

2 Results on singular SDEs

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Singular Diffusions

↙ distribution

$$dX_t = b(t, X_t) dt + d\text{"Noise"}$$

↙ here: α -stable process, $\alpha \in (1, 2]$
($\alpha = 2$: Brownian motion)

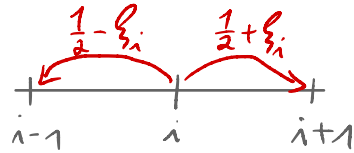
Questions:

- Why are we interested?
- How to solve them?
- Qualitative aspects?

Singular Diffusions as Processes in Random Media

Sinai's random walk in random environment:

$(\xi_j)_{j \in \mathbb{Z}}$ iid, centered, values in $(-\frac{1}{2}, \frac{1}{2})$



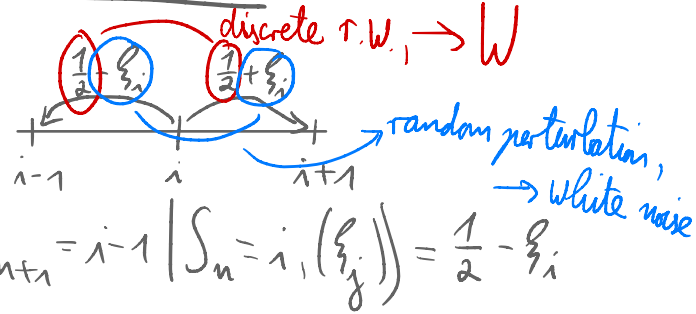
$$\mathbb{P}(S_{n+1} = i+1 | S_n = i, (\xi_j)) = \frac{1}{2} + \xi_i \quad \mathbb{P}(S_{n+1} = i-1 | S_n = i, (\xi_j)) = \frac{1}{2} - \xi_i$$

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Continuous analogue: Brox diffusion

$$dX_t = \xi(X_t)dt + dW_t$$

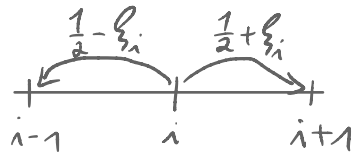
white noise = B'

\leftarrow 1d B.m.

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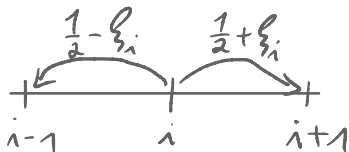
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- scaling limit if ξ_i "small"
- easier to analyze because of scaling invariance
- $\alpha^{-2} X_{e^\alpha}$ converges, $\alpha \rightarrow \infty$

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Continuous analogue: Brox diffusion

$$dX_t = \zeta(X_t) dt + dW_t$$

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1d B.m.

very slow!

- faster for $W \rightarrow$ Lévy noise?
(compare Imbeller - Pavlyukovich '06)
- scaling limit if ξ_i "small"
- easier to analyze because of scaling invariance
- $\alpha^{-2} X_{e^\alpha}$ converges, $\alpha \rightarrow \infty$

Singular diffusions as random polymer measures

Random directed polymer measure

Alberts - Khanin - Quastel '14, Delamare - Diehl '16

$$dP = \frac{1}{Z} \exp \left(\int_0^T \xi(s, B_s) ds \right) dW$$

↑ space-time white noise ↑ Wiener measure

⇒ under P the coordinate process solves

$$dX_t = \partial_x h(t, X_t) dt + dW_t$$

↑ solution to KPZ, non-differentiable

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1d

Cannizzaro - Chauk '18: Similar story for space white noise, $d=1,2,3$.

Singular diffusions as stochastic characteristics I

$$\text{KPZ equ.: } \partial_t h = -\left(\frac{1}{2}\partial_{xx} h + \frac{1}{2}(\partial_x h)^2 + \xi\right), \quad h(T) = \bar{h}$$

Gubinelli-P.17

$$\Rightarrow h(t, x) = \sup_v \mathbb{E}_{t, x} \left[\int_t^T \xi(s, X_s^v) ds + \bar{h}(X_T^v) \right]$$

$$dX_t^v = v_t dt + dW_t$$

$$\text{Optimal } v_t = \partial_x h(t, X_t^v)$$

↖ distributional drift

→ leads to good L^∞ bounds for h

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Expect:

- Feynman-Kac type representations for general linear singular SPDEs (parabolic)
- Nonlinear equ. via BSDEs, McKean-Vlasov, stoch. control

What's new here?

- Extension of Carrizoso - Chouk '18 to Lévy noise
- "Quantitative" heat kernel bounds for singular diffusions in the "Young regime"

Solution theories for singular diffusions

- 1d / gradient drift: Itô - McKean construction / Dirichlet forms

Brox '86 / Mathieu '94

- 1d via martingale problem / Zvonkin transform

Flandoli - Russo - Wolf '03 / Bass - Chen '01

$$\rightarrow dX_t = b(X_t)dt + dW_t$$

$$\Rightarrow \text{formally: } f(X_t) - f(X_0) - \int_0^t \mathcal{L}f(X_s) ds = M_t.$$

$$\hookrightarrow \mathcal{L}f = b \partial_x f + \frac{1}{2} \partial_{xx} f$$

f smooth $\Rightarrow \mathcal{L}f$ distribution

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$$\rightarrow \mathcal{L}f = b \partial_x f + \frac{1}{2} \partial_{xx} f$$

Idea: find non-smooth f s.t. $\mathcal{L}f \in C_b$. f smooth $\Rightarrow \mathcal{L}f$ distribution

E.g. solve $\mathcal{L}f = g$ (linear ODE in 1d)

• $d \geq 1$ and $b \in C(\mathbb{R}_+, C^{-\alpha})$, $\alpha < \frac{1}{2}$ (Young regime):

Flandoli-Issoglio-Russo '17 via regularity analysis for

$$\partial_t u = \mathcal{L}u + f, \quad u(0) = \varphi$$

• $d = 1$ and $b \in C(\mathbb{R}_+, C^{-\alpha})$, $\alpha < \frac{2}{3}$ (rough path regime):

Delarue-Diel '16 via rough path analysis of $\partial_t u = \mathcal{L}u + f$

(similar to Hairer's rough stochastic PDEs)

• $d \geq 1$ and $b \in C(\mathbb{R}_+, C^{-\alpha})$, $\alpha < \frac{2}{3}$ (rough regime) \rightarrow

Cannizzaro-Chouk '18 via paraccontrolled distrib.

$\alpha < 1$ would be
subcritical
(RS?)

Extension of Carrizoso - Chouk to Lévy noise

$$dX_t = b(t, X_t)dt + dL_t$$

$\leftarrow \alpha$ -stable Lévy process, generator $\approx -(-\Delta)^{\frac{\alpha}{2}}$

Formal martingale problem:

$$f(t, X_t) - f(0, X_0) - \int_0^t (\partial_s + \mathcal{L}) f(s, X_s) ds = \text{Mart.}$$

$$= (\partial_s + b(s, x) \cdot \nabla - (-\Delta)^{\frac{\alpha}{2}}) f$$

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\rightarrow Kolmogorov backward equation: $\in C^{-\gamma} \Rightarrow$ guess $u \in C^{\alpha-\gamma} \Rightarrow \nabla u \in C^{\alpha-\gamma-1}$

gains C^α

$$(\partial_t - (-\Delta)^{\frac{\alpha}{2}}) u = -b \cdot \nabla u + f$$

$\Rightarrow b \cdot \nabla u$ ok if $-\gamma + (\alpha - \gamma - 1) \geq 0 \Leftrightarrow \gamma < \frac{\alpha - 1}{2}$

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Young regime,

Athreya - Butkovsky - Mytnik '20

Ling - Zhao '13

de Raynal - Manzeri '13

Lévy Brox would be

$$\xi \in C^{-\frac{1}{2}}, \frac{1}{2} > \frac{\alpha-1}{2}$$

\rightarrow Kolmogorov backward equation: $\in C^{-\gamma} \Rightarrow$ guess $u \in C^{\alpha-\gamma} \Rightarrow \nabla u \in C^{\alpha-\gamma-1}$

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$$(\partial_t - (-\Delta)^{\frac{\alpha}{2}}) u = b \cdot \nabla u + f$$

$\Rightarrow b \cdot \nabla u$ ok if

$$-\gamma + (\alpha - \gamma - 1) > 0 \Leftrightarrow \gamma < \frac{\alpha-1}{2}$$

Paracontrolled ansatz

$$(\partial_t - (-\Delta)^{\alpha/2})u = -\nabla u \cdot b + f = \underbrace{-\nabla u \otimes b + f}_{\in C^{\alpha-2\gamma-1} \text{ (better)}} + \underbrace{-\nabla u \otimes b}_{\in C^{-\gamma}}$$

"paraproducts"

isolates singular part of the product

Thm. (Kremp-P. '20)

For $\gamma < \frac{2\alpha-2}{3}$, $b \in C_T C^{-\gamma}$, $b \circ \nabla (\partial_t - (-\Delta)^{\alpha/2})^{-1} b \in C_T C^{\alpha-2\gamma-1}$:

$\exists!$ paracontrolled solution u . Similar to *Comizzaro-Chouk* with $\Delta \rightarrow -(-\Delta)^{\alpha/2}$

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"paraproducts"

white noise: $\gamma = \frac{1}{2} +$, ok if $\alpha > \frac{7}{4}$. Critical regularity: $\alpha = \frac{3}{2}$ (via RS?)

isolates singular part of the product

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$\exists!$ paracontrolled solution u . Similar to Camillo-Chek with $\Delta \rightarrow -(-\Delta)^{\alpha/2}$

Construction of the process

$$dX_t^n = b_n(t, X_t^n) dt + dL_t$$

$b_n \rightarrow b$ in "rough path norm"

Need: • Tightness of (X^n) .

• Any limit point X is s.t. $\forall u$ solving $(\partial_t - (-\Delta)^{\alpha/2})u = -b \cdot \nabla u + f$:

$$u(t, X_t) - u(0, X_0) - \int_0^t f(s, X_s) ds = \text{Mart.}$$

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Kremp-P. '20: • Tightness via stoch. calculus for Poisson random measures

(Campbell's formula), get $E\left[\left|\int_s^t b_n(r, X_r^n) dr\right|^2\right] \lesssim |t-s|^{1-\frac{1}{\alpha}}$

• Martingale property easy once we know $P(\Delta X_t \neq 0) = 0 \forall t \geq 0$.

Brox diffusion with Lévy noise

$$dX_t = \zeta(X_t) dt + dL_t, \quad d=1$$

↪ take **periodic white noise** (avoid weighted spaces)

• Clear: $\zeta \in C^{-\frac{1}{2}-}$

• Kremp-P. '20: • can construct $\zeta \circ \nabla (\partial_t - (-\Delta)^{\alpha/2})^{-1} \zeta \in C^{\alpha-2-}$

~~no~~ renormalization

if $\alpha > \frac{3}{2}$

$\alpha = \frac{3}{2}$: critical regularity

• get **periodic** Brox diffusion for $\alpha > \frac{7}{4}$.

Next: heat kernel estimates in the Brownian case

Motivation: Mass asymptotics for 2d PAM

$$\partial_t u = \Delta u + u \xi, \quad \xi \text{ 2d white noise}$$

$$U(t) := \int_{\mathbb{R}^d} u(t, x) dx$$

Goal: study $U(t)$ for $t \rightarrow \infty$.
(a.s. limit)

Known results

• ξ smooth Gaussian potential $\Rightarrow \frac{1}{t(\log t)^{\frac{1}{2}}} \log U(t) \rightarrow \sqrt{2d \mathbb{E}[\xi(0)^2]}$

Carmona-Molchanov '95, Gärtner-König-Molchanov '00: 2nd order asymptotics

• $\mathbb{E}[\xi(x)\xi(y)] = |x-y|^{-\alpha}, 0 < \alpha < \min\{2, d\}$

$\Rightarrow \frac{1}{t(\log t)^{\frac{2}{4-\alpha}}} \log U(t) \rightarrow C_\alpha$ Chen '14

• ξ 1d white noise $\Rightarrow \frac{1}{t(\log t)^{2/3}} \log U(t) \rightarrow c$ Chen '14

• $\xi = 2_t B^H$ $\Rightarrow \frac{1}{t(\log t)^{\frac{1}{1+H}}} \log U(t) \rightarrow c$ Chakraborty-Chen - Gao-Tindel '18
fractional B.m., Hurst index $H \in (0, 1)$

$$\frac{1}{t(\log t)^{\frac{1}{2}}}$$

$$\frac{1}{t(\log t)^{\frac{2}{4-\alpha}}}$$

$$\alpha \in (0, 2 \wedge d)$$

$$\frac{1}{t(\log t)^{2/3}}$$

$$\frac{1}{t(\log t)^{\frac{1}{1+H}}}$$

$$H \in (0, 1)$$

Note: Always $\frac{1}{t(\log t)^\gamma} \log U(t) \rightarrow c$, $\gamma \in [\frac{1}{2}, 1)$

$\gamma = 1$ is natural barrier, $\gamma \geq 1$ needs renormalization.

Expect: $\xi \in C^{-k} \Rightarrow \gamma = \frac{1}{2-k} \vee \frac{1}{2}$

$$\leadsto \gamma < 1 \Leftrightarrow k < 1$$

$$\frac{1}{t(\log t)^{\frac{1}{2}}}$$

$$\frac{1}{t(\log t)^{\frac{2}{4-\alpha}}}$$

$$\alpha \in (0, 2 \text{d})$$

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2d white noise:

$\leadsto \gamma < 1 \Leftrightarrow k < 1$

$k = 1$

What breaks at $K_0 = 1$?

$$\partial_t u = \Delta u + u \xi, \quad u(0) = \delta_0$$

All proofs based on Feynman-Kac: $U(t) = \mathbb{E} \left[\exp \left(\int_0^t \xi(w_s) ds \right) \right]$

But: $\xi \in C^{-1-} \Rightarrow$ cannot construct $\int_0^t \xi(w_s) ds$ 

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But: $\xi \in C^{-1-}$ \Rightarrow cannot construct ~~$\int_0^t \xi(w_s) ds$~~ \Downarrow

König-vanZuijlen-P.'20: Still can show

$$\lim_{t \rightarrow \infty} \frac{1}{t \log t} \log U(t) \rightarrow \mathcal{X} \quad \text{for 2d white noise}$$

via "transformed Feynman-Kac formula"

Making the Feynman-Kac formula work

Let $-\Delta Z = \xi$

$$\Rightarrow Z(B_t) = Z(B_0) + \int_0^t \nabla Z(B_s) \cdot dB_s + \underbrace{\int_0^t \Delta Z(B_s) ds}_{= - \int_0^t \xi(B_s) ds}$$

" $\frac{1}{2} = 1$ "
↓

"It's trick"

$$\Rightarrow \int_0^t \xi(B_s) ds = Z(B_0) - Z(B_t) + \int_0^t \nabla Z(B_s) \cdot dB_s$$

ok if $\xi \in C^{-1+\epsilon}$ ($\Rightarrow Z \in C^{1+\epsilon}$)

here: just fails

Making Feynman-Kac work $\int_0^t \xi(B_s) ds = Z(B_0) - Z(B_t) + \int_0^t \nabla Z(B_s) \cdot dB_s$

$$\Rightarrow U(t) = \mathbb{E} \left[\exp \left(Z(B_0) - Z(B_t) + \int_0^t \nabla Z(B_s) \cdot dB_s \right) \right]$$

$$= \mathbb{E} \left[\exp \left(Z(0) - Z(B_t) + \left(\int_0^t \nabla Z(B_s) \cdot dB_s - \frac{1}{2} \int_0^t |\nabla Z(B_s)|^2 ds \right) + \frac{1}{2} \int_0^t |\nabla Z(B_s)|^2 ds \right) \right]$$

Making Feynman-Kac work $\int_0^t \xi(B_s) ds = Z(B_0) - Z(B_t) + \int_0^t \nabla Z(B_s) \cdot dB_s$

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$$= \mathbb{E} \left[\exp \left(Z(0) - Z(X_t) + \frac{1}{2} \int_0^t |\nabla Z(X_s)|^2 ds \right) \right] + \frac{1}{2} \int_0^t |\nabla Z(B_s)|^2 ds$$

$$X_t = \int_0^t \nabla Z(X_s) ds + B_t \quad (\text{weakly})$$

Making Feynman-Kac work $\int_0^t \xi(B_s) ds = Z(B_0) - Z(B_t) + \int_0^t \nabla Z(B_s) \cdot dB_s$

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$$X_t = \int_0^t \nabla Z(X_s) ds + B_t \quad (\text{weakly})$$

$|\nabla Z|^2$ ill-defined ($Z \in C^{1+\varepsilon}$)

Probabilistic argument: construct $|\nabla Z|^2$ as (random) distribution in $C^{-\varepsilon}$

Rigorous Feynman-Kac for 2d PAM

It's trick once more:

$$-\Delta Z = \zeta, \quad (-\Delta - \nabla Z \cdot \nabla) Y = -|\nabla Z|^2, \quad W = Z + Y$$

$$\begin{aligned} \hookrightarrow \epsilon C^{-\epsilon} &\Rightarrow Y \in C^{2-\epsilon} \\ &\Rightarrow |\nabla Y|^2 \text{ ok} \end{aligned}$$

$$dX_t = \nabla W(X_t) dt + dB_t$$

$$\Rightarrow U(t) = \mathbb{E} \left[\exp \left(W(0) - W(X_t) - \frac{1}{2} \int_0^t |\nabla Y|^2(X_s) ds \right) \right]$$

→ Get asymptotics of $U(t)$ if we understand X well.
(upper + lower tail bounds)

Heat kernel estimates for singular diffusions

$$dX_t = b(t, X_t)dt + dB_t \quad b \in C(\mathbb{R}_+, C^{-\alpha}), \alpha < \frac{1}{2}$$

→ not hard to show: \exists transition density

$$\Gamma_{s,t}(x,y) dy = \mathbb{P}(X_t \in dy | X_s = x)$$

Zhang-Zhao '17: $\exists M > 1, \forall T > 0 \exists C > 0$ depending only on T and

$$\|b\|_{C([0,T], C^{-\alpha})}$$

(if $\alpha < \frac{1}{2}$)

$$\frac{1}{C} \gamma\left(\frac{t}{M}, x-y\right) \leq \Gamma_t(x,y) \leq C \gamma(Mt, x-y)$$

↖ standard Gaussian kernel

↙ based on
"Zvonkin
transform"

Heat kernel estimates for singular diffusions

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$$\Gamma_t(x, y) dy = \mathbb{P}_x(X_t \in dy)$$

We need explicit dependence on T , $\|b\|$.

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(if $\alpha < \frac{1}{2}$)

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standard Gaussian kernel

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"Quantitative" heat kernel bounds for singular diffusions

$$dX_t = b(t, X_t)dt + dB_t$$

$$\int_{s,t} (x,y) dy = \mathbb{P}(X_t \in dy | X_s = x)$$

Thm. (van Zuijlen - P. '20)

Young regime
 $\forall \alpha \in (0, \frac{1}{2})$, $c > 1 \quad \exists C > 0, M > 1$:

$$\frac{1}{C} \exp(-C(t-s) \|b\|_{C_t}^{\frac{2}{1-\alpha}}) \eta\left(\frac{t-s}{M}, x-y\right) \leq \int_{s,t} (x,y)$$

$$\leq C \exp(C(t-s) \|b\|_{C_t}^{\frac{2}{1-\alpha}}) \eta(c(t-s), x-y)$$

$$\forall b \in C(\mathbb{R}_+, C^{-\alpha}), \forall 0 \leq s < t$$

Idea of the proof Fix $s=0$, $\Gamma_t := \Gamma_{0,t}$

$$\partial_t \Gamma_t = \underbrace{\frac{1}{2} \Delta \Gamma_t + b \cdot \nabla \Gamma_t}_{\mathcal{L} \Gamma_t}, \quad \Gamma_t = \delta_y$$

Parametrix method: $\partial_t \eta_t = \frac{1}{2} \Delta \eta_t$, $\eta_0 = \delta_y$

compare Friedman's book

$\Rightarrow \tau_t := \Gamma_t - \eta_t$ solves $(\partial_t - \mathcal{L}) \tau_t = \underbrace{-b \cdot \nabla \tau_t}_{=: \Phi_t}$, $\tau_0 = 0$

space-time "convolution"

$$\tau = \Gamma * \Phi$$

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$$\partial_t \Gamma_t = \underbrace{\frac{1}{2} \Delta \Gamma_t + b \cdot \nabla \Gamma_t}_{\mathcal{L} \Gamma_t}, \quad \Gamma_t = \delta_y$$

Parametrix method: $\partial_t \eta_t = \frac{1}{2} \Delta \eta_t$, $\eta_0 = \delta_y$

$\Rightarrow \tau_t := \Gamma_t - \eta_t$ solves $(\partial_t - \mathcal{L}) \tau_t = \underbrace{-b \cdot \nabla \tau_t}_{=: \Phi_t}$, $\tau_0 = 0$

$\tau = \Gamma * \Phi$
space-time "convolution"

$$\Gamma = \eta + \tau = \eta * \Phi + \Gamma * \Phi^{*2}$$

$$\dots \Rightarrow \Gamma = \sum_{n=0}^{\infty} \eta * \Phi^{*n}$$

↪ control each term inductively

$$\Gamma = \sum_{n=0}^{\infty} \gamma * \Phi^{*n}$$

↳ control each term inductively

≈ 8 pages of computations:
get upper bound for $|\Gamma - \gamma|$

$$\Gamma = \sum_{n=0}^{\infty} \gamma * \Phi^{*n}$$

↳ control each term inductively

≈ 8 pages of computations:

get upper bound for $|\Gamma - \gamma|$

Chapman-Kolmogorov

+ "chaining"

⇒ lower bound (cf. Stroock's PDE book)

Summary

- Singular diffusions arise as processes in random media
+ as stochastic characteristics for singular SPDEs.
- Weak solution theory via martingale problem.
- Here 2 new results:
 - Brox diffusion with Lévy noise
 - quant. heat kernel estimates in Brownian Young regime

Thank You