



- 1 (a) Let $x(t)$ be given by $\dot{x}(t) = 1, x(0) = x_0$ and define $z(t) = u(x(t), t)$. Then $\dot{z}(t) = u_t + u_x \cdot 1 = -u(x(t), t) = -z(t)$, and the characteristic equations are

$$\begin{aligned}\dot{x} &= 1, & x(0) &= x_0, \\ \dot{z} &= -z, & z(0) &= u_0(x_0).\end{aligned}$$

This gives $x = x_0 + t$ and $z(t) = u_0(x_0)e^{-t}$. But x_0 is not given, we want to compute $u(x, t)$. We calculate x_0 from $x(t)$, $x_0 = x - t$ and thus the solution is $u(x, t) = u_0(x - t)e^{-t}$.

- (b) Let $x(t)$ be given by $\dot{x}(t) = 1, x(0) = x_0$ and define $z(t) = u(x(t), t)$. Then $\dot{z}(t) = u_t + u_x \cdot 1 = x(t)$, and the characteristic equations are

$$\begin{aligned}\dot{x} &= 1, & x(0) &= x_0, \\ \dot{z} &= x, & z(0) &= u_0(x_0).\end{aligned}$$

The solutions are $x(t) = x_0 + t, z(t) = z_0 + x_0t + \frac{1}{2}t^2$. This gives $u(x, t) = u_0(x - t) + xt - \frac{1}{2}t^2$.

Remark: For equations of the type $u_t + f(u, x)u_x = g(u, x)$ let $x(t) = f(z, x), z(t) = u(x(t), t)$. Then

$$\begin{aligned}\dot{z} &= u_t(x(t), t) + u_x(x(t), t)\dot{x}(t) \\ &= u_t + f(u, x)u_x = g(u, x) \\ &= g(z, x),\end{aligned}$$

and the characteristic equations are $\dot{x} = f(z, x), \dot{z} = g(z, x)$. We need to be able to solve the equations and find x_0, z_0 expressed with x, t, u_0 .

- 2 a) From the general conservation law it follows immediately that

$$\frac{d}{dt} \int_R (\Phi\rho) dV + \int_{\partial R} \mathbf{j} \cdot \mathbf{n} d\sigma = \int_R q dV.$$

We assume that the wells are δ -functions (which is OK assuming the reservoir is large compared to the wells),

$$q(\mathbf{x}, t) = \sum_{n=1}^2 q_n(t) \delta(\mathbf{x} - \mathbf{x}_n).$$

Thus

$$\frac{d}{dt} \int_R (\Phi\rho) dV + \int_{\partial R} \mathbf{j} \cdot \mathbf{n} d\sigma = q_1 + q_2.$$

If ρ is constant and the properties of R are unchanging in time, then

$$\frac{d}{dt} \int_R (\Phi \rho) dV = 0,$$

and the balance of the water is

$$\int_{\partial R} \mathbf{j} \cdot \mathbf{n} d\sigma = q_1 + q_2.$$

b) When

$$[p] = \frac{\text{kg} \cdot \text{m}}{\text{s}^2} \frac{1}{\text{m}^2} = \frac{\text{kg}}{\text{s}^2 \cdot \text{m}}.$$

$$[\nabla p] = \frac{\text{kg} \cdot \text{m}}{\text{s}^2} \frac{1}{\text{m}^3} = \frac{\text{kg}}{\text{s}^2 \cdot \text{m}^2}.$$

the dimension matrix is

	\mathbf{j}	K	μ	∇p
kg	0	0	1	1
m	1	2	-1	-2
s	-1	0	-1	-2

We see that the rank is 3, and we have a dimensionless combination, which we find easily. Strictly speaking, we derive the relation for scalar quantities, and use physical insight to find the given form. It may be noted that K in general will not be a scalar, but a 2nd order tensor (This happens typically in layered rocks where the pores are mostly in one direction).

c) We assume that the rod lies along the x -axis and has a constant cross section A . The conservation law for a section of the rod is then

$$\frac{d}{dt} \int_a^b (S_i \Phi) A dx + j_i(b) A - j_i(a) A = 0, \quad i = o, v.$$

We get the differential formulation in the usual way,

$$\Phi \frac{\partial S_i}{\partial t} + \frac{\partial j_i}{\partial x} = 0, \quad i = o, v.$$

We eliminate $\frac{\partial p}{\partial x}$ by adding the equations for the flux:

$$- \left(k_o(S_o) \frac{K}{\mu_o} + k_v(S_v) \frac{K}{\mu_v} \right) \frac{\partial p}{\partial x} = q,$$

i.e.

$$\frac{\partial p}{\partial x} = - \frac{q}{\left(k_o(S_o) \frac{K}{\mu_o} + k_v(S_v) \frac{K}{\mu_v} \right)}.$$

By inserting this in the equation for $S \equiv S_v$, we get

$$\begin{aligned} \Phi \frac{\partial S}{\partial t} + \frac{\partial j_v}{\partial x} &= \\ &= \Phi \frac{\partial S}{\partial t} + \frac{\partial}{\partial x} \left(k_v(S) \frac{K}{\mu_v} \frac{q}{\left(k_o(1-S) \frac{K}{\mu_o} + k_v(S) \frac{K}{\mu_v} \right)} \right) \\ &= \Phi \frac{\partial S}{\partial t} + \frac{\partial}{\partial x} f(S) = 0. \end{aligned}$$

- d) By inserting the given quantities, we see that $f(S) = qS^2$ and the equation reduces to

$$\frac{\partial S}{\partial t} + \frac{q}{\Phi} \frac{\partial S^2}{\partial x} = 0.$$

This is a regular first order hyperbolic equation with kinematic velocity

$$c(S) = \frac{2q}{\Phi} S.$$

To study this more carefully, we set up the equation for a characteristic starting point x_0 , where $0 < x_0 < L$:

$$x = x_0 + \frac{2q}{\Phi} S(x_0)t = x_0 + \frac{2q}{\Phi} \left(1 - \frac{x_0}{L}\right) t.$$

We leave to the reader to plot the characteristics. If we do this correctly, we see that all the characteristics meet in the point $x_s = L$, $t_s = \frac{\Phi L}{2q}$. Furthermore, the solution $S = 1$ above the characteristic that starts in the origin, i.e. $x = \frac{2qt}{\Phi}$. For a point (x, t) below this characteristic, we have

$$x = x_0 + \frac{2q}{\Phi} \left(1 - \frac{x_0}{L}\right) t,$$

i.e.

$$x_0 = L \frac{2qt - x\Phi}{2qt - \Phi L}$$

and

$$S(x, t) = 1 - \frac{2qt - x\Phi}{2qt - \Phi L}.$$

- 3** (1.13 in Ockendon, Howison, Lacey, Movchan, Applied Partial Differential Equations)

- a) Rankine-Hugoniot: for $u_t + \frac{\partial}{\partial x} f(u) = 0$:

$$\frac{dS}{dt}(u^+ - u^-) = f(u^+) - f(u^-)$$

Here $f(u) = \frac{1}{3}u^3$, so

$$\frac{dS}{dt} = \frac{1}{3} \frac{(u^+)^3 - (u^-)^3}{u^+ - u^-}.$$

Characteristic equations: $z(t) = u(x(t), t)$

$$\begin{cases} \dot{x} = f'(z) = z^2, & x(0) = x_0 \\ \dot{z} = 0, & z(0) = u(x(0), 0) = u_0(x_0) \end{cases}$$

Solution:

$$\begin{cases} z = \text{const} = u_0(x_0) \\ x = x_0 + t(u_0(x_0))^2 \end{cases}$$

Hence characteristics are straight lines.

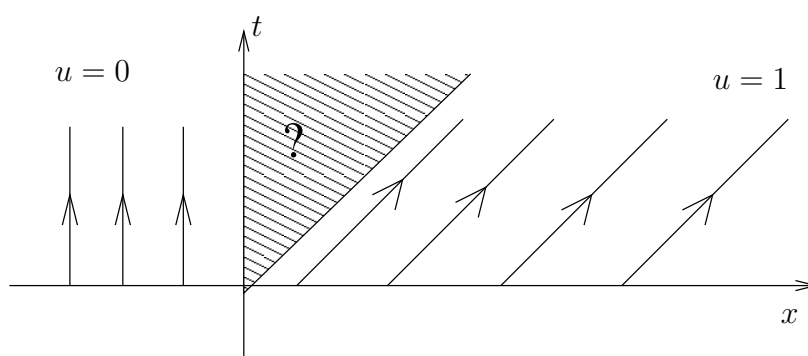
b) The problem consists of 2 Riemann problems:

$$(1) \quad u_t + u^2 u_x = 0, \quad u(x, 0) = \begin{cases} 0 & x < 0 \\ 1 & x > 0 \end{cases} \quad (x < 1)$$

$$(2) \quad u_t + u^2 u_x = 0, \quad u(x, 0) = \begin{cases} 1 & x < 1 \\ 0 & x > 1 \end{cases} \quad (x > 0)$$

For (1):

$$x = x_0 + t \cdot \begin{cases} 0 & x_0 < 0 \\ 1 & x_0 > 0 \end{cases} \quad (x_0 < 1)$$



There is a "dead sector" and we know we have to fill it with a rarefaction fan:

$$(3) \quad u(x, t) = \phi\left(\frac{x}{t}\right)$$

i.e. a function constant at rays out of $(0, 0)$. Insert (3) into (1) to find that

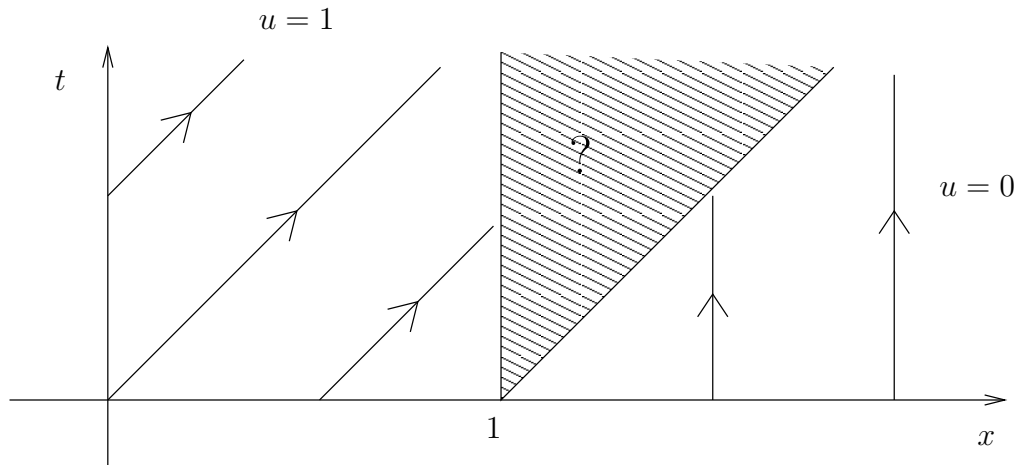
$$\begin{aligned} \phi'\left(\frac{x}{t}\right) \left(-\frac{x}{t^2}\right) + \phi^2 \phi'\left(\frac{x}{t}\right) \frac{1}{t} &= 0 \\ \Rightarrow \phi^2 &= \frac{x}{t}, \quad \Rightarrow \phi = \sqrt{\frac{x}{t}} \end{aligned}$$

Solution:

$$u(x, t) = \begin{cases} 0, & x < 0 \\ \sqrt{\frac{x}{t}}, & 0 < x < t \\ 1, & x > t \end{cases} \quad (\text{but not for big})$$

For (2):

$$x = x_0 + t \cdot \begin{cases} 1 & x_0 < 1 \\ 0 & x_0 > 1 \end{cases} \quad (x_0 > 0)$$



Characteristics collide at $x = 1$, hence there is a shock starting at $(x, t) = (1, 0)$ and moving with speed

$$\frac{dS}{dt} = \frac{1}{3} \frac{0^3 - 1^3}{0 - 1} = \frac{1}{3}.$$

Hence

$$S(t) = 1 + \frac{1}{3}t$$

The shock solution then becomes

$$u(x, t) = \begin{cases} 1 & x < S(t) \\ 0 & x > S(t) \end{cases} \quad (x > t)$$

Total solution

$$u(x, t) = \begin{cases} 0 & x < 0 \\ \sqrt{\frac{x}{t}} & 0 < x < t \\ 1 & t < x < 1 + \frac{1}{3}t \\ 0 & x > 1 + \frac{1}{3}t \end{cases}$$

Obs: Solution only defined as long as $t \leq 1 + \frac{1}{3}t$ or $t \leq \frac{3}{2}$.

- c) For $t > \frac{3}{2}$, the left state of shock changes from (1) to $\sqrt{\frac{x}{t}} = \sqrt{\frac{S(t)}{t}} (< 0)$ and hence the speed of the shock changes to

$$\begin{aligned} \frac{dS}{dt} &= \frac{1}{3} \frac{(u^+)^3 - (u^-)^3}{u^+ - u^-} \\ &= \frac{1}{3} \frac{0^3 - \left(\frac{S(t)}{t}\right)^{3/2}}{0 - \sqrt{\frac{S(t)}{t}}} \\ &= \frac{1}{3} \frac{S(t)}{t} \end{aligned}$$

Since

$$S\left(t = \frac{3}{2}\right) = \frac{3}{2},$$

we solve (separable equations) and find that

$$S(t) = \left(\frac{3}{2}\right)^{\frac{2}{3}} t^{1/3}$$

The new solution is then

$$u(x, t) = \begin{cases} 0, & x < 0 \\ \sqrt{\frac{x}{t}}, & 0 < x < S(t) \\ 0, & S(t) < x \end{cases}$$

and this solution persists until $t = +\infty$.