

- 1 a) Let us write

$$U_T = (a, b) \times (0, T),$$

so the closure is

$$\bar{U}_T = [a, b] \times [0, T]$$

As usual, we denote by  $\Gamma_T$  the parabolic boundary, consisting of the bottom and vertical sides of the boundary of  $[a, b] \times [0, T]$ . Thus,

$$\Gamma_T = \{(x, t) \in \bar{U}_T : x = a \text{ or } x = b \text{ or } t = 0\}.$$

Now suppose  $u \in C_1^2(U_T) \cap C(\bar{U}_T)$  is a solution of the heat equation  $u_t = u_{xx}$  in  $U_T$ . Define

$$M = \max_{\bar{U}_T} u.$$

Then the weak maximum principle says that *the maximum  $M$  is attained at some point on the parabolic boundary  $\Gamma_T$* . In symbols,

$$\max_{\Gamma_T} u = M.$$

- b) It is trivial to verify that  $u = -2xt - x^2$  satisfies  $u_t = xu_{xx}$ . To check whether  $u$  violates the weak maximum principle, we need to locate the maximum of  $u$  in the rectangle  $R = [-2, 2] \times [0, 1]$ . We first check interior points, so we solve

$$u_t = u_x = 0,$$

which has only one solution, namely  $x = t = 0$ , where  $u = 0$ . Now we check the boundary. For  $t = 0$  we have  $u \leq 0$ . For  $x = -2$  and  $0 \leq t \leq 1$ , we have  $u = 4(t - 1) \leq 0$ . For  $x = 2$  and  $0 \leq t \leq 1$ , we have  $u = -4(t + 1) < 0$ . Finally, for  $t = 1$  we have  $u = -2x - x^2$ , which has maximum  $u = 1$  at  $x = -1$ . So clearly, the maximum of  $u$  over  $R$  is  $u = 1$ , and it occurs *only* at the top part of the boundary, at the point  $(x, t) = (-1, 1)$ , hence the weak maximum principle is violated.

- 2 We have two jumps in the initial data. The first jump (at  $x = 0$ ) gives a rarefaction solution, the other jump (at  $x = 1$ ) gives a shock solution, which by Rankine-Hugoniot (with  $F(u) = u^2/2$ ) moves with speed  $1/2$  to the right. Combining the rarefaction and shock solutions, we get

$$(1) \quad u(x, t) = \begin{cases} 0 & \text{for } x \leq 0, \\ \frac{x}{t} & \text{for } 0 \leq x \leq t, \\ 1 & \text{for } t \leq x \leq 1 + \frac{t}{2}, \\ 0 & \text{for } x \geq 1 + \frac{t}{2}. \end{cases}$$

Note, however, that this is valid only until the time  $t$  when the rarefaction and the shock meet:

$$t = 1 + \frac{t}{2} \quad \text{i.e.,} \quad t = 2.$$

At this time we have

$$u(x, 2) = \begin{cases} 0 & \text{for } x \leq 0, \\ \frac{x}{2} & \text{for } 0 \leq x \leq 2, \\ 0 & \text{for } x \geq 2. \end{cases}$$

So now we solve the initial value problem starting from  $t = 2$  with these initial data. Then we get a shock emanating from  $x = 2$ , moving to the right along a path

$$x = \xi(t),$$

and to the left of the shock we have the continuation of the rarefaction wave,

$$u = u_l = \frac{x}{t},$$

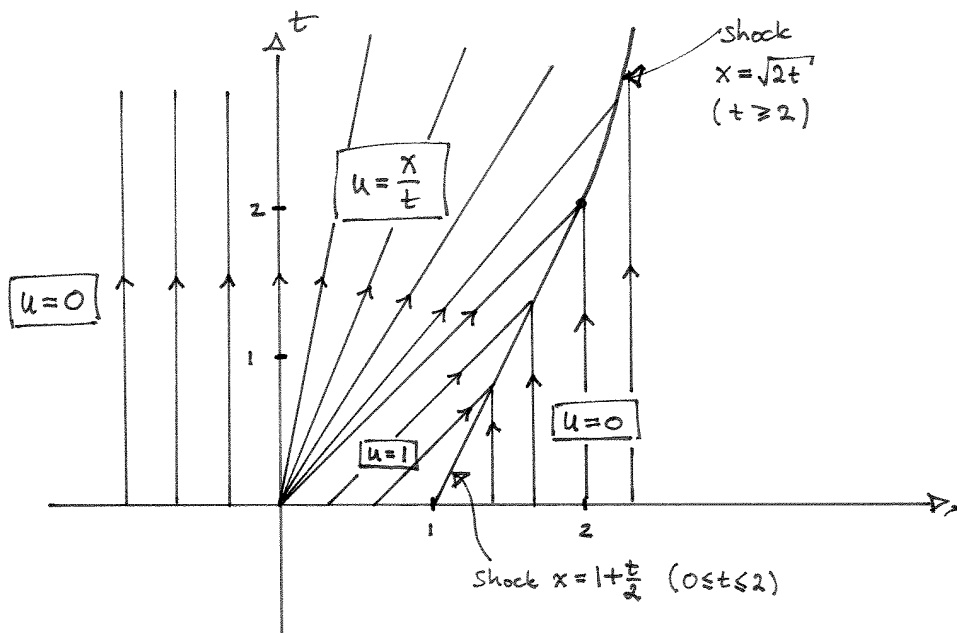
while  $u = u_r = 0$  to the right. Thus, by Rankine-Hugoniot,

$$\frac{d\xi}{dt} = \frac{1}{2}u_l = \frac{\xi}{2t},$$

which we separate and solve to get  $\xi = \sqrt{2t}$ , since the initial condition is  $\xi(2) = 2$ . We conclude that for  $t \geq 2$ ,

$$(2) \quad u(x, t) = \begin{cases} 0 & \text{for } x \leq 0, \\ \frac{x}{t} & \text{for } 0 \leq x \leq \sqrt{2t}, \\ 0 & \text{for } x \geq \sqrt{2t}, \end{cases}$$

which together with (1) for  $0 \leq t \leq 2$  provides the complete solution.



- 3] Suppose  $u, v : \bar{U} \times [0, T]$  are two smooth solution of the initial/boundary value problem. Then  $w = u - v$  is a solution of the same problem, but with  $f, g, h = 0$ . Now define

$$E(t) = \frac{1}{2} \int_U w_t^2 + |\nabla w|^2 dx,$$

where the integrand is understood to be evaluated at  $(x, t)$ . Then

$$\begin{aligned} E'(t) &= \int_U w_t w_{tt} + \nabla w \cdot \nabla w_t \, dx \\ &= \int_U w_t (w_{tt} - \Delta w) \, dx + \int_{\partial U} \frac{\partial w}{\partial \nu} w_t \, dx \\ &= 0 \end{aligned}$$

where we used integration by parts and the fact that

$$\frac{\partial w}{\partial \nu} = 0 \quad \text{on } \partial U \times [0, T].$$

Thus,  $0 \leq E(t) \leq E(0)$  for  $0 \leq t \leq T$ , but  $E(0) = 0$ , since  $f = g = 0$ . Thus  $Dw = 0$  in  $U \times [0, T]$ , hence  $w = \text{const.}$  in  $U \times [0, T]$ , but this constant must be zero, since we have zero initial data.

4 For simplicity let's switch to  $(x_1, x_2, x_3)$  instead of  $(x, y, z)$ . By Kirchhoff's formula,

$$u(x, t) = D_t \left( \frac{t}{4\pi} \int_{|y|=1} f(x + ty) \, dS(y) \right),$$

where

$$f(x) = x_1^2 + x_2^2.$$

Note that

$$f(x + ty) = x_1^2 + x_2^2 + 2t(x_1 y_1 + x_2 y_2) + t^2(y_1^2 + y_2^2).$$

Now we use:

$$\begin{aligned} \int_{|y|=1} y_i \, dS(y) &= 0, \\ \int_{|y|=1} 1 \, dS(y) &= 4\pi, \\ \int_{|y|=1} y_i^2 \, dS(y) &= \frac{1}{3} \int_{|y|=1} \sum_{j=1}^3 y_j^2 \, dS(y) = \frac{1}{3} \int_{|y|=1} 1 \, dS(y) = \frac{4\pi}{3}. \end{aligned}$$

It follows that

$$\frac{t}{4\pi} \int_{|y|=1} f(x + ty) \, dS(y) = t(x_1^2 + x_2^2) + \frac{2t^3}{3},$$

hence

$$u = \underline{\underline{x_1^2 + x_2^2 + 2t^2}}.$$

5 a) A weak solution of the given Dirichlet problem is a function  $u \in H_0^1(U)$  such that

$$(3) \quad \underbrace{\int_U \nabla u \cdot \nabla v \, dx}_{=B[u,v]} = \int_U f v \, dx \quad \text{for all } v \in H_0^1(U).$$

By the Poincaré inequality, which reads (here  $C$  depends only on  $U$ )

$$\|u\|_{L^2(U)} \leq C \|Du\|_{L^2(U)} \quad \text{for all } u \in H_0^1(U),$$

we have that  $B[u, u] = 0 \implies u = 0$ , so  $B[u, v]$  is an inner product on  $H_0^1(U)$ . Moreover, the norm  $\|u\| = \sqrt{B[u, u]} = \|Du\|_{L^2(U)}$  is equivalent to the standard norm on  $H_0^1(U)$ , again by Poincaré's inequality. Indeed,

$$\|u\| \leq \|u\|_{H_0^1(U)} = \|u\|_{L^2(U)} + \|Du\|_{L^2(U)} \leq C \|u\|.$$

Finally, existence of a weak solution follows by the Riesz representation theorem applied to the Hilbert space  $H_0^1(U)$ , since  $v \mapsto \int_U f v dx$  is a bounded linear functional on  $H_0^1(U)$ :

$$\left| \int_U f v dx \right| \leq \|f\|_{L^2(U)} \|u\|_{L^2(U)} \leq C \|f\|_{L^2(U)} \|u\|,$$

for all  $v \in H_0^1(U)$ .

**b)** Suppose  $u \in H_0^1(U)$  is a weak solution and  $\text{supp } u \subset\subset U$ . Let  $k \in \{1, \dots, n\}$ . We take

$$v = D_k^{-h} D_k^h u$$

in (3). Then

$$\underbrace{\int_U \nabla u \cdot \nabla D_k^{-h} D_k^h u dx}_{=A} = \underbrace{\int_U f D_k^{-h} D_k^h u dx}_{=B}.$$

By "integration by parts" for difference quotients (valid by the support assumption on  $u$ ),

$$A = - \int_U \nabla D_k^h u \cdot \nabla D_k^h u dx = - \|\nabla D_k^h u\|_{L^2(U)}^2 = - \|D_k^h \nabla u\|_{L^2(U)}^2.$$

By Cauchy's inequality with  $\varepsilon$ , we have

$$|B| \leq \frac{1}{4\varepsilon} \|f\|_{L^2(U)}^2 + \varepsilon \|D_k^{-h} D_k^h u\|_{L^2(U)}^2 \leq \frac{1}{4\varepsilon} \|f\|_{L^2(U)}^2 + \varepsilon \|\nabla D_k^h u\|_{L^2(U)}^2,$$

where we used Theorem 3(i) in §5.8.2 of Evans (which holds with constant  $C = 1$ , actually) to get the last inequality.

We conclude:

$$\|D_k^h \nabla u\|_{L^2(U)}^2 = -A = -B \leq \frac{1}{4\varepsilon} \|f\|_{L^2(U)}^2 + \varepsilon \|D_k^h \nabla u\|_{L^2(U)}^2,$$

hence, choosing  $\varepsilon = 1/2$ ,

$$\|D_k^h \nabla u\|_{L^2(U)}^2 \leq \|f\|_{L^2(U)}^2.$$

Since this holds for all small enough  $h \neq 0$ , it follows by Theorem 3(ii) in §5.8.2 that the second order weak partial derivatives exist, and

$$\|u_{x_j x_k}\|_{L^2(U)}^2 \leq \|f\|_{L^2(U)}^2 \quad (j, k = 1, \dots, n).$$

Therefore  $u \in H^2(U)$ .