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by

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Order conditions for the semi-Lagrangian exponential integrators

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1 Introduction

Semi-Lagragian methods have been shown [1, 4, 2] to play an important role in the computation of flows of vector fields in exponential integrators designed for convection dominated problems of the convection-diffusion type. In this paper we examine some of the issues regarding the order conditions for the semi-Lagrangian exponential integrators, starting with a preliminary work by the authors in [2].

Suppose from the semi-discretization of a convection-diffusion model one obtains an ordinary differential equation (with initial data y_0) of the form

$$y_t = C(y)y + Ay, \quad y(0) = y_0,$$
(1.1)

with $y = y(t) \in \mathbb{R}^N$ for $t \in [0, T]$. The $N \times N$ matrices C(y) and A represents the discrete convection and linear diffusion operators respectively.

The methods then take the following general format

for
$$i = 1 : s$$
 do

$$Y_i = \varphi_i y_n + h \sum_j a_{i,j} \varphi_{i,j} A Y_j,$$

$$\varphi_i = \exp(h \sum_k \alpha_{i,J}^k C(Y_k)) \dots \exp(h \sum_k \alpha_{i,1}^k C(Y_k)),$$

$$\varphi_{i,j} = \varphi_i \varphi_j^{-1}$$
end

$$y_{n+1} = \varphi_{n+1} y_n + h \sum_i b_i \varphi_{n+1,i} A Y_i,$$

$$\varphi_{n+1} = \exp(h \sum_k \beta_J^k C(Y_k)) \dots \exp(h \sum_k \beta_1^k C(Y_k)),$$

$$\varphi_{n+1,i} = \varphi_{n+1} \varphi_i^{-1},$$

where $\{a_{i,j}, b_i\}$ are coefficients of a *s*-stage Runge-Kutta (RK) method and $\alpha_{i,l}^j$ and β_l^j are coefficients of a commutator-free (CF) Lie group method (studied in [3, 6]) defined on a RK method with coefficients $\{\hat{a}_{i,j}, \hat{b}_i\}$ such that

$$\hat{a}_{i,j} = \sum_{l=1}^{J} \alpha_{i,l}^{j}, \quad \hat{b}_{i} = \sum_{l=1}^{J} \beta_{l}^{j}.$$
 (1.2)

Thus given a partition RK method with Butcher tableaus

$$\begin{array}{c|c} \mathbf{c} & \mathcal{A} \\ \hline & \mathbf{b} \end{array}, \qquad \begin{array}{c|c} \hat{\mathbf{c}} & \hat{\mathcal{A}} \\ \hline & \hat{\mathbf{b}} \end{array}, \tag{1.3}$$

we treat the diffusion with the *s*-stage RK method $\{\mathcal{A}, \mathbf{b}, \mathbf{c}\}$ (preferably implicit) and the convection with a CF method based on the RK $\{\hat{\mathcal{A}}, \hat{\mathbf{b}}, \hat{\mathbf{c}}\}$.

Note: Here and in the rest of the literature, we shall write \sum_{j} (without explicit limits of summation) to actually mean $\sum_{j=1}^{s}$. All tables have been put in the appendix.

In the study of the under conditions we treat (for the sake of convenience) the numerical solution y_{n+1} as an extra state value

$$Y_{s+1} = \varphi_{s+1}y_n + h\sum_j a_j\varphi_{s+1,j}AY_j, \quad a_{s+1,j} = b_j,$$

with

$$\varphi_{s+1} = \exp(h\sum_{k} \alpha_{s+1,J}^k C(Y_k)) \dots \exp(h\sum_{k} \alpha_{s+1,1}^k C(Y_k)), \quad \alpha_{s+1,l}^k = \beta_l^k$$

and $\varphi_{s+1,j} = \varphi_{s+1}\varphi_j^{-1}$.

2 Deriving the order conditions

Taking the q^{th} derivatives with respect to h of the exact solution to (1.1) and of the stage values of the numerical solution we obtain the recursive formulas

$$y^{(q)} = \sum_{k=0}^{q-1} {\binom{q-1}{k}} \frac{d^{q-1-k}}{dh^{q-1-k}} \left(C(y) + A\right) y^{(k)},\tag{2.1}$$

$$Y_i^{(q)} = \varphi_i^{(q)} y_0 + q \sum_j a_{i,j} \sum_{k=0}^{q-1} {\binom{q-1}{k}} \varphi_{i,j}^{(q-1-k)} A Y_j^{(k)}.$$
 (2.2)

Order conditions for order $p = 1, 2, 3, \ldots$ are recursively from the equations

$$Y_{s+1}^{(q)}|_{h=0} = y^{(q)}|_{h=0}, \quad q = 0, 1, \dots, p.$$
(2.3)

We often will simplify higher order conditions using conditions of lower order whenever neccessary. The computation of the derivatives of Y_i requires the use of φ_i and $\varphi_{i,j}$ and their derivatives.

Now let us consider the matrix-valued functions,

$$C_{i,J-l} := h \sum_{k} \alpha_{i,J-l}^{k} C(Y_k), \quad l = 0, 1, \dots, J-1,$$

and

$$\tilde{C}_{i,J-l} := -h \sum_{k} \alpha_{i,l+1}^{k} C(Y_k), \quad l = 0, 1, \dots, J-1$$

We donote by $B_{i,J-l}$ either of $C_{i,J-l}$ or $\tilde{C}_{i,J-l}$, for $l = 0, 1, \ldots, J-1$, and consider

$$\psi_i(h) := \exp(B_{i,J}) \cdot \exp(B_{i,J-1}) \cdot \ldots \cdot \exp(B_{i,1}).$$

Depending on the choice of $B_{i,J-l} = C_{i,J-l}$ or $\tilde{C}_{i,J-l}$, we have $\psi_i = \varphi_i$ or φ_i^{-1} , respectively. We will also make use of

$$\varphi_i^l(h) := \exp(B_{i,J}) \cdot \exp(B_{i,J-1}) \cdot \ldots \cdot \exp(B_{i,J-l}).$$

We obtain¹

$$\dot{\psi}_i = \sum_{l=0}^{J-1} \operatorname{Ad}_{\psi_i^l} \left(\operatorname{dexp}_{B_{i,J-l}}(\dot{B}_{i,J-l}) \right) \cdot \psi_i$$

So we can write

$$\dot{\psi}_i = S_i(h)\psi_i$$
 with $S_i(h) := \sum_{l=0}^{J-1} \operatorname{Ad}_{\psi_i^l} \left(\operatorname{dexp}_{B_{i,J-l}}(\dot{B}_{i,J-l}) \right)$,

and as a direct consequence we have

$$\psi_i^{(r)} = \sum_{k=0}^{r-1} \binom{r-1}{k} \left(\frac{d^{r-1-k}}{dh^{r-1-k}} S_i(h) \right) \psi_i^{(k)}.$$
(2.4)

Now we have the following proposition for finding the derivatives of $S_i(h)$:

Proposition 2.1. Given that $Z^0 = Z^0(h)$ and W = W(h) are two matrix-valued differentiable functions then

$$\frac{d^r}{dh^r}Ad_WZ^0 = Ad_WZ^r,$$
(2.5)

with

$$Z^{r} = [W^{-1}\dot{W}, Z^{r-1}] + \dot{Z}^{r-1}.$$
(2.6)

The proof is by induction.

By differentiating from (2.6) we obtain

$$\dot{Z}^{r} = [\dot{W}^{-1}\dot{W} + W^{-1}\ddot{W}, Z^{r-1}] + [W^{-1}\dot{W}, \dot{Z}^{r-1}] + \ddot{Z}^{r-1},$$
(2.7)

and using (2.6) and (2.7), assuming W(0) = I, we obtain

$$\begin{cases} \frac{d}{dh} \operatorname{Ad}_{W} Z^{0}|_{h=0} = Z^{1}(0) = [\dot{W}(0), Z^{0}(0)] + \dot{Z}^{0}(0) \\ \dot{Z}^{1}(0) = [-\dot{W}(0)^{2} + \ddot{W}(0), Z^{0}(0)] + [\dot{W}(0), \dot{Z}^{0}(0)] + \ddot{Z}^{0}(0) \\ \frac{d^{2}}{dh^{2}} \operatorname{Ad}_{W} Z^{0}|_{h=0} = Z^{2}(0) = [\dot{W}(0), Z^{1}(0)] + \dot{Z}^{1}(0) \\ \dot{Z}^{2}(0) = [-\dot{W}(0)^{2} + \ddot{W}(0), Z^{1}(0)] + [\dot{W}(0), \dot{Z}^{1}(0)] + \ddot{Z}^{1}(0) \\ \frac{d^{3}}{dh^{3}} \operatorname{Ad}_{W} Z^{0}|_{h=0} = Z^{3}(0) = [\dot{W}(0), Z^{2}(0)] + \dot{Z}^{2}(0) \\ \vdots \end{cases}$$

$$(2.8)$$

¹We recall that $\operatorname{dexp}_{w}(u) := \frac{e^{z}-1}{z} \Big|_{z=\operatorname{ad}_{w}} (u) = u + 1/2! [w, u] + 1/3! [w, [w, u]] + \dots, \operatorname{ad}_{w}(u) := [w, u]$ (matrix commutator of w and u) and $\operatorname{Ad}_{\psi}(u) := \psi u \psi^{-1}$.

Further assuming that $Z^0 = \text{dexp}_{-B}(\dot{B})$ for some matrix-valued function B = B(h), expanding the right-hand side and differentiating we obtain

$$\begin{cases} Z^{0}(0) = \dot{B}(0) \\ \dot{Z}^{0}(0) = \ddot{B}(0) \\ \ddot{Z}^{0}(0) = \ddot{B}(0) - \frac{1}{2}[\dot{B}(0), [\dot{B}(0), \ddot{B}(0)]] \\ \ddot{Z}^{0}(0) = B^{iv}(0) - [\dot{B}(0), \ddot{B}(0)] + \frac{1}{2}[\dot{B}(0), \ddot{B}(0)] \\ \vdots \end{cases}$$
(2.9)

We can now obtain derivatives of S_i and ψ_i . By setting $W = \psi_i^l$ and $B = B_i^{J-l}$ we can calculate the derivatives of S_i using the steps in (2.5)-(2.9). We obtain

$$\begin{cases} S_{i}(0) = \sum_{l=0}^{J-1} \dot{B}_{i}^{J-l}(0), \\ \frac{dS_{i}}{dh}\Big|_{h=0} = \sum_{l=0}^{J-1} \sum_{r=0}^{J-l-1} ([\dot{B}_{i}^{J-r}(0), \dot{B}_{i}^{J-l}(0)] + \ddot{B}_{i}^{J-l}(0)), \\ \frac{d^{2}S_{i}}{dh^{2}}\Big|_{h=0} = 2 \sum_{l=0}^{J-1} \sum_{r=0}^{J-l-1} [\dot{B}_{i}^{J-r}(0), \ddot{B}_{i}^{J-l}(0)] + \sum_{l=0}^{J-1} \sum_{r=0}^{J-l-1} [\ddot{B}_{i}^{J-r}(0), \dot{B}_{i}^{J-l}(0)] + \\ + \sum_{l=0}^{J-1} (\ddot{B}_{i}^{J-r}(0) - \frac{1}{2} [\dot{B}_{i}^{J-l}(0), \ddot{B}_{i}^{J-l}(0)]), \\ \vdots \end{cases}$$

$$(2.10)$$

Analogously the derivatives of $S_i^l := \sum_{r=0}^{J-l-1} \operatorname{Ad}_{\psi_i^r} \left(\operatorname{dexp}_{B_{i,J-r}}(\dot{B}_{i,J-r}) \right)$ are obtained as in the forgoing formulae but substituting J-1 as upper index in the external summation with J-l-1. In table 1 we report the values of the derivatives of φ_i and φ_j^{-1} at h = 0, which are obtained from (2.4) and (2.10) by recursion, starting with $\psi_i(0) = I$. In table 2 we report the values of the derivatives of $\varphi_{i,j}$ at h = 0, which are obtained using table 1 and the formula

$$\varphi_{i,j}^{(m)} = \sum_{r=0}^{m} \binom{m}{r} \varphi_i^{(m-r)} (\varphi_j^{-1})^{(r)}.$$
(2.11)

The derivatives of Y_i , reported in table 3, are obtained using the results in tables 1 and 2, and the recursion formula (2.2), starting with $Y_i(0) = y_0$.

3 Order conditions for orders 1-3

We now present a detailed analysis for deriving the third order conditions.

From (2.4) we obtain that

$$\dot{\varphi}_i(0) = \sum_{l=0}^{J-1} \dot{B}_{i,J-l}(0) = \sum_{l=0}^{J-1} \dot{C}_{i,J-l}(0) \sum_{l=0}^{J-1} \sum_k \alpha_{i,J-l}^k C = (\sum_k \hat{a}_{i,k})C.$$
(3.1)

Analogously one computes $\dot{\varphi_j^{-1}}(0)$. These expressions are reported in table 1 and can be used to obtain

q	$\varphi_i^{(q)}(0)$
0	Ι
1	$C\hat{c}_i$
2	$2\sum_{k}\hat{a}_{i,k}C'(\hat{c}_{k}Cy_{0}+c_{k}Ay_{0})+\hat{c}_{i}^{2}C^{2}$
3	$ \begin{array}{l} 4\sum_{l=0}^{J-1}\sum_{r=0}^{J-1-l}\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}[C,C'(\hat{c}_{m}C+c_{m}A)] + \\ 2\sum_{l=0}^{J-1}\sum_{r=0}^{J-1-l}\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}[C'(\hat{c}_{k}C+c_{k}A),C] + \\ -\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{m}\alpha_{i,J-l}^{m}[C,C'(\hat{c}_{m}C+c_{m}A)] + \\ 3\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}C''(\hat{c}_{k}C+c_{k}A,\hat{c}_{k}C+c_{k}A) + \\ 6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{j}\hat{a}_{k,j}C'(C'(\hat{c}_{j}C+c_{j}A)) + \\ 3\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\hat{c}_{k}^{k}C'(C''2) + \\ 6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\hat{c}_{k}^{k}C'(C'(\hat{c}_{j}C+c_{j}A)) + \\ 4\hat{c}_{i}\sum_{k}a_{i,k}(C'(\hat{c}_{k}C+c_{k}A))C + 2\hat{c}_{i}\sum_{k}a_{i,k}C(C'(\hat{c}_{k}C+c_{k}A)) + \hat{c}_{i}^{3}C^{3} \end{array} $
q	$(\varphi_j^{-1})^{(q)}(0)$
0	Ι
1	$-C\hat{c_j}$
2	$-2\sum_{k}\hat{a}_{j,k}C'(\hat{c}_{k}Cy_{0}+c_{k}Ay_{0})+\hat{c}_{j}^{2}C^{2}$
	$ \begin{bmatrix} 4 \sum_{l=0}^{J-1} \sum_{r=0}^{J-l-1} \sum_{k} \alpha_{j,r+1}^{k} \sum_{m} \alpha_{i,r+1}^{m} [C, C'(\hat{c}_{m}C + c_{m}A)] + \\ 2 \sum_{l=0}^{J-1} \sum_{r=0}^{J-l-1} \sum_{k} \alpha_{i,r+1}^{k} \sum_{m} \alpha_{i,r+1}^{m} [C'(\hat{c}_{k}C + c_{k}A), C] + \end{bmatrix} $

Table 1: Derivatives of φ_i and its inverse at h = 0.

$$\dot{Y}_i|_{h=0} = (\sum_k \hat{a}_{i,k})Cy_0 + (\sum_k a_{i,k})Ay_0 = \hat{c}_i Cy_0 + c_i Ay_0,$$
(3.2)

see table 3. Also from (2.1) we obtain

$$\dot{y}(0) = Cy_0 + Ay_0. \tag{3.3}$$

Imposing $\dot{Y}_{s+1}|_{h=0} = \dot{y}(0)$ we obtain the following order conditions for order 1,

$$\sum_{k} \hat{a}_{s+1,k} = 1, \sum_{k} a_{s+1,k} = 1.$$

These correspond to requiring that the two RK methods (1.3) are consistent.

For order 2 from (2.2) we have that

$$\ddot{Y}_i|_{h=0} = \varphi_i^{(2)}|_{h=0} y_0 + 2\sum_j a_{i,j} \left(\varphi_{i,j}^{(1)}(0) + \varphi_{i,j}(0)A\dot{Y}_j(0)\right)$$

q	$arphi_{i,j}^{(q)}(0)$			
0	Ι			
1	$(\hat{c}_i - \hat{c}_j)C$			
2	$2\sum_{k}(\hat{a}_{i,k} - \hat{a}_{j,k})C'(\hat{c}_{k}C + c_{k}A) + (\hat{c}_{i} - \hat{c}_{j})^{2}C^{2}$			
3	$\begin{split} & 4\sum_{l=0}^{J-1}\sum_{r=0}^{J-l-1}(\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}-\sum_{k}\alpha_{i,r+1}^{k}\sum_{m}\alpha_{i,l+1}^{m})[C,C'(\hat{c}_{m}C+c_{m}A)] + \\ & 2\sum_{l=0}^{J-1}\sum_{r=0}^{J-l-1}(\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}-\sum_{k}\alpha_{i,r+1}^{k}\sum_{m}\alpha_{i,l+1}^{m})[C'(\hat{c}_{k}C+c_{k}A),C] + \\ & -\sum_{l=0}^{J-1}(\sum_{k}\alpha_{i,J-l}^{k}\sum_{m}\alpha_{i,J-l}^{m}-\sum_{k}\alpha_{i,l+1}^{k}\sum_{m}\alpha_{i,l+1}^{m})[C,C'(\hat{c}_{m}C+c_{m}A)] + \\ & 3\sum_{k}(\hat{a}_{i,k}-\hat{a}_{j,k})C''(\hat{c}_{k}C+c_{k}A,\hat{c}_{k}C+c_{k}A) + \\ & 6\sum_{k}(\hat{a}_{i,k}\sum_{m}\hat{a}_{k,m}-\hat{a}_{j,k}\sum_{m}\hat{a}_{k,m})C'(C'(\hat{c}_{m}C+c_{m}A)) + \\ & 3\sum_{k}(\hat{a}_{i,k}-\hat{a}_{j,k})\hat{c}_{k}^{2}C'(C^{2}) + \\ & 6\sum_{k}(\hat{a}_{i,k}-\hat{a}_{j,k})\sum_{m}a_{k,m}C'(A(\hat{c}_{m}C+c_{m}A)) + \\ & 4\sum_{k}(\hat{c}_{i}\hat{a}_{i,k}+\hat{c}_{j}\hat{a}_{j,k})C'(\hat{c}_{k}C+c_{k}A)C + 2\sum_{k}(\hat{c}_{i}\hat{a}_{i,k}+\hat{c}_{j}\hat{a}_{j,k})CC'(\hat{c}_{k}C+c_{k}A) + \\ & (\hat{c}_{i}^{3}-\hat{c}_{j}^{3})C^{3}+3(\hat{c}_{i}\hat{c}_{j}^{2}-\hat{c}_{j}\hat{c}_{i}^{2})C^{3}-6\hat{c}_{j}\sum_{k}\hat{a}_{i,k}C'(\hat{c}_{k}C+c_{k}A)C + \\ & -6\hat{c}_{i}\sum_{k}\hat{a}_{j,k}CC'(\hat{c}_{k}C+c_{k}A) \end{split}$			

Table 2: Derivatives of $\varphi_{i,j}$ at h = 0.

with $\varphi_{i,j}(0) = I$ and $\varphi_{i,j}^{(1)}(0) = \dot{\varphi}_i(0) - \dot{\varphi}_j(0) = (\hat{c}_i - \hat{c}_j)C$. Using $\dot{Y}_j(0)$ from table 3 we obtain

$$\ddot{Y}_i|_{h=0} = \varphi_i^{(2)}|_{h=0} y_0 + 2\sum_j a_{i,j}((\hat{c}_i - \hat{c}_j)CAy_0 + \hat{c}_jACy_0 + c_jA^2y_0).$$
(3.4)

From (2.4) and (2.10) we obtain

$$\varphi_{i}^{(2)}|_{h=0} = \frac{dS_{i}(h)}{dh}\varphi_{i}(h)|_{h=0} + (S_{i}(h)^{2}\varphi_{i}(h))|_{h=0}
= 2\sum_{k}\hat{a}_{i,k}\hat{c}_{k}C'(Cy_{0}) + 2\sum_{k}\hat{a}_{i,k}c_{k}C'(Ay_{0}) + (\sum_{j}\hat{a}_{i,j})^{2}C^{2},$$
(3.5)

reported in table 1. Substituting the results in (3.4) we obtain

$$\ddot{Y}_{i}|_{h=0} = (2\sum_{k} \hat{a}_{i,k} \hat{c}_{k} C'(Cy_{0}) + 2\sum_{k} \hat{a}_{i,k} c_{k} C'(Ay_{0}) + c_{i}^{2} C^{2}) y_{0} + 2\sum_{j} a_{i,j} ((\hat{c}_{i} - \hat{c}_{j}) CAy_{0} + \hat{c}_{j} ACy_{0} + c_{j} A^{2} y_{0},$$

$$(3.6)$$

reported in table 3. Using (2.1) and substituting for $\dot{y}(0)$ from (3.3) we obtain

$$y^{(2)}|_{h=0} = C'(\dot{y}(0))y_0 + (C+A)^2y_0 = C'((C+A)y_0)y_0 + (C+A)^2y_0, \qquad (3.7)$$

where C'(y)(w) is obtained by differentiating C(y) such that

$$(C'(y)(w))_{i,j} := \sum_{k=1}^{N} \frac{\partial c_{i,j}}{\partial y^k} (y) w_k, \quad c_{i,j} = (C(y))_{i,j}, \quad y = [y^1, \dots, y^N]^T.$$

Taking i = s + 1 and matching coefficients in $\ddot{Y}_i|_{h=0}$ and $y^{(2)}|_{h=0}$ we obtain the four order conditions for order 2,

$$\sum_{j} \hat{a}_{s+1,j} \hat{c}_{j}, \quad \sum_{j} a_{s+1,j} c_{j},$$
$$\sum_{j} a_{s+1,j} \hat{c}_{j}, \quad \sum_{j} \hat{a}_{s+1,j} c_{j}.$$

 $Y_{i}^{(q)}(0)$ q0 y_0 $(\sum_{j} \hat{a}_{i,j})Cy_0 + (\sum_{j} a_{i,j})Ay_0$ 1 $2\sum_{j}\hat{a}_{i,j}C'(\hat{c}_{j}C+c_{j}A)y_{0}+\hat{c}_{i}^{2}C^{2}y_{0}+2\hat{c}_{i}c_{i}CAy_{0}-2(\sum_{j}a_{i,j}\hat{c}_{j})CAy_{0}+2\sum_{j}a_{i,j}A(\hat{c}_{j}C+c_{j}A)y_{0}+2\sum_{j}a_{i,j}A(\hat{c}_{j}C+c_{j}A)y_{0}+2\sum_{j}a_{j}A(\hat$ 2 $\frac{1}{4\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}[C,C'(\hat{c}_{m}C+c_{m}A)]y_{0}}$ $2\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}[C'(\hat{c}_{k}C+c_{k}A),C]y_{0}+$ $-\sum_{l=0}^{J-1}\sum_{i=1}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{r}\alpha_{i,J-l}^{r}[C,C'(\hat{c}_{r}C+c_{r}A)]y_{0}+$ $3\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}C''(\hat{c}_{k}C+c_{k}A,\hat{c}_{k}C+c_{k}A)y_{0}+$ $6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{j}\hat{a}_{k,j}C'(C'(\hat{c}_{j}C+c_{j}A))y_{0}+$ $3\sum_{l=0}^{J-1}\sum_{k}^{k}\alpha_{i,J-l}^{k}\hat{c}_{k}^{2}C'(C"2)y_{0} +$
$$\begin{split} & 3\sum_{l=0}^{j-1}\sum_{k}\alpha_{i,J-l}^{\kappa}c_{k}^{c}C'(C''2)y_{0} + \\ & 6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}(\hat{c}_{k}c_{k}-\sum_{j}a_{k,j}\hat{c}_{j})C'(CA)y_{0} + \\ & 6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{j}a_{k,j}C'(A(\hat{c}_{j}C+c_{j}A))y_{0} + \\ & 4\hat{c}_{i}\sum_{k}a_{i,k}(C'(\hat{c}_{k}C+c_{k}A))Cy_{0} + 2\hat{c}_{i}\sum_{k}a_{i,k}C(C'(\hat{c}_{k}C+c_{k}A))y_{0} + \hat{c}_{i}^{3}C^{3}y_{0} \\ & 6\sum_{k}\hat{a}_{i,k}(\sum_{j}a_{i,j})C'(\hat{c}_{k}C+c_{k}A)Ay_{0} + 3(\sum_{k}\hat{a}_{i,k})^{2}(\sum_{j}a_{i,j})C^{2}Ay_{0} + \\ & 6\sum_{k}\hat{a}_{i,k})(-\sum_{j}a_{i,j}\hat{c}_{j}C^{2}Ay_{0} + \sum_{j}a_{i,j}CA(\hat{c}_{j}C+c_{j}A)y_{0} + \\ & 3\sum_{j}a_{i,j}\hat{c}_{j}^{2}C^{2}Ay_{0} - 6\sum_{j}a_{i,j}\sum_{k}\hat{a}_{j,k}C'(\hat{c}_{k}C+c_{k}A)Ay_{0} - 6\sum_{j}a_{i,j}\hat{c}_{j}CA(\hat{c}_{j}C+c_{j}A)y_{0} + \\ & 6\sum_{j}a_{i,j}\sum_{m}\hat{a}_{j,m}AC'(\hat{c}_{m}C+c_{m}A)y_{0} + 3\sum_{j}a_{i,j}\hat{c}_{j}^{2}AC^{2}y_{0} + \\ & 6\sum_{j}a_{i,j}(\hat{c}_{j}c_{j}-\sum_{m}a_{j,m}\hat{c}_{m})ACAy_{0} + 6\sum_{j}a_{i,j}\sum_{m}a_{j,m}A^{2}(\hat{c}_{m}C+c_{m}A)y_{0} \end{split}$$
3

Table 3: Derivatives of Y_i at h = 0.

Note: The matrix-valued function C = C(y) and it's derivatives are linear with respect to y.

For order 3 we proceed as follows: First from (2.1) we have

$$y^{(3)}|_{h=0} = C''(y_0)(\dot{y}(0), \dot{y}(0))y_0 + C'(y_0)(\ddot{y}(0))y_0 + 2C'(y_0)(\dot{y}(0))(C(y_0) + A)y_0 + (C(y_0) + A)C'(y_0)(\dot{y}(0))y_0 + (C(y_0) + A)^3y_0,$$
(3.8)

where we have used C''(y)(w,z) obtained by differentiating C(y) such that

$$(C''(y))(w,z))_{i,j} := \sum_{k=1}^{N} \sum_{m=1}^{N} \frac{\partial^2 c_{i,j}}{\partial y^k \partial y^m}(y) w_k z_m, \quad c_{i,j} = (C(y))_{i,j}.$$

In short we will write C, C', C'', \ldots for $C(y_0), C'(y_0), C''(y_0), \ldots$ respectively. Substituting for $\dot{y}(0)$ and $\ddot{y}(0)$ from (3.3) and (3.7) respectively, we obtain

$$y^{(3)}|_{h=0} = C''((C+A)y_0, (C+A)y_0)y_0 + C'(C'((C+A)y_0)y_0 + (C+A)^2y_0)y_0 + 2C'((C+A)y_0)(C+A)y_0 + (C+A)C'((C+A)y_0)y_0 + (C+A)^3y_0.$$
(3.9)

We now consider the third derivative of the numerical solution. From (2.2) we obtain

$$Y_{i}^{(3)}\Big|_{h=0} = \varphi_{i}^{(3)}\Big|_{h=0} y_{0} + 3\sum_{j} a_{i,j} \varphi_{i}^{(2)}\Big|_{h=0} Ay_{0} + 6\sum_{j} a_{i,j} \varphi_{i}^{(1)}\Big|_{h=0} A\dot{Y}_{j}\Big|_{h=0} + 3\sum_{j} a_{i,j} A \ddot{Y}_{j}\Big|_{h=0}.$$
(3.10)

We need to find $\varphi_i^{(3)}|_{h=0}$ via (2.4) using the expressions for $\varphi_i^{(2)}|_{h=0}$ and $\varphi_i^{(1)}|_{h=0}$ which have already been found and reported in table 1. Using earlier row entries of table 1 and (2.11) we also compute $\varphi_{i,j}^{(2)} = \varphi_i(0) + 2\ddot{\varphi}_i(0)\dot{\varphi}_j^{-1}(0) + \dot{\varphi}_j^{-1}(0)$, reported in table 2. From (2.4) it follows that

$$\varphi_i^{(3)}\Big|_{h=0} = \left. \frac{d^2 S_i}{dh^2} \right|_{h=0} + 2 \left. \frac{dS_i}{dh} \right|_{h=0} \varphi_i^{(1)} \Big|_{h=0} + S_i \Big|_{h=0} \left. \varphi_i^{(2)} \right|_{h=0}.$$
(3.11)

We obtain

$$\varphi_{i}^{(3)}\Big|_{h=0} = \frac{d^{2}S_{i}}{dh^{2}}\Big|_{h=0} + 2(2\sum_{k}\hat{a}_{i,k}\hat{c}_{k}C'(Cy_{0}) + 2\sum_{k}\hat{a}_{i,k}c_{k}C'(Ay_{0}))\sum_{k}\hat{a}_{i,k}C + 2\sum_{k}\hat{a}_{i,k}C(\sum_{k}\hat{a}_{i,k}\hat{c}_{k}C'(Cy_{0}) + \sum_{k}\hat{a}_{i,k}c_{k}C'(Ay_{0})) + \sum_{k}\hat{a}_{i,k}C(\sum_{j}\hat{a}_{i,j})^{2}C^{2}.$$

$$(3.12)$$

We have

$$\dot{C}_{i,J-l}(0) = \sum_{k} \alpha_{i,J-l}^{k} C,$$

$$\ddot{C}_{i,J-l}(0) = 2 \sum_{k} \alpha_{i,J-l}^{k} C'(\hat{c}_{k}C + c_{k}A),$$

$$\ddot{C}_{i,J-l}(0) = 3 \sum_{k} \alpha_{i,J-l}^{k} C''(\hat{c}_{k}C + c_{k}, \hat{c}_{k}C + c_{k}) + C'(\ddot{Y}_{k}(0)).$$

(3.13)

We use (2.10) to find $\frac{d^2S_i}{dh^2}|_{h=0}$, setting $B_{i,J-l} = C_{i,J-l}$, and using the derivatives computed in (3.13) with $\ddot{Y}_k(0)$ from table 3.

Finally we get

$$\frac{d^{2}S_{i}}{dh^{2}}|_{h=0} = 4\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}[C,C'(\hat{c}_{m}C+c_{m}A)] + 2\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\alpha_{i,J-r}^{k}\sum_{m}\alpha_{i,J-l}^{m}[C'(\hat{c}_{k}C+c_{k}A),C] + 3\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}C''(\hat{c}_{k}C+c_{k}A,\hat{c}_{k}C+c_{k}A) + 6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{j}\hat{a}_{k,j}C'(C'(\hat{c}_{j}C+c_{j}A)) + 3\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\hat{c}_{k}^{2}C'(C^{2}) + 6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\hat{c}_{k}c_{k}-\sum_{j}a_{k,j}\hat{c}_{j})C'(CA) + 6\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{j}a_{k,j}C'(A(\hat{c}_{j}C+c_{j}A)) + -\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{j}a_{k,j}C'(A(\hat{c}_{j}C+c_{j}A)) + -\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}^{k}\sum_{j}a_{k,j}C'(A(\hat{c}_{j}C+c_{j}A)) + -\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}\sum_{j}\alpha_{i,J-l}\sum_{j}\alpha_{i,J-l}C'(A(\hat{c}_{j}C+c_{j}A)) + -\sum_{l=0}^{J-1}\sum_{k}\alpha_{i,J-l}\sum_{j}\alpha_{i,J-l}\sum_{j}\alpha_{i,J-l}C'(A(\hat{c}_{j}C+c_{j}A))].$$

Using (3.11) we obtain $\varphi_i^{(3)}(0)$ as reported in table 1 and from (3.10) we obtain $Y_i^{(3)}(0)$ reported in table 3.

By imposing $Y_{s+1}^{(3)}|_{h=0} = y^{(3)}|_{h=0}$ we obtain the conditions for order 3 (recalling that $\alpha_{s+1,J-l}^k = \beta^k$ and $\sum_{l=0}^{J-1} \beta_{J-l}^k = \hat{b}_k$) reported in table 4.

condition	elementary differential			
commutators				
$4\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\beta_{J-r}^{k}\sum_{m}\beta_{J-l}^{m}\hat{c}_{m} + -2\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\beta_{J-r}^{k}\sum_{m}\beta_{J-l}^{m}\hat{c}_{k} + -\sum_{l=0}^{J-1}\sum_{k}\beta_{J-l}^{k}\sum_{r}\beta_{J-l}^{r}\hat{c}_{r} = 0$	[C'(C), C]			
$4\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\beta_{J-r}^{k}\sum_{m}\beta_{J-l}^{m}c_{m} + -2\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\beta_{J-r}^{k}\sum_{m}\beta_{J-l}^{m}c_{k} + -\sum_{l=0}^{J-1}\sum_{k}\beta_{J-l}^{k}\sum_{r}\beta_{J-l}^{r}c_{r} = 0$	[C'(A),C]			
higher order differentials				
$3\sum_{k} b_k \hat{c}_k^2 = 1$	C''(C,C)			
$3\sum_{k} b_k \hat{c}_k c_k = 1$	C''(C,A)			
$3\sum_k \hat{b}_k \hat{c}_k c_k^2 = 1$	C''(A, A)			
$3\sum_{k}\hat{b}_{k}c_{k}\hat{c}_{k} = 1$	C''(A,C)			
$6\sum_{k}\hat{b}_{k}\sum_{j}\hat{a}_{k,j}\hat{c}_{j} = 1$	C'(C'(C))			
$6\sum_{k}\hat{b}_{k}\sum_{j}\hat{a}_{k,j}c_{j}=1$	C'(C'(A))			
$3\sum_{k}^{n} \hat{b}_{k} \hat{c}_{k}^{2} = 1$	$C'(C^2)$			
$6\sum_{k=1}^{n} \hat{b}_{k}(\hat{c}_{k}c_{k} - \sum_{j} a_{k,j}\hat{c}_{j}) = 1$	C'(CA)			
$6\sum_{k}^{n} \hat{b}_{k} \sum_{j} a_{k,j} \hat{c}_{j} = 1$	C'(AC)			
$6\sum_{k}^{n}\hat{b}_{k}\sum_{j}^{n}a_{k,j}c_{j}=1$	$C'(A^2)$			
products of lower order differentials				
$\overline{6\sum_{k}\hat{b}_{k}(\sum_{j}b_{j})\hat{c}_{k}-6\sum_{j}b_{j}\sum_{k}\hat{a}_{j,k}\hat{c}_{k}=2}$	C'(C)A			
$\int \frac{1}{2} \sum_{k} \hat{b}_{k} (\sum_{j} b_{j}) c_{k} - 6 \sum_{j} b_{j} \sum_{k} \hat{a}_{j,k} c_{k} = 2$	C'(A)A			
$3 - 6\sum_{i} b_{i}\hat{c}_{i} + 3\sum_{i} b_{i}\hat{c}_{i}^{2} = 1$	$C^2 A$			
$6\sum_{i} b_{i}\hat{c}_{i} - 6\sum_{i} b_{i}\hat{c}_{i}^{2} = 1$	CAC			
$\int \sum_{i} b_{j} c_{i} - 6 \sum_{i} b_{j} \hat{c}_{i} c_{j} = 1$	CA^2			
$\int \sum_{i} b_{i} \sum_{m} \hat{a}_{i,m} \hat{c}_{m} = 1$	AC'(C)			
$\int \sum_{i} b_{i} \sum_{m} \hat{a}_{i,m} c_{m} = 1$	AC'(A)			
$3\sum_{i} b_{i}\hat{c}_{i}^{2} = 1$	AC^2			
$\int \sum_{j=1}^{2} b_j (\hat{c}_j c_j \sum_{m=1}^{2} a_{j,m} \hat{c}_m) = 1$	ACA			
$6\sum_{i}b_{i}\sum_{m}a_{i}\hat{c}_{m}=1$	A^2C			
$6\sum_{j} b_{j} \sum_{m} a_{j,m} c_{m} = 1$	A^3			

Table 4: Conditions of order 3

4 Extra coupling conditions for order 4

We consider coefficients of elementary differentials preceded by an A in both the expressions for the fourth derivatives (q = 4) of the exact and numerical solutions (2.1) and (2.2) respectively. That means matching the terms in $Ay^{(3)}|_{h=0}$ and $4\sum_{j}\sum_{j}a_{i,j}\varphi_{i,j}(0)AY_{j}^{(3)}(0)$,

since $\varphi_{i,j}(0) = I$. We obtain

$$\begin{aligned} Ay^{(3)}\Big|_{h=0} &= A[C''(C,C) + C''(C,A) + C''(A,C) + C''(A,A) + C'(C'(C)) + C'(C'(A)) \\ &+ C'(C^2) + C'(CA) + C'(AC) + C'(A^2) + 2C'(C)C + 2C'(C)A + 2C'(A)C \\ &+ 2C'(A)A + CC'(C) + CC'(A) + AC'(C) + AC'(A) + (C+A)^3]y_0. \end{aligned}$$

$$(4.1)$$

We substitute for $\varphi_{i,j}(0)$ and $Y_j^{(3)}(0)$ in $4\sum_j \sum_j a_{i,j}\varphi_{i,j}^{(0)}AY_j^{(3)}(0)$, and select elementary differentials whose coefficients contain the CF coefficients $\alpha_{s+1,J-l}^k := \beta_{J-l}^k$. These include $ACC'(C)y_0$, $ACC'(A)y_0$, $AC'(C)Cy_0$, $AC'(A)Cy_0$, arising from the terms

$$16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}A[C,C'(\hat{c}_{m}C+c_{m}A)]y_{0} + \\8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}A[C,C'(\hat{c}_{k}C+c_{k}A)]y_{0} + \\-4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}A[C,C'(\hat{c}_{r}C+c_{r}A)]y_{0} + \\16\sum_{j}\sum_{k}a_{i,j}\hat{c}_{j}\hat{a}_{j,k}AC'(\hat{c}_{k}C+c_{k}A)Cy_{0} + 8\sum_{j}\sum_{k}a_{i,j}\hat{c}_{j}\hat{a}_{j,k}ACC'(\hat{c}_{k}C+c_{k}A)y_{0}$$

$$(4.2)$$

Comparing coefficients of elementary differentials $ACC'(C)y_0$, $ACC'(A)y_0$, $AC'(C)Cy_0$, $AC'(A)Cy_0$ we obtain respectively the order conditions

$$16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{m} - 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{k}$$

$$-4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}\hat{c}_{r} + 8\sum_{j}\sum_{k}a_{i,j}\hat{c}_{j}\hat{a}_{j,k}\hat{c}_{k} = 1,$$

$$(4.3)$$

$$16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{m} - 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{k} - 4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}c_{r} + 8\sum_{j}\sum_{k}a_{i,j}\hat{c}_{j}\hat{a}_{j,k}c_{k} = 1,$$

$$(4.4)$$

$$-16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{m} + 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{k} + 4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}\hat{c}_{r} + 16\sum_{j}\sum_{k}a_{i,j}\hat{c}_{j}\hat{a}_{j,k}\hat{c}_{k} = 2,$$

$$(4.5)$$

$$-16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{m} + 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{k} + 4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}c_{r} + 16\sum_{j}\sum_{k}a_{i,j}\hat{c}_{j}\hat{a}_{j,k}c_{k} = 1,$$

$$(4.6)$$

Simplifying (4.3)-(4.6) via the use of third order conditions we obtain the order conditions (4.7)-(4.10). Assuming that the RK tableaus (1.3) fulfill the order conditions for a classical partitioned RK method of order 4, and that the b's are different from the \hat{b} 's, i.e., $b_j \neq \hat{b}_j$ for some $j = 1, \ldots, s$, then the order conditions in (4.7)-(4.10) will result in a new or extra set of coupling conditions (involving the α and β coefficients) for our method, which are not included in the set of order conditions for the classical particular RK methods and CF methods. This is not the case for orders 1 through 3 where all our order conditions are only a subset of those for the particular RK and CF methods.

$$16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{m} - 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{k} - 4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}\hat{c}_{r} = 0,$$

$$(4.7)$$

$$16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{m} - 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{k} - 4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}c_{r} = 0,$$

$$(4.8)$$

$$-16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{m} + 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}\hat{c}_{k} + 4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}\hat{c}_{r} = 0,$$

$$(4.9)$$

$$-16\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{m} + 8\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}\sum_{m}a_{i,j}\alpha_{j,J-r}^{k}\alpha_{j,J-l}^{m}c_{k} + 4\sum_{j}\sum_{l=0}^{J-1}\sum_{r=0}^{l}\sum_{k}a_{i,j}\alpha_{j,J-l}^{k}\alpha_{j,J-l}^{r}c_{r} = 0,$$

$$(4.9)$$

where i = s + 1 so that $a_{i,j} = a_{s+1,j} = b_j$, j = 1, ..., s.

In figure 1(a) we have a numerical test showing the order in time for a fourth order method (DIRK-CF4). This method is constructed using the additive particulation IMEX RK method of Kennedy and Carpenter [5] (here named as IMEX4) wherein we derive from the corresponding explicit tableau the commutator-free coefficients via (1.2) satisfying the

CF order conditions as described by Owren [6]. It is mportant to note, however, that our new method DIRK-CF4 automatically satisfy the coupling conditions (4.7)-(4.10) as part of the particle RK order conditions since in this choice of IMEX RK scheme the b's and \hat{b} 's are the same (see [5]). The figure also shows a comparison between the DIRK-CF4 and its counterpart IMEX4. The numerical experiment is performed on the viscous Burgers equation

$u_t + uu_x = \nu u_{xx}$

over a domain [0,1] with initial condition $u(x,0) = u_0(x) = \sin(\pi x)$, and Dirichlet homogenoues boundary conditions. We integrate on the interval [0,T] $(T = 1, \nu = 0.05$ in this case) with time steps in the range $\{\Delta t = 2^{-n} | n = 4, 5, \dots, 9\}$. The spatial discretization is the standard centered differences on a uniform grid of mesh step $\Delta x = 1/32$. The error is measured as a grid-point error in the 2-norm, and the reference (exact) solution is computed as in [2]. Figure 1(b) shows the numerical order tests performed for some of the first to third order methods derived in [2].



Figure 1: Numerical order tests using Burgers equation, $u_t + uu_x = \nu u_{xx}$, on Dirichlet homogeneous BCs, $x \in [0, 1], \nu = 0.05, u_0 = \sin \pi x, T = 1, \Delta x = 1/32$. Plot of the 2- norm error as a function of $\Delta t = 2^{-n}, n = 4, 5, \ldots, 9$. (a)Order test for the fourth order DIRK-CF4 and IMEX4 (b) Order test for first order DIRK-CF1, second order DIRK-CF2 and DIRK-CF2L, third order DIRK-CF3G.

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