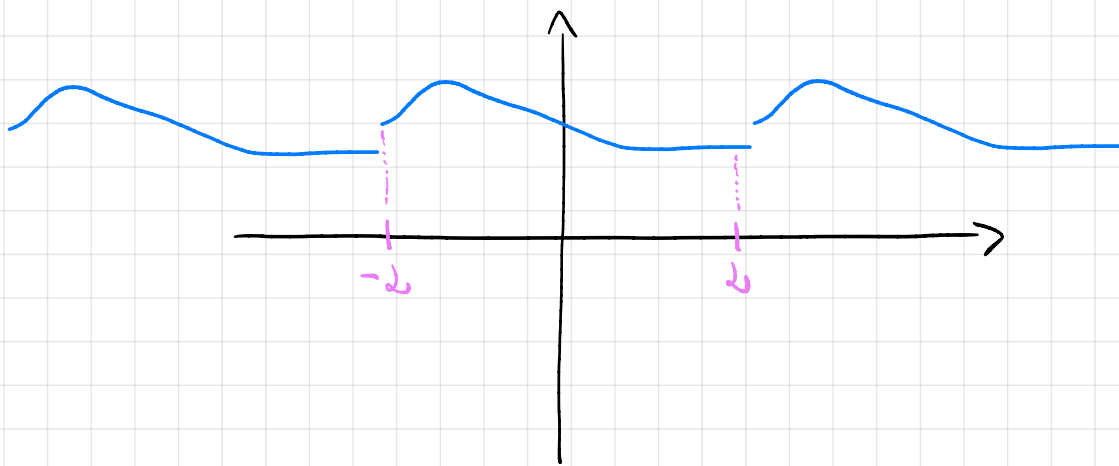
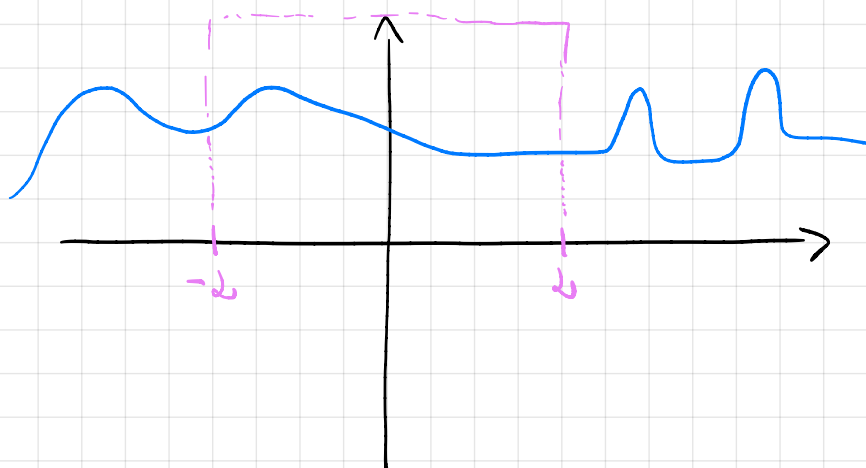


Fourier transform

Questions: Fourier series allows us to handle $f: \mathbb{R} \rightarrow \mathbb{C}$ which are 2d periodic. What about non-periodic functions?

Idea: Restrict f to interval $(-L, L)$, treat it as "2d" periodic function using Fourier series and try to understand what happens if $L \rightarrow \infty$.



Let's have a look at the resulting Fourier series:

$$f|_{(-L, L)} \sim \sum_{n=-\infty}^{\infty} \frac{1}{2L} \int_{-L}^L f(y) e^{-iy \frac{n\pi}{L}} dy e^{ix \frac{n\pi}{L}}$$

$$\Delta \omega := \frac{\pi}{L}$$

$$\omega_n := n \cdot \Delta \omega$$

$$= \sum_{n=-\infty}^{\infty} \Delta \omega \frac{1}{2\pi} \int_{-L}^L f(y) e^{-iy \omega_n} dy e^{ix \omega_n}$$

This looks like a Riemann sum for the infinite integral of $\hat{f}_L(\omega) := \frac{1}{2\pi} \int_{-L}^L f(y) e^{-iy \omega} dy$ Function in ω !

So with a lot of handwaving, we expect that for $\lambda \rightarrow \infty$ we obtain

$$(*) \quad f(x) \sim \int_{-\infty}^{\infty} \frac{1}{2\pi} \int_{-\infty}^{\infty} f(y) e^{-i\omega y} dy e^{i\omega x} d\omega.$$

Definition (Fourier transform)

Assume that $\int_{-\infty}^{\infty} |f(x)| dx$ exists and $< \infty$. Then

the Fourier transform $\mathcal{F}(f)$ or $\hat{f}$ is defined by

$$\hat{f}(\omega) := \mathcal{F}(f)(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx$$

If $g(\omega)$ is absolutely integrable, then the inverse Fourier transform is defined by

$$\check{g}(x) = \mathcal{F}^{-1}(g) := \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(\omega) e^{i\omega x} d\omega.$$

Now $(*)$ translates to

Theorem

Let f and $\hat{f}$ be absolutely integrable, then

$$f(x) = \mathcal{F}^{-1}(\hat{f}) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{i\omega x} d\omega.$$

Proof. Omitted.

Warning (Different versions of $\mathcal{F}, \mathcal{F}^{-1} \dots$)

There exists a lot of different definitions for $\mathcal{F}$, as there is no agreement on whether 1 has to absorb the $\frac{1}{2\pi}$ in the definition.

Common other choices.

$$\begin{array}{ccc} \boxed{\int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx} & \xrightarrow{\mathcal{F}^{-1}} & \boxed{\frac{1}{2\pi} \int_{-\infty}^{\infty} f(\omega) e^{i\omega x} d\omega} \\ \int_{-\infty}^{\infty} f(x) e^{-i2\pi\omega x} dx & & \int_{-\infty}^{\infty} f(\omega) e^{i2\pi\omega x} d\omega \end{array}$$

Either in front of $\mathcal{F}$ or $\mathcal{F}^{-1}$.

So if you see a Fourier transformed function, check which def. was used?

Examples of Fourier transforms

Example 1: $f(x) = \chi_{[a,b]} = \begin{cases} 1 & \text{if } x \in [a,b] \\ 0 & \text{else.} \end{cases}$

$$\hat{f}(0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \chi_{[a,b]} \cdot 1 dx = \frac{1}{\sqrt{2\pi}} (b-a).$$

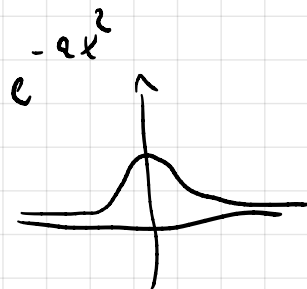
If $\omega \neq 0$ we compute that

$$\begin{aligned} \hat{f}(\omega) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \chi_{[a,b]}(x) e^{-i\omega x} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_a^b e^{-i\omega x} dx = \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{-i\omega} e^{-i\omega x} \Big|_{x=a}^b \\ &= \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{-i\omega} \cdot \left(e^{-i\omega b} - e^{-i\omega a} \right) \underbrace{e^{i\omega \frac{(a+b)}{2}} \cdot e^{-i\omega \frac{(a+b)}{2}}}_{=1} \\ &= \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{\omega} \cdot \left(\frac{e^{-i\omega \frac{(b-a)}{2}} - e^{i\omega \frac{(b-a)}{2}}}{-i \cdot 2} \right) \cdot e^{-i\omega \frac{(a+b)}{2}} \\ &= \sqrt{\frac{2}{\pi}} \cdot \frac{\sin(\omega \cdot \frac{(b-a)}{2})}{\omega} \cdot e^{-i\omega \frac{(a+b)}{2}} \end{aligned}$$

Example 2: let $a > 0$, and set $f(x) = e^{-a|x|}$. Then

$$\mathcal{F}(e^{-a|x|}) = \sqrt{\frac{2}{\pi}} \cdot \frac{a}{a^2 + \omega^2}$$

See plenary exercise for complete calculation



Example 3: let $a > 0$, then

$$\mathcal{F}(e^{-ax^2}) = \frac{1}{\sqrt{2a}} e^{-\frac{\omega^2}{2a}}. \quad \text{See plenary exercise}$$

"The Fourier transform of a Gaussian is again a Gaussian"

Important rules for computing the Fourier transform

Theorem (linearity of $\mathcal{F}$)

Let $a, b \in \mathbb{C}$ and $f, g \in L^1(\mathbb{R})$. Then

$$\mathcal{F}(af + bg) = a \mathcal{F}(f) + b \mathcal{F}(g)$$

Proof. Exercise, follows directly from linearity of the integral.

Theorem

Assume that both f and $f' \in L^1$ and that $f(x) \rightarrow 0$ as $x \rightarrow \pm\infty$.

Then it holds that

$$\mathcal{F}(f') = i\omega \mathcal{F}(f).$$

Proof:

$$\frac{1}{\sqrt{2\pi}} \mathcal{F}(f') = \lim_{L \rightarrow \infty} \int_{-L}^L f'(x) e^{-i\omega x} dx$$

integration by part

$$= \lim_{L \rightarrow \infty} \left(f(x) e^{-i\omega x} \Big|_{-L}^L - \int_{-L}^L f(x) (-i\omega) e^{-i\omega x} dx \right)$$

$$= i\omega \int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx = \sqrt{2\pi} i\omega \mathcal{F}(f).$$

Definition (convolution)

Let $f, g \in L^1(\mathbb{R})$. Then the convolution $f * g$

$$(f * g)(x) := \int_{-\infty}^{\infty} f(y) g(x-y) dy$$

is well-defined

and $f * g \in L^1(\mathbb{R})$.

Remark: That $f * g$ is well-defined needs actually a little proof which we skip here.

Theorem (commutativity of the convolution).

$$(f * g) = (g * f).$$

Proof: Similar as in the 2π periodic case.

Theorem

Let $f, g \in L^1(\mathbb{R})$. Then

$$\mathcal{F}(f * g) = \sqrt{2\pi} \cdot \mathcal{F}(f) \cdot \mathcal{F}(g)$$

We can also write this as

$$(f * g)^\wedge(\omega) = \sqrt{2\pi} \hat{f}(\omega) \cdot \hat{g}(\omega).$$

Remark: Last theorem tells us that on the ω -side / in the frequency domain, convolution translates into multiplication.