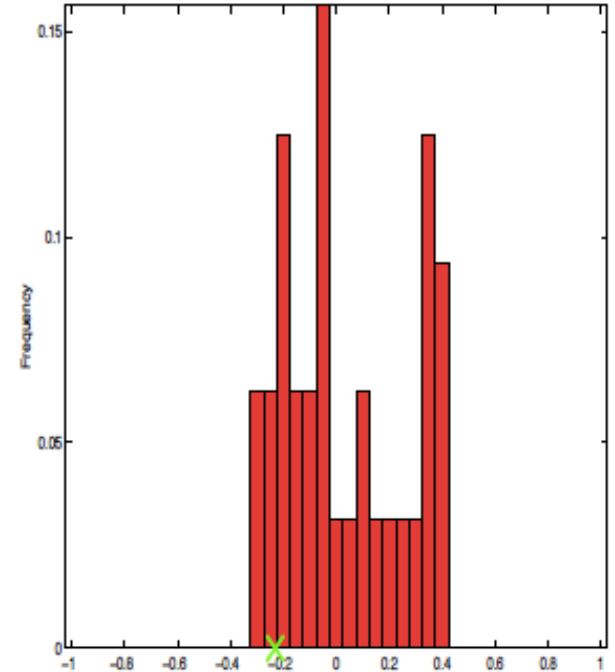
Marginal posterior probability densities



variable (20,64)

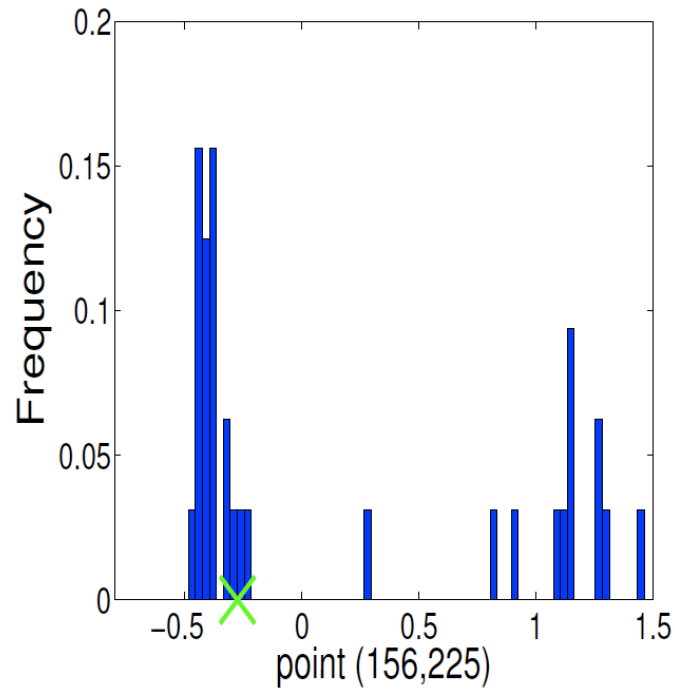


variable (56,64)

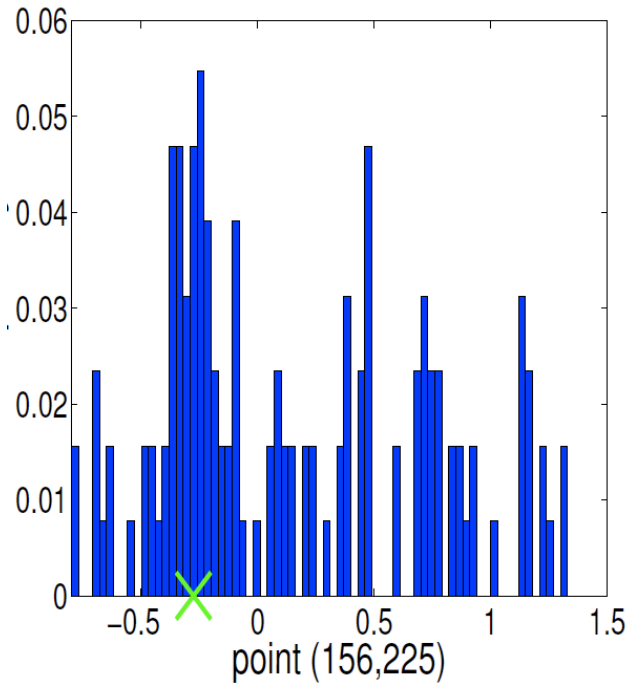


variable (197,64)

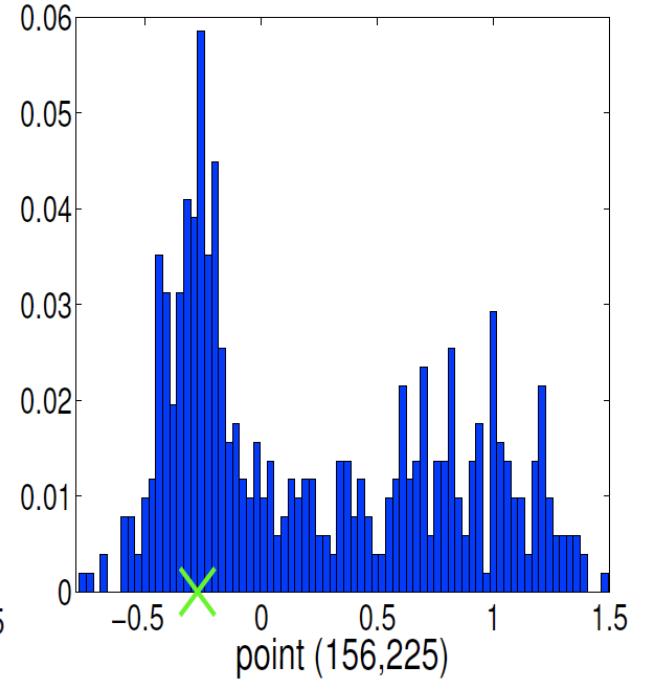
Convergence of the pdf's



32 particles

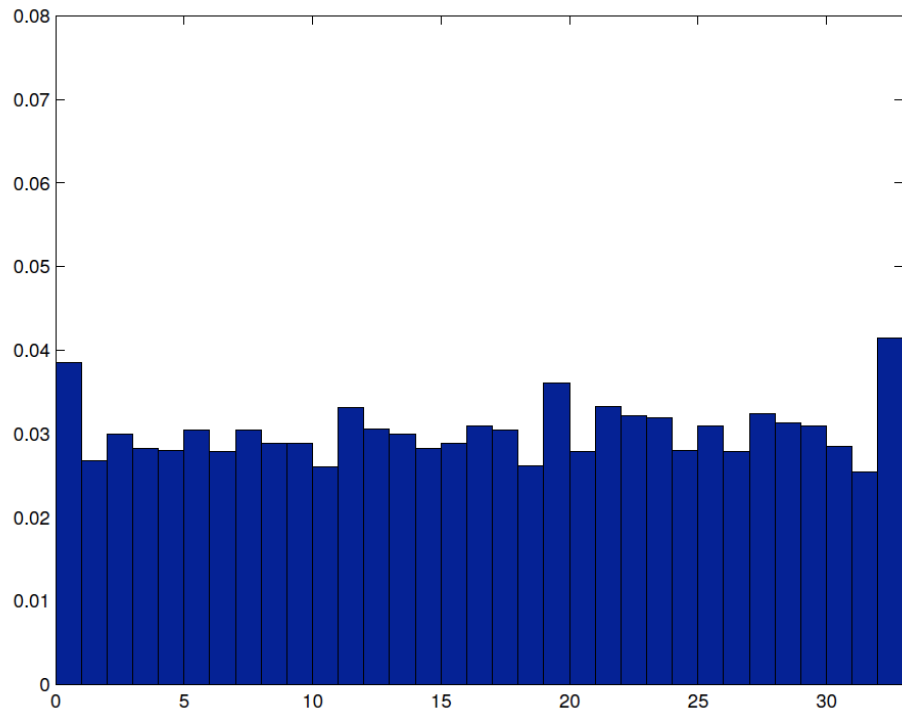


128 particles

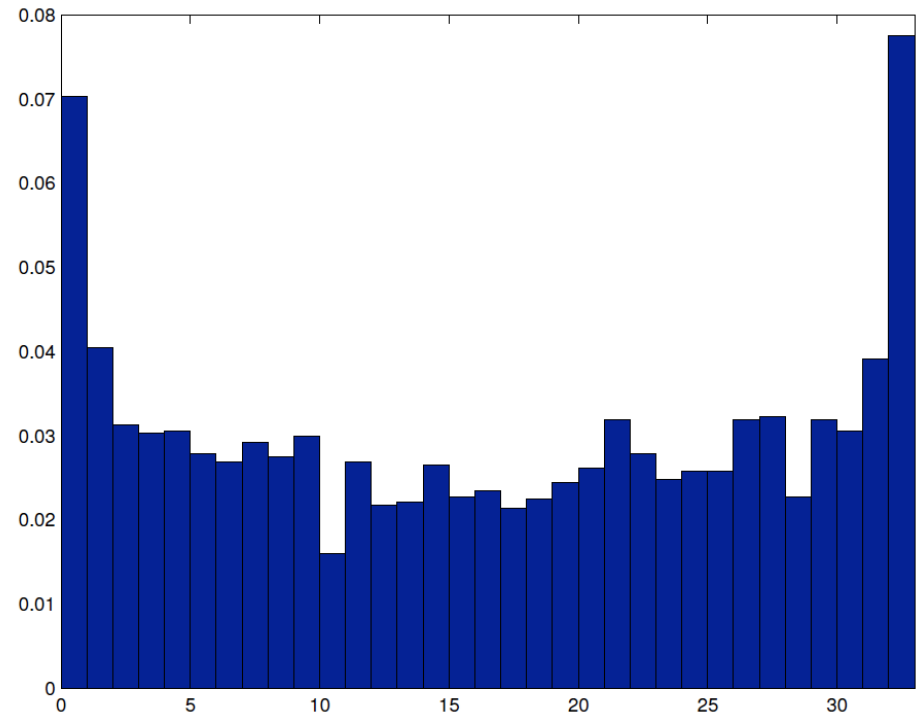


512 particles

Rank histograms

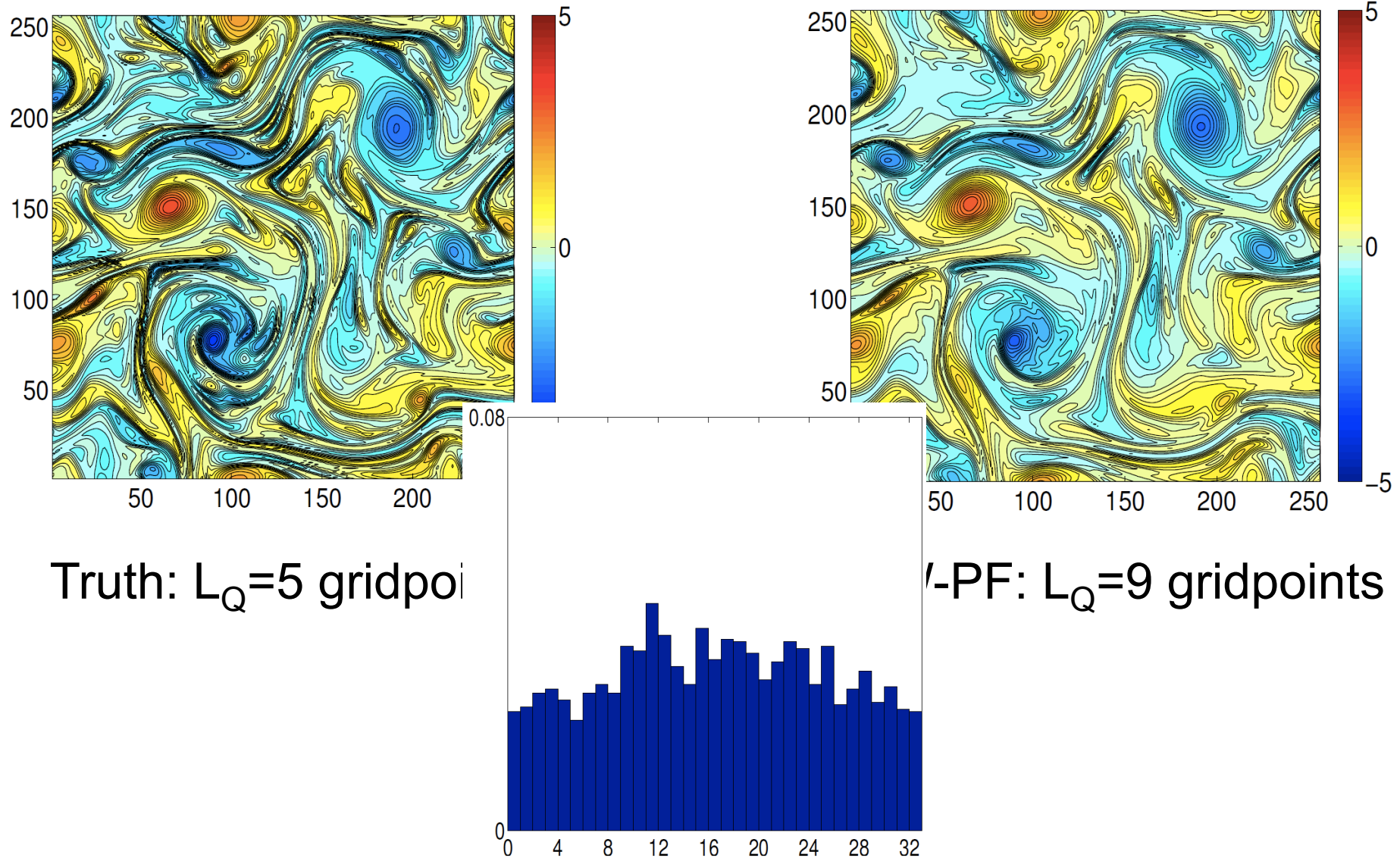


Full state observed

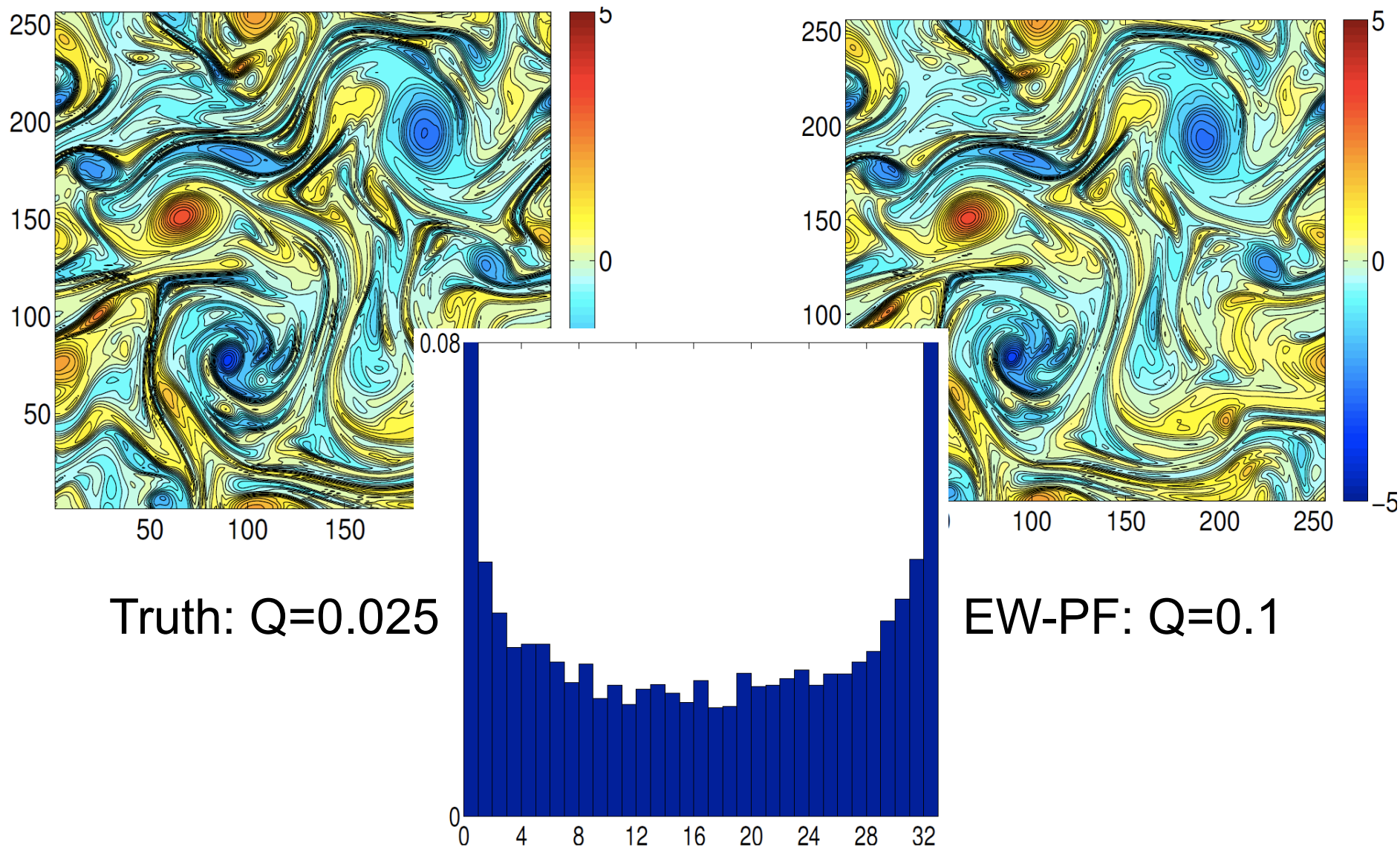


$\frac{1}{4}$ of half state observed

Miss-specification of process noise



Miss-specification of process noise



Conclusions

- Particle filters do not need state covariances.
- Degeneracy is related to number of observations, **not** to size of the state space.
- Proposal density allows enormous freedom
- Equivalent-weights scheme solves dimensionality problem?
- Other efficient schemes are being derived.
- Present work: numerical weather prediction, climate forecasting

We need more people !

- The Data Assimilation group at reading consists of 30 scientists
- We still have room in the

*Data Assimilation and Inverse Methods
in Geosciences MSc programme*

References

All references can be found on my website

<http://www.met.reading.ac.uk/~xv901096/research/publications.html>

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