



Exam in TMA4305 Partial differential Equations

Suggested solutions: May 2009

Problem 1 Characteristic equations

$$\begin{cases} \dot{x} = x, & x(0) = x_0, \\ \dot{y} = -y, & y(0) = 1, \\ \dot{z} = -z, & z(0) = u(x_0, 1) = h(x_0), \end{cases}$$

with solutions

$$x = x_0 e^t, \quad y = e^{-t}, \quad z = h(x_0) e^{-t}.$$

In terms of (x, y) we get $x_0 = x e^{-t}$ and $e^{-t} = y$, thus,

$$u(x, y) = z(t(x, y); x_0(x, y)) = h(xy)y.$$

Problem 2 Characteristic equations

$$\begin{cases} \dot{t} = 1, & t(0) = 0, \\ \dot{x} = e^z, & x(0) = x_0, \\ \dot{z} = 0, & z(0) = u(x_0, 0) = \begin{cases} 1, & x < 0, \\ 2, & x > 0. \end{cases} \end{cases}$$

We solve for z before x , and get

$$\begin{cases} t = t, \\ x = x_0 + t e^{u(x_0, 0)} = \begin{cases} x_0 + te, & x_0 < 0, \\ x_0 + te^2, & x_0 > 0, \end{cases} \\ z = u(x_0, 0) = \begin{cases} 1, & x < 0, \\ 2, & x > 0. \end{cases} \end{cases}$$

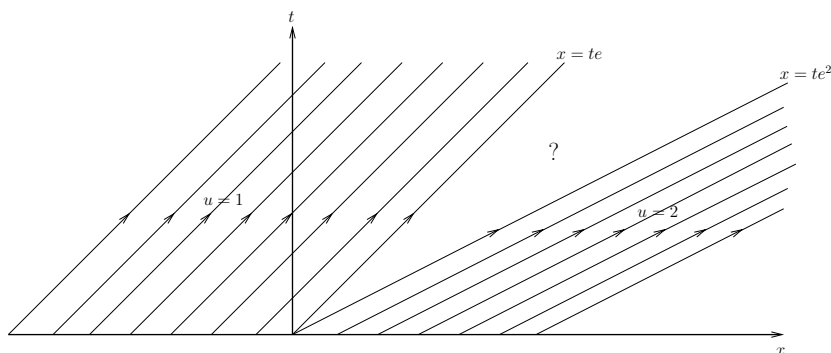


Figure 1: The projected characteristic curves of Problem 2.

See Figure 1 for a sketch of the solution.

The rarefaction fan is given by $u(x, t) = \phi(x/t)$ satisfying the PDE

$$-\frac{x}{t^2}\phi' + e^\phi \frac{1}{t}\phi' = 0.$$

We get $e^\phi = \frac{x}{t}$, that is, $\phi = \ln \frac{x}{t}$. (Note that $\phi' = 0$ implies ϕ constant and does not give a rarefaction fan.)

The total solution is

$$u(x, t) = \begin{cases} 1, & x \leq te, \\ \ln \frac{x}{t}, & te \leq x \leq te^2, \\ 2, & x \geq te^2. \end{cases}$$

Remarks:

- Continuous, piecewise smooth solutions are weak solutions by Rankine-Hugoniot.
- This is the entropy solution since no projected characteristics collide.

Problem 3 We will need the following Green's identity (can be derived using the Divergence Theorem):

$$\int_{\partial\Omega} v \frac{\partial u}{\partial \nu} dS = \int_{\Omega} (v\Delta u + \nabla v \cdot \nabla u) dx, \quad (1)$$

Assume that there are two solutions, u and v , and let $w = u - v$. Note that w is in $C^2(\Omega)$ and solves

$$\begin{aligned} \Delta w - cw &= 0, & \text{in } \Omega, \\ \frac{\partial w}{\partial \nu} + \lambda w &= 0, & \text{on } \partial\Omega. \end{aligned} \quad (2)$$

By Green's identity (1) we have

$$\int_{\partial\Omega} w \frac{\partial w}{\partial \nu} = \int_{\Omega} (w\Delta w + |\nabla w|^2),$$

and using equation (2) we get

$$\begin{aligned} - \int_{\partial\Omega} \lambda w^2 &= \int_{\Omega} (cw^2 + |\nabla w|^2), \\ \int_{\Omega} (cw^2 + |\nabla w|^2) + \int_{\partial\Omega} \lambda w^2 &= 0. \end{aligned}$$

Since $w \in C^2(\overline{\Omega})$, this implies that

$$\begin{cases} cw^2 + |\nabla w|^2 = 0, & \text{in } \Omega, \\ \lambda w^2 = 0, & \text{on } \partial\Omega. \end{cases} \quad (3)$$

Recall that c and λ are non-negative and $w \in C(\overline{\Omega})$. If $c + \lambda > 0$, it then follows from (3) that $w \equiv 0$ in $\overline{\Omega}$. Hence, $u \equiv v$ in $\overline{\Omega}$ and the solutions are unique.

Problem 4

a) 1. Observe that in Ω

$$Lv_\epsilon = Lu + \epsilon L(x^2 + y^2) \geq 0 + \epsilon (2(1 + y^2) + 2(1 + x^2) + 2x^2 + 2y^2) > 0. \quad (4)$$

2. Since v_ϵ is continuous and $\overline{\Omega}$ compact, there is $(\bar{x}, \bar{y}) \in \overline{\Omega}$ so that $v_\epsilon(\bar{x}, \bar{y}) = \max_{\overline{\Omega}} v_\epsilon$.

3. If $(\bar{x}, \bar{y}) \in \Omega$, then at (\bar{x}, \bar{y})

$$v_{\epsilon x} = 0 = v_{\epsilon y}, \quad \sum_{i=1}^2 \xi_i v_{\epsilon x_i x_j} \xi_j \leq 0 \quad \text{for all } \xi \in \mathbb{R}^2.$$

In particular, take first $\xi = (1, 0)$ and then $\xi = (0, 1)$ to get

$$v_{\epsilon xx} \leq 0, \quad v_{\epsilon yy} \leq 0 \quad \text{at } (\bar{x}, \bar{y})$$

Hence,

$$Lv_\epsilon(\bar{x}, \bar{y}) \leq (1 + \bar{x}^2) \cdot 0 + (1 + \bar{y}^2) \cdot 0 + 0 + 0 = 0,$$

which contradicts (4). Therefore, $(\bar{x}, \bar{y}) \in \partial\Omega$, and for $(x, y) \in \Omega$

$$v_\epsilon(x, y) \leq v_\epsilon(\bar{x}, \bar{y}) = \max_{\partial\Omega} v_\epsilon \quad (5)$$

4. We now have

$$u(x, y) \leq v_\epsilon(x, y) \leq \max_{\partial\Omega} v_\epsilon \leq \max_{\partial\Omega} (u + \epsilon(x^2 + y^2)) \leq \max_{\partial\Omega} u + 2\epsilon(\text{diam}\Omega)^2.$$

Sending $\epsilon \rightarrow 0$, we get

$$\max_{\bar{\Omega}} u \leq \max_{\partial\Omega} u.$$

b) Assume there are two solutions, $u, v \in C^2(\Omega) \cap C(\bar{\Omega})$. Let $w = u - v$ and note that $w \in C^2(\Omega) \cap C(\bar{\Omega})$ solves

$$\begin{aligned} Lw &= 0, & \text{in } \Omega, \\ w &= 0, & \text{on } \partial\Omega. \end{aligned}$$

Hence, the assumptions in a) are satisfied by $\pm w$, and the weak maximum principle yields

$$\begin{aligned} \max_{\bar{\Omega}} w &= \max_{\partial\Omega} w = 0, \\ \max_{\bar{\Omega}} (-w) &= \max_{\partial\Omega} (-w) = 0. \end{aligned}$$

Therefore, $|w| = 0$ in Ω and $u \equiv v$, which proves uniqueness.

Problem 5

a)

$$\begin{aligned} F(u + tv) &= \int \frac{1}{2} (\Delta(u + tv))^2 - f(u + tv) \\ &= \int \frac{1}{2} \Delta u + t \Delta u \Delta v + \frac{1}{2} t^2 \Delta v^2 - fu + tfv \\ &= F(u) + t \int (\Delta u \Delta v - fv) + \frac{1}{2} t^2 \int \Delta v^2. \end{aligned}$$

$$D_v F(u) = \lim_{t \rightarrow 0} \frac{F(u + tv) - F(u)}{t} = \int_{\Omega} (\Delta u \Delta v - fv).$$

The Euler-Lagrange equation:

$$0 = D_v F(u) = \int_{\Omega} (\Delta u \Delta v - f v), \quad \text{for all } v \in H_0^2(\Omega).$$

b) By a) u satisfy

$$\int_{\Omega} \Delta u \Delta v = \int_{\Omega} f v, \quad \text{for all } v \in H_0^2(\Omega)$$

Since $C_0^\infty(\Omega) \subset H_0^2(\Omega)$;

$$\int_{\Omega} \Delta u \Delta \phi = \int_{\Omega} f \phi, \quad \text{for all } \phi \in C_0^\infty(\Omega). \quad (6)$$

Integration by parts;

$$\begin{aligned} \int_{\Omega} \Delta u \Delta \phi &= \sum_i \int_{\Omega} \Delta u \phi_{x_i x_i} \\ &\stackrel{\text{i.b.p.}}{=} \sum_i \left(\int_{\Omega} \partial_{x_i} (\Delta u \phi_{x_i}) - \int_{\Omega} \Delta u_{x_i} \phi_{x_i} \right) \\ &\stackrel{\text{div. thm.}}{=} \sum_i \left(\int_{\Omega} \Delta u \phi_{x_i} v_i - \int_{\Omega} \Delta u_{x_i} \phi_{x_i} \right) \\ &\stackrel{\phi \in C_0^\infty}{=} 0 - \sum_i \int_{\Omega} \Delta u_{x_i} \phi_{x_i} \\ &= - \sum_i \left(\int_{\Omega} \partial_{x_i} (\Delta u_{x_i} \phi) - \int_{\Omega} \Delta u_{x_i x_i} \phi \right) \\ &= 0 + \sum_i \int_{\Omega} \Delta u_{x_i x_i} \phi \\ &= \int_{\Omega} \Delta^2 u \phi. \end{aligned}$$

By (6),

$$\int_{\Omega} (\Delta^2 u - v) \phi = 0, \quad \text{for all } \phi \in C_0^\infty(\Omega),$$

and since $\Delta^2 u - v$ is continuous, it follows from the variational lemma that $\Delta^2 u = v$ in Ω .

c)

$$\begin{aligned}
F(u) &= \frac{1}{2}|u|_{2,2}^2 - \int f v \\
&\stackrel{\text{eq. (6)}}{\geq} \frac{1}{2C_\Omega}|u|_{2,2}^2 - \int f v \\
&\stackrel{ab \leq \frac{a^2}{2\epsilon} + \frac{b^2}{2}}{\geq} \frac{1}{2C_\Omega}|u|_{2,2}^2 - \frac{1}{2\epsilon}\|v\|_2^2 - \frac{\epsilon}{2}\|f\|_2^2 \\
&\stackrel{\epsilon=2C_\Omega}{=} \frac{1}{4}|u|_{2,2}^2 - C_\Omega\|f\|_2^2,
\end{aligned}$$

hence, F is coersive with $C_2 = \frac{1}{4}$ and $C_2 = C_\Omega\|f\|_2^2$.

d) The direct method:

1. Since F is coersive;

$$I := \inf_{v \in H_0^2} F(v) \geq C_1\|u\|_{2,2}^2 - C_2 \geq -C_2 > -\infty.$$

2. Take a minimixing sequence $\{u_i\} \subset H_0^2(\Omega)$ satisfying $F(u_i) \rightarrow I$ and $|F(u_i)| \leq I+1$ for all i .
3. $\|u\|_{2,2} \leq M$ for all i since

$$C_1\|u\|_{2,2}^2 \leq C_2 + F(u_i) \leq C_2 + I + 1 =: M^2 \leq \infty.$$

By weak compactness there exists a subsequence $\{u_{i_k}\} \subset \{u_i\}$, $u \in H_0^2(\Omega)$, so that $u_{i_k} \rightharpoonup u$ in $H_0^2(\Omega)$.

4. $F(u) \leq \liminf_{k \rightarrow \infty} F(u_{i_k})$: We will use that in any Hilbert space $(X, \|\cdot\|)$

$$u_i \rightharpoonup u \quad \implies \quad \|u\| \leq \liminf \|u_{i_k}\|,$$

and that

$$\begin{aligned}
u_i \rightharpoonup u \text{ in } L^2 &\implies \int f u_i \rightarrow \int f u. \\
F(u) &= \frac{1}{2}|u|_{2,2}^2 - \int f u \\
&\leq \liminf_{k \rightarrow \infty} \frac{1}{2}|u_{i_k}|_{2,2}^2 - \lim \int f u_{i_k} \\
&= \liminf F(u_{i_k}).
\end{aligned}$$

We can now concluded that

$$I \leq F(u) \leq \liminf F(u_{i_k}) = \lim F(u_{i_k}) = I,$$

and hence $F(u) = I$.

Conclusion: There exists a $u \in H_0^2(\Omega)$ so that $F(u) = I = \inf_{v \in H_0^2} F(v)$.